

Solutions

(1)

Shock waves

ex: Burger's equation $u_t + uu_x = 0$ (*)
(inviscid fluid dyn)

$u(c(t), t)$ $x = c(t)$ (a characteristic curve)

$$\frac{d}{dt} u(c(t), t) = u_x(c(t), t) c'(t) + u_t(c(t), t)$$

if $c(t) = f(x_0)t + x_0$ then $u_x(f(x_0)t + x_0, t) f'(x_0) + u_t(f(x_0)t + x_0, t)$
(at $t=0$ $(x_0, 0)$ with $u(x_0, 0) = f(x_0)$) at $t=0$

$$u_x(x_0, 0) f'(x_0) + u_t(x_0, 0) = u_x(x_0, 0) u(x_0, 0) + u_t(x_0, 0) = 0$$

as u solution of (*)

$$\frac{d}{dt} u(c(t), t) = 0$$

value of u constant along char. curve

in this case a wave traveling at speed equal to initial height $f(x_0) = u(x_0, 0)$ of wave



KdV equation again

$$u_t = \frac{3}{2} u u_x + \frac{1}{4} u_{xxx}$$

Note: similarity to Burger's eq nonlin

$u_t + u u_x = 0$ where have shock waves

also similarity to lin eq, $u_t = u_{xxx}$ that has sol's

$$u_k(x, t) = \cos(kx - k^3 t) = \cos(k(x - k^2 t))$$

linear comb'n's
w/ different
frequencies
separate out in
time

dispersive equation
where speed of
traveling wave depends
on frequency k

③ but surprisingly KdV has family of solutions that avoid both problems

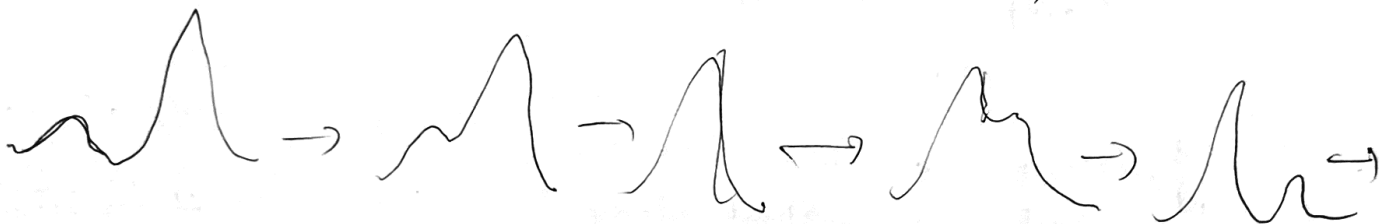
$$u_k(x,t) = \frac{8k^2}{(e^{kx+k^3t} + e^{-kx-k^3t})^2}$$

Solitary wave
traveling at speed k^2
and height $2k^2$

1-soliton $\rightarrow$

n-soliton solution

example $u(x,t) = \frac{24(e^{2x+2t} + 6e^{6x+18t} + 4e^{4x+16t} + 4e^{8x+20t} + e^{10x+25t})}{(1 + 3e^{2x+2t} + e^{6x+18t} + 3e^{4x+16t})^2}$



two interacting waves ... scattering theory of solitons
how multiple waves interact
and bounce off each other
(replaces the simpler superposition of the linear case)

Weierstrass $\wp$ -function

$$y^2 = 4x^3 - k_1x - k_2 \quad \text{ell. curve} \quad (k_1, k_2) \longleftrightarrow \Lambda$$

$$\wp(z; \Lambda, \omega_2) \quad \Lambda = \omega_1\mathbb{Z} + \omega_2\mathbb{Z}$$

$$\wp(z; \Lambda) = \frac{1}{z^2} + \sum_{\lambda \in \Lambda \setminus \{0\}} \left(\frac{1}{(z-\lambda)^2} - \frac{1}{\lambda^2} \right)$$

nonlin ODE

$$(w')^2 = 4w^3 - k_1w - k_2$$

has sol'n

$$w(z) = \wp(z+r; \Lambda(k_1, k_2))$$

parameterization of the elliptic curve

constant of integrals

KdV travelling waves
 $u(x,t) = w(x+ct)$

$$u_t = \frac{3}{2} u u_x + \frac{1}{4} u_{xxx}$$

$$c w' = \frac{3}{2} w w' + \frac{1}{4} w'''$$

integrate

$$c w = \frac{3}{4} w^2 + \frac{1}{4} w'' + r_1$$

mult. by w' & integrate

$$\frac{c}{2} w^2 = \frac{1}{4} w^3 + \frac{1}{8} (w')^2 + r_1 w + r_2$$

$$\Rightarrow (w')^2 = -2w^3 + 4c w^2 - 8r_1 w - 8r_2$$

$$k_1 = \frac{4}{3}(c^2 - 3r_1)$$

$$k_2 = \frac{8c^3}{27} - \frac{4cr_1}{3} - 2r_2$$

$$\Rightarrow u(x,t) = -2\wp(x+ct+w, k_1, k_2) + 2c/3$$

Solution of KdV for any choice of w, k_1, k_2, c

= travelling waves of any speed
 and profile $\wp$ of an arbitrary ell. curve



More general algebro-geometric picture:

choice of any algebraic curve C
 (not only elliptic curve)

$\Rightarrow J(C)$ Jacobian
 (for elliptic curve
 $J(E) \cong E$)

+ choice of pt. on Jacobian

$\Rightarrow$ Solutions of KdV if C
 solution of KP obtained

superelliptic

② a - show (a(x)) for $x \in \mathbb{R}$

b) continuous at x_0

c) $\lim_{x \rightarrow x_0} a(x) = a(x_0)$

d) interval - endpoints for

$$a(x) = \frac{x}{x^2 + 1} \in \mathbb{R}$$

$$\lim_{x \rightarrow x_0} a(x) = \frac{x_0}{x_0^2 + 1}$$

$$a(x_0) = \frac{x_0}{x_0^2 + 1}$$

e) for large x , graph of $a(x)$ has
a local maximum

7. function

$$a(x, y) = \frac{x}{x^2 + y^2} \quad x, y \in \mathbb{R} \quad \text{such that}$$

$$a(x, y) = \frac{x}{x^2 + y^2} \log a(x, y) = \frac{a(x, y)}{x^2 + y^2}$$

where

$$a(x, y) = \frac{x}{x^2 + y^2} \log \left(\frac{x}{x^2 + y^2} \right) = \frac{x}{x^2 + y^2} \log x - \frac{x}{x^2 + y^2} \log(x^2 + y^2)$$

$$\frac{\partial}{\partial x} \left(\frac{x}{x^2 + y^2} \log x \right) = \frac{1}{x^2 + y^2} \log x - \frac{x}{(x^2 + y^2)^2}$$

$$= \frac{1}{x^2 + y^2} \log x - \frac{x}{(x^2 + y^2)^2}$$

(5) n-solitons 1) $u(x,t)$ sol'n of KdV

2) continuous in x, t

3) $\lim_{x \rightarrow \pm\infty} u(x,t) = 0 \quad \forall t$

4) "rational-exponential" form

$$u(x,t) = \frac{\sum_{i=1}^m c_i e^{a_i x + b_i t}}{\sum_{j=1}^m g_j e^{\alpha_j x + \beta_j t}}$$

$$m, n \in \mathbb{N}$$

$$a_i, b_i, \alpha_j, \beta_j \in \mathbb{R}$$

5) for large $|t|$ graph of $u(x,t)$ has m local maxima

τ function

$$\tau(x,t) = \sum_{i=1}^m c_i e^{a_i x + b_i t} \quad \text{such that}$$

$$u(x,t) = 2 \partial_x^2 \log \tau(x,t) = \frac{2\tau \tau_{xx} - 2\tau_x^2}{\tau^2}$$

is solution

$$u(x,t) = 2 \partial_x^2 \log (e^{kx+k^3t} + e^{-kx-k^3t})$$

$$= 2 \partial_x^2 \frac{k e^{kx+k^3t} - k e^{-kx-k^3t}}{e^{kx+k^3t} + e^{-kx-k^3t}}$$

$$= \frac{8k^2}{(e^{kx+k^3t} + e^{-kx-k^3t})^2}$$

1. function f is called *linear*

⑤

linear (L) (L) (L) is not linear

non linear

$$M_y = \frac{1}{2} \left(1 + \frac{1}{2} \right) = 1.5$$

function

$$f(x) = \frac{1}{2} \left(1 + \frac{1}{2} \right) = 1.5$$

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function

linear

non linear

linear

linear



linear

linear



τ -functions for n -soliton solutions

⑥

choose $\{c_i\}$ $\{k_i\}$ $i=1, \dots, n$ real constants
 $n \times n$ matrix M

$$M_{ij} = \frac{2^{i+j}}{2x^{i+j}} \left(e^{k_j x + k_j^3 t} + c_j e^{-k_j x - k_j^3 t} \right)$$

Wronskian matrix

$$W(\underline{f}) = \begin{pmatrix} f_1 & \dots & f_n \\ f_1' & \dots & f_n' \\ \vdots & & \vdots \\ f_1^{(n-1)} & \dots & f_n^{(n-1)} \end{pmatrix}$$

$$W(\underline{f})$$

$$f_j = e^{k_j x + k_j^3 t} + c_j e^{-k_j x - k_j^3 t}$$

$\det W(\underline{f}) = 0$ iff some lin comb of f_i
 $\sum_i a_i f_i = 0$

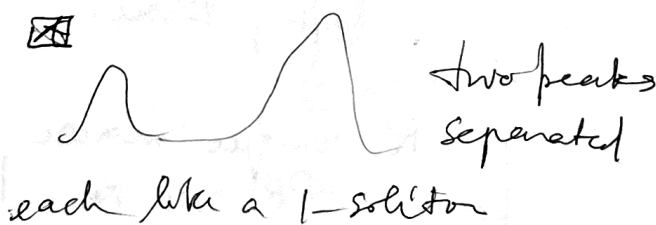
if $\tau(x,t) = \det M \neq 0$

then $u(x,t) = 2 \partial_x^2 \log \tau(x,t)$

Sol'n of KdV n -soliton sol'n

2-soliton sol'n ☒

for $|t|$ large $\rightsquigarrow$



surface $u(x,t)$ plane



scattering of one wave by other
 incoming/outgoing waves



⑦ Consider again Weyl algebra
of differential operators

now not nec. polyn coeff's
instead of

$$W_n(\mathbb{R}[x]) \text{ take } W_n(C^\infty(\mathbb{R})) \text{ on } \mathbb{R}^n$$

D diff op in one-var x linear

fix f in $\text{Ker}(D)$ iff

$$D = Q \circ \left(\partial - \frac{f'}{f} \right)$$

because $\left(\partial - \frac{f'}{f} \right) f = f' - \frac{f'}{f} \cdot f = 0$

so if $D = Q \circ \left(\partial - \frac{f'}{f} \right)$ then $f \in \text{Ker}(D)$

conversely: if n dim arbitrary

$V \subset \text{Ker}(D)$ k -dim subspace of $\text{Ker}(D)$

B basis for V

$\Rightarrow D = Q \circ R$ ← quater produced from $B = \{f_1, \dots, f_n\}$
diff op 1-variable $R(f) = \frac{W(f_1, \dots, f_n, f)}{W(f_1, \dots, f_n)}$

R unique monic nth order

ODE w/ Kernel = $V(B)$ span of B

$B' \cup B$ basis of $\text{Ker } D$
 $Q = (\text{leading coeff of } D) \cdot \text{unique monic op w/ Ker } B' \text{ under } B$
 $= \text{Image of } B' \text{ under } B$

then if $f \in \ker(D) \Rightarrow f$ also in $\ker(D - Q \circ R)$

②

$$(D - Q \circ R)(f) = D(f) - Q(R(f)) \\ = Q(\underbrace{R(f)}_{=0}) = 0$$

Strictly less order than D
(as Q has less order than D)

so $\ker$ should be smaller $\Rightarrow D - Q \circ R \equiv 0$

Isospectrality: matrices $n \times n$

C with $Cv = \lambda v$

A, M with $AC = MA$

$\Rightarrow w = Av$ satisfies $Mw = \lambda w$

so C & M have same spectrum

} $M = A(A^{-1})$
if A invertible
obvious
but true more generally

~~$M(t)$~~ $M(t) = A(t) C A(t)^{-1}$ continuous isosp. defn

if invertible, $M(t) = A(t) C A(t)^{-1}$

$$B = \dot{A}(t) A(t)^{-1}$$

$$\dot{M} = [B, M]$$

from $\dot{M} = \frac{d}{dt}(A C A^{-1}) \rightarrow$

for differential operators:

D_0 linear ordinary diff op

V fin dim invariant subspace $D_0 V \subset V$

⑨ K unique linear diff op

w/ kernel $= V$

$\exists$ D ordinary diff op

$$K \circ D_0 = D \circ K$$

$\forall \psi_0$ eigenvector $D_0 \psi_0 = \lambda_0 \psi_0$

$\psi_0 \notin V \Rightarrow \psi = K(\psi_0)$ eigenv of D w/ λ_0 eigenvalue

for $\phi \in V$ $K \circ \underbrace{D_0}_{V = \ker(K)}(\phi) = 0$

so $K \circ D_0 = 0 \circ K$ factors

$$\begin{aligned} D(\psi) &= D(K(\psi_0)) \\ &= K(D_0 \psi_0) = \lambda_0 K(\psi_0) \end{aligned}$$

example: $D_0 = \partial^2 + 2\partial + 1$

$V = \text{span}(x^2(1+x)e^x, (1+3x)e^x)$ inv. subsp.

$$\leadsto K = \partial^2 - \frac{6x^3 + 15x^2 + 8x + 1}{3x^3 + 3x^2 + x} \partial + \frac{3x^3 + 12x^2 + 16x + 4}{3x^3 + 3x^2 + x}$$

move ord diff op deg 2 into V as Kernel

$$= \partial^2 + q(x)\partial + c_0(x)$$

look for $D = \partial^2 + \alpha(x)\partial + \beta(x)$ s.t.

$$K \circ D_0 = D \circ K$$

compare terms of same order

$$\begin{aligned} \alpha(x)q(x) &= -2q(x) \\ 1 + \alpha(x) - 2q(x) &= \\ c_0(x) + \beta(x) - 2q(x) &+ 2q'(x) \end{aligned}$$

$$\begin{cases} \beta(x) = 1 - 2q'(x) \\ \alpha(x) = -2 \end{cases}$$

Example,

$$D_0 = \partial^2$$

$$D_0 \psi = 4\psi$$

$$\psi_0 = e^{2x}$$

$$V_{\text{span}} \{x + t^2, e^{x+t} + e^{-x}\}$$

$$\Rightarrow K = \partial^2 + \frac{(e^{t+2x} + 1)(t^2 + x)}{e^{t+2x}(t^2 + x - 1) + t^2 + x + 1}$$

$$D = \partial^2 - \frac{2(4t^4 e^{t+2x} + 8t^2 x e^{t+2x} + 4x^2 e^{t+2x})}{(t^2 e^{t+2x} - t^2 - e^{t+2x} + x e^{t+2x} + 1)}$$

$$\psi(x, t) = - \frac{e^{2x} (e^{t+2x} (2t^2 + 2x - 3) - 3(2t^2 + 2x + 1))}{-e^{t+2x}(t^2 + x - 1) + t^2 + x + 1}$$

$$\partial \psi = 4\psi$$

Lax form of KdV and other soliton equations

operator $H = \partial^2 + u(x, t)$

(Hamiltonian Schrödinger eq
(as stationary part) w/ potential $u(x, t)$ & time-dependent)

- if $u(x, t)$ is a ~~per~~ KdV n-soliton

↗ H ~~is~~ isospectral when t -variable changes

Gardner-Greene-Kruskal-Miura

- this observation combined w/ techniques for reconstruction of ~~spect~~ potential from spectral data $\rightarrow$ inverse scattering technique for solitons

a related approach: Lax pairs & Lax equation

①

$$D = \sum_v c_v(x,t) \partial_x^v$$

coeffs dep on t-param. differentiation only w/ x
ordinary diff op

$$\dot{D} = \sum \frac{\partial c_v(x,t)}{\partial t} \partial_x^v$$

if $D = \partial^2 + u(x,t) \partial$

$$M = \partial^3 + \frac{3}{2} u(x,t) \partial + \frac{3}{4} u_x(x,t)$$

then u solutions of KdV iff D, M satisfy

$$M \circ D = \partial^5 + \frac{5}{2} u \partial^3 + \frac{15}{4} u_x \partial^2 + \frac{3}{2} (u^2 + 2u_{xx}) \partial + \frac{9}{4} u u_x + u_{xxx}$$

$$D \circ M = \partial^5 + \frac{5}{2} u \partial^3 + \frac{15}{4} u_x \partial^2 + \frac{3}{2} (u^2 + 2u_{xx}) \partial + \frac{3}{4} (u u_x + u_{xxx})$$

$$\Rightarrow M \circ D - D \circ M = \frac{3}{2} u u_x + \frac{1}{4} u_{xxx}$$

if a PDE can be written in this form it has a Lax form w/ D, M its Lax pair

→ KdV has a Lax form

Use this to compute solutions of KdV !

$$\dot{D} = [M, D]$$

Lax equation

Just $u(x,t)$ as ∂^2 constant coeff w/ D

as ∂^2 constant coeff w/ D

KdV

example : $V = \text{span}(x, x^3 + 6t)$

(12)

$$D_0 = \partial^2 \quad \text{so } D_0 x = 0 \quad \text{so } D_0 V \subset V$$

$$D_0(x^3 + 6t) = 6x$$

$$K = \partial^2 - \frac{3x^3}{x^3 - 3t} \partial + \frac{3x}{x^3 - 3t}$$

monic second order ordinary K with $x, x^3 + 6t$ in kernel

then want $D = \partial^2 + u(x, t)$

s.t. $K \circ D_0 = D \circ K$ $\leftarrow$ find $u(x, t)$

computing 2nd order terms to match

$$\text{l.h.s.} = -\frac{3x}{x^3 - 3t} \quad \text{r.h.s.} = \frac{9x(x^3 - 3t)}{(x^3 - 3t)^2} + u(x, t)$$

$$\Rightarrow u(x, t) = -\frac{6(6xt + x^4)}{(x^3 - 3t)^2}$$