

① Sturm-Liouville

Differential operators (linear) single variable $x \in \mathbb{R}$

$$\mathcal{L} = a_n(x) \frac{\partial^n}{\partial x^n} + \dots + a_1(x) \frac{\partial}{\partial x} + a_0(x)$$

acts on a suff. smooth function $y(x)$
producing a new function $(\mathcal{L}y)(x)$

many physics-derived problems have second order
diff op's (as accel. & velocity)

$$\mathcal{L} = P(x) \frac{\partial^2}{\partial x^2} + R(x) \frac{\partial}{\partial x} - Q(x) \quad (*)$$

homog. eq. $\mathcal{L}y = 0$ two non-trivial lin indep solutions
 $y_1(x)$ $y_2(x)$
gen. sol'n $a y_1(x) + b y_2(x)$

Self adjoint differential operators

Sturm
Liouville
operators

$$\mathcal{L}y := \frac{d}{dx} \left(p(x) \frac{dy}{dx} \right) - q(x)y \quad (**)$$

y may be $\mathbb{C}$ -valued $\left\{ \begin{array}{l} p(x) \text{ real and at least once differentiable} \\ q(x) \text{ real and at least continuous} \end{array} \right.$

no more restrictive than $(*)$: if $p(x) \neq 0$ on domain

$$\frac{d^2}{dx^2} + \frac{R(x)}{P(x)} \frac{d}{dx} - \frac{Q(x)}{P(x)} = e^{-\int_0^x \frac{R(t)}{P(t)} dt} \frac{d}{dx} \left(e^{\int_0^x \frac{R(t)}{P(t)} dt} \frac{d}{dx} \right) - \frac{Q(x)}{P(x)}$$

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with $p(x) = e^{\int_0^x R(t)/p(t) dt}$ and $q(x) = \frac{C(x)p(x)}{p(x)}$

get (*) is of form (**)

so general case

Advantage of form (**): operator L is self adjoint

$$\langle f, g \rangle = \int_a^b \overline{f(x)} g(x) dx$$

if use right boundary conditions at a, b

$$\langle Lf, g \rangle = \int_a^b \left(\frac{d}{dx} (p(x) \frac{d\overline{f}}{dx}) - q(x) \overline{f(x)} \right) g(x) dx$$

$$= p(x) \frac{d\overline{f}}{dx} \Big|_a^b - \int_a^b p(x) \frac{d\overline{f}}{dx} \frac{dg}{dx} - q(x) \overline{f(x)} g(x) dx$$

$$= \left(p(x) \frac{d\overline{f}}{dx} - p(x) \overline{f} \frac{dg}{dx} \right) \Big|_a^b + \int_a^b \left(\frac{d}{dx} (p(x) \frac{dg}{dx}) - q(x) g(x) \right) \overline{f(x)} dx$$

$$= \left(p(x) \frac{d\overline{f}}{dx} - p(x) \overline{f} \frac{dg}{dx} \right) \Big|_a^b + \langle f, Lg \rangle$$

so boundary condition that makes L self-adj (symmetric)

is

$$(***) \quad p(x) \left(\frac{d\overline{f}}{dx} g(x) - \overline{f(x)} \frac{dg}{dx} \right) \Big|_a^b = 0$$

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examples of boundary conditions that satisfy this are

$$\begin{aligned} (b_1, b_2) &\neq (0, 0) \\ (c_1, c_2) &\neq (0, 0) \end{aligned}$$

ex 1:

$$b_1 f'(a) + b_2 f(a) = 0$$

$$c_1 f'(b) + c_2 f(b) = 0$$

for nontrivial b_i, c_i

(same b_i, c_i for all f 's in space of functions)
 $\Rightarrow$ both f & g

vanishes separately at a & b

ex 2:

if $p(a) = p(b)$

then periodic $f(a) = f(b)$

also works

$$f'(a) = f'(b)$$

$$f'(a) = f'(b)$$

} all funct's

ex 3:

$$p(a) = p(b) = 0$$

special case (singular)

Weight functions: instead of inner product

$$\langle f, g \rangle = \int_a^b \bar{f}(x) g(x) dx$$

Consider more general inner products with

a weight function $w(x) > 0$

more gen w only fin many zeros in $[a, b]$

$$\langle f, g \rangle = \int_a^b \bar{f}(x) g(x) w(x) dx$$

Examples:

with $w(x) \equiv 1$ and $[a, b] = [0, \pi]$ the functions

$f_n(x) = \sin(nx)$ are pairwise orthogonal
 $n \in \mathbb{N}$

[4] Example: $f_0(x) = 1$, $f_1(x) = 2x$, $f_2(x) = 4x^2 - 1$
 $f_3(x) = 8x^3 - 4x$ $f_4(x) = 16x^4 - 12x^2 + 1$
 $f_5(x) = 32x^5 - 32x^3 + 6x \dots$ orthogonal

with $w(x) = \sqrt{1-x^2}$ and $[a, b] = [-1, 1]$

(Chebyshev polynomials
of second kind)

Consider the eigenvalue problem for L
but written in the form (weighted)

$$L y(x) = \lambda w(x) y(x)$$

note: if $y(x)$ an eigenfunction of this problem
(I have $w(x) > 0$) then

$\tilde{y}(x) = \sqrt{w(x)} \cdot y(x)$ eigenfunction of
usual eigenvalue problem for

$$\tilde{L} \tilde{y} = \lambda \tilde{y} \quad \tilde{L} \tilde{y} = \frac{1}{\sqrt{w}} L \left(\frac{\tilde{y}}{\sqrt{w}} \right)$$

• if $Ly = \lambda w y$ then

$$\lambda \langle f, f \rangle_w = \langle f, \lambda w f \rangle = \langle f, Lf \rangle = \langle Lf, f \rangle$$

↑
weighted inn prod

$$= \bar{\lambda} \langle f, f \rangle_w \Rightarrow \lambda \in \mathbb{R}$$

• then if y eigenfunction $\bar{y}$ also (p.g. w real, λ real)

[5]

- eigen functions w/ different eigenvalues $\lambda_i \neq \lambda_j$ but same weight are orthogonal (w-in prod)

$$\lambda_j \langle y_i, y_j \rangle_w = \langle y_i, \lambda_j y_j \rangle = \langle \lambda_j y_i, y_j \rangle = \lambda_j \langle y_i, y_j \rangle_w$$

$$\lambda_j \neq \lambda_i \Rightarrow \langle y_i, y_j \rangle_w = 0$$

- orthonormal eigen functions

$$Y_i(x) = \frac{y_i(x)}{\sqrt{\int_a^b |y_i(x)|^2 w(x) dx}} = \|y_i\|_w$$

weighted norm

in particular interested in cases

where $\exists$ sequence of eigenvalues

$$\lambda_1, \lambda_2, \dots, \lambda_n, \dots \quad |\lambda_n| \rightarrow \infty$$

and set $\{Y_i\}_{i \geq 1}$ or orthonormal eigenfunctions

that give a basis of funct. space w/ a partic. choice of boundary cond. at a, b

i.e. can expand functions

$$f(x) = \sum_{n=1}^{\infty} a_n Y_n(x) \quad \text{with}$$

generalization of Fourier series

$$a_n = \int_a^b \overline{Y_n(x)} f(x) w(x) dx = \langle Y_n, f \rangle_w$$

[6] eigenvalue / eigenfunction problem
is problem of solution of "regular" Sturm Liouville problem

$$\frac{d}{dx} \left(p(x) \frac{dy}{dx} \right) + (\lambda w(x) - q(x)) y = 0 \quad \text{for } x \in [a, b]$$

$$\begin{cases} b_1 y(a) + b_2 y'(a) = 0 \\ c_1 y(b) + c_2 y'(b) = 0 \end{cases}$$

$$(b_1, b_2) \neq (0, 0)$$

$$(c_1, c_2) \neq (0, 0)$$

p, p', q, w contin
on $[a, b]$

$$p, w > 0 \quad \text{on } [a, b]$$

↑
"separated"
boundary conditions

Example of regular Sturm-Liouville problem:

$$\begin{aligned} y'' + \lambda y &= 0 & 0 < x < L \\ y(0) = y(L) &= 0 & (p(x) = w(x) = 1, \\ & & q(x) = 0) \end{aligned}$$

$$\text{then } \lambda_n = \frac{n^2 \pi^2}{L^2} \quad \text{and}$$

$$y_n(x) = \sin \frac{n\pi x}{L} \quad \text{eigenfunctions } n = 1, 2, 3, \dots$$

other example: start w/ equation

$$\frac{1}{2} H'' - x H' = -\lambda H \quad x \in \mathbb{R}$$

[7]

and look for solutions that grow
at most polynomially fast at $|x| \rightarrow \infty$

make the change of variable t to put
in Sturm-Liouville form

$$p(x) = - \int_0^x 2t dt = -x^2$$

then $\frac{d}{dt} \left(e^{-x^2} \frac{dH}{dt} \right) = -2\lambda e^{-x^2} H(x)$

is the equivalent Sturm-Liouville form

Hermite equation

self-adjointness condition

zero boundary
term because
 $e^{-x^2} |f(x)| \rightarrow 0$
at $x \rightarrow \pm\infty$

$$\int_{\mathbb{R}} \bar{f}(x) \frac{d}{dx} \left(e^{-x^2} \frac{dg}{dx} \right) dx = \int_{\mathbb{R}} \frac{d\bar{f}}{dx} e^{-x^2} \frac{dg}{dx} dx$$

and convergence
of integrals by
decay cond.

$$= \int_{\mathbb{R}} \frac{d}{dx} \left(e^{-x^2} \frac{d\bar{f}}{dx} \right) g(x) dx$$

$$H_0(x) = 1 \quad H_1(x) = 2x$$

$$H_2(x) = 4x^2 - 2 \quad \dots$$

$$H_n(x) = (-1)^n e^{x^2} \frac{d^n}{dx^n} e^{-x^2}$$

eigenfunctions w/ eigenvalues $n \geq 0$
non-neg integers

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orthogonal

$$\begin{aligned} \langle H_m, H_n \rangle_{e^{-x^2}} &= \int_{\mathbb{R}} H_m(x) H_n(x) e^{-x^2} dx \\ &= \delta_{m,n} \cdot 2^n \cdot \sqrt{\pi} \cdot n! \end{aligned}$$

Regular Sturm-Liouville problem

$$\begin{aligned} \frac{d}{dx} \left(p(x) \frac{dy}{dx} \right) - (q(x) - \lambda w(x)) y &= 0 \\ \begin{cases} b_1 y'(a) + b_2 y(a) = 0 \\ c_1 y'(b) + c_2 y(b) = 0 \end{cases} & \quad x \in [a, b] \end{aligned}$$

Thm:

$\Rightarrow$ eigenvalues form an increasing sequence

$$\lambda_0 < \lambda_1 < \lambda_2 < \lambda_3 < \dots$$

$$\lim_{n \rightarrow \infty} \lambda_n = \infty$$

unique eigenfunction y_n (up to normalization)
of λ_n and y_n has exactly $n-1$ zeros
for $a < x < b$ interval

(9) cases that are not regular

Singular Sturm-Liouville problems

Example :

$$x^2 y'' + x y' + (x^2 - m^2) y = 0 \quad 0 < x < a$$
$$y(a) = 0$$

not regular because $p(x) = w(x) = x$ $q(x) = -\frac{m^2}{x}$

- $p(x)$ & $w(x)$ not strictly positive:
zero at $x = 0$

- $q(x)$ not continuous at $x = 0$

- missing boundary condition
at $x = 0$

Notation J interval

$L(J, \mathbb{C}) :=$ vect sp of $\mathbb{C}$ -valued
Lebesgue measurable funct's
on J

$L_{\text{loc}}(J, \mathbb{C}) :=$ " on all compact subintervals
 $[\alpha, \beta] \subseteq J$

if J compact $L_{\text{loc}} = L$

10 Notation $A, B \in M_{2 \times 2}(\mathbb{C})$

$$A:B \in M_{2 \times 4}(\mathbb{C}) \quad \underbrace{\begin{pmatrix} A & B \end{pmatrix}}_{\text{columns}}$$

S Sturm Liouville problem (Regular)

SLP (*)

$$\begin{cases} -(py')' + qy = \lambda wy & J = (a, b) \\ \text{w/ conditions at boundary:} \\ A Y(a) + B Y(b) = 0 & Y = \begin{pmatrix} y \\ py' \end{pmatrix} \quad \text{"quasi derivative" of } y \\ A, B \in M_2(\mathbb{C}) \quad AEA^* = BEB^* & E = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \\ \text{rk}(A:B) = 2 \end{cases}$$

$$r = \frac{1}{p}, \quad q, \quad w \in L^1(J, \mathbb{R})$$

$w > 0$

(Note : r, q, w can vanish on subintervals
 $\leadsto$ these counted as "single zero")

"trivial solution" on $I \subset J$ $y|_I \equiv 0, \quad py'|_I \equiv 0$

$f \in L^1$ changes sign on I if $\exists A, B \subset I$ pos meas:
 real val. $f|_A > 0 \quad f|_B < 0$

II write SL eq also as $Y' = (P - \lambda W) Y$

$$Y = \begin{pmatrix} y \\ p y' \end{pmatrix} \quad P = \begin{pmatrix} 0 & 1/p \\ q & 0 \end{pmatrix} \quad W = \begin{pmatrix} 0 & 0 \\ w & 0 \end{pmatrix}$$

boundary conditions $A Y(a) + B Y(b) = 0$

cases:

①

separated

↑ simpler case like discussed before:

$$\begin{aligned} a_i, b_i &\in \mathbb{R} \\ (a_1, a_2) &\neq (0, 0) \\ (b_1, b_2) &\neq (0, 0) \end{aligned}$$

$$\begin{aligned} a_1 y(a) + a_2 p y'(a) &= 0 \\ b_1 y(b) + b_2 p y'(b) &= 0 \end{aligned}$$

Since ∂ -condition invariant under mult. by a ~~non~~ sing matrix

here can rescale and parameterize by

$$0 < \alpha < \pi$$

$$0 < \beta \leq \pi$$

$$\cos \alpha y(a) - \sin \alpha p y'(a) = 0$$

$$\cos \beta y(b) - \sin \beta p y'(b) = 0$$

↑ (for convenience in some computations)

② real coupled

$$Y(b) = K Y(a)$$

$$K \in SL_2(\mathbb{R})$$

③ complex coupled

$$Y(b) = e^{i\gamma} K Y(a)$$

$$\begin{aligned} K &\in SL_2(\mathbb{R}) \\ -\pi &< \gamma < 0 \quad \text{or} \quad 0 < \gamma < \pi \end{aligned}$$

Fact Any SLP $(*)$ is always equiv. to ①, ②, or ③

② all self-adjoint SLP w/ $w > 0$ weight

Thm: for SLP $(*)$ ($w/w > 0$ a.e. on J)
(no sign restriction on p, q)

(1) all eigenvalues are real, isolated
w/ no (finite) accumulation pts, countably ∞ set

(2) if p changes sign on J
 $\Rightarrow$ seq. of eigenvalues unbounded below
and above

$$\dots \leq \lambda_{-m} \leq \dots \leq \lambda_{-1} \leq \lambda_0 \leq \lambda_1 \leq \lambda_2 \leq \dots \leq \lambda_n \leq \dots$$

$$\lambda_n \rightarrow \infty, \quad \lambda_{-n} \rightarrow -\infty$$

each λ_k either simple or double (not more)

($\Rightarrow$ for each fixed $\lambda = \lambda_k$ exactly 2 lin. indep. solutions of eq.)

(3) if p changes sign on J
and ∂ -conditions are separated
then all λ_k are simple

(4) if $p > 0$ on J and ∂ -cond are real coupled
then eigenvalues are bounded below

$$-\infty < \lambda_0 \leq \lambda_1 \leq \lambda_2 \leq \dots \leq \lambda_n \leq \dots$$

$\lambda_n \rightarrow \infty$

[13]

w/ eigenv. simple or double

$\{u_n\}$ real valued eigenfunctions for λ_n
(for real coupled ∂ -cond.)

$\Rightarrow$ # zeros of u_n in open $\bar{J}$ is

- 0 or 1 for $n=0$
- $n-1, n$, or $n+1$ for $n \geq 1$

(5) if $p > 0$ and ∂ -cond is complex coupled
all eigenvalues are simple

if $\{u_n\}$ eigenfunctions for λ_n

number of zeros of $\operatorname{Re}(u_n)$ on $[a, b]$

- 0 or 1 for $n=0$
- $n-1, n, n+1$ for $n \geq 1$

Same for $\operatorname{Im}(u_n)$

but zeros not in same places
 u_n no zero on (a, b)

(6) if $p > 0$ and ∂ -conditions
are separated then eigenvalues are simple

$\{u_n\}$ eigenfunctions: u_n 's are unique
(up to normalization)

and u_n has exactly n zeros on $\bar{J}$

[14] (7) if $p > 0$ any 2-condition (1), (2), (3) have asymptotic formula:

$$\frac{\lambda_n}{n^2} \xrightarrow{n \rightarrow \infty} c = \pi^2 \left(\int_a^b \sqrt{\frac{w}{p}} \right)^{-2}$$

(8) if p changes sign $p^\pm = p$ where ≥ 0 and 0 else

$$\frac{\lambda_n}{n^2} \rightarrow \pi^2 \left(\int_a^b \sqrt{\frac{w}{p_+}} \right)^{-2} \quad n \rightarrow \infty$$

$$\frac{\lambda_n}{n^2} \rightarrow \pi^2 \left(\int_a^b \sqrt{\frac{w}{p_-}} \right)^{-2} \quad n \rightarrow -\infty$$

Main ideas in the proof of this theorem:

(see below) {

- operator L and spectrum (eigenvalue probl.) in a Hilbert space $L^2(J, w)$ Weighted Leb. meas $w(x)dx$

$\Rightarrow$ eigenvalues are real and there is an ∞ seq. of them

- eigenvalues isolated w/ no accumulation point in $\mathbb{R}$:
 characterization as zeros of an entire function

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$Y(t)$ matrix solution of SLP (*) (2×2)

$$\text{rk } Y(t) = \text{rk } Y(u) \quad \forall t, u \in J$$

for fixed $u \in J$:

primary fundamental matrix $\phi(t, u) = Y(t)Y(u)^{-1}$

$\phi(u, u) = I$ (also write as $\phi(t, u, \underline{P})$ or $\phi(t, u, \underline{\lambda})$)

$$\phi' = (P - \lambda W) \phi \quad \text{on } J \quad \phi(u, u, \lambda) = I$$

$a \leq u \leq b$

characteristic function:

$$\delta(\lambda) = \det(A + B \phi(a, b, \lambda))$$

can def for arbitrary $\lambda \in \mathbb{C}$

is an entire function of λ

(for fixed a, b, A, B, P)

Note:
 $\exists!$ also implies continuous dependence of solution on parameters so cont. in $a, b, \dots$

because ϕ is entire in $\lambda \in \mathbb{C}$

$\frac{d}{d\lambda}$: obtained from variation of parameters formula

$$\text{for } Y' = PY + F$$

for $Y' = PY$
 ϕ primary fund matrix
 C constant matrix then
 ϕC solution of $Y' = PY$

$$Y(t) = \phi(t, u, P)C + \int_u^t \phi(t, s, P) F(s) ds \quad \text{solution of } Y' = PY + F, Y(u) = C$$

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$$P \mapsto P + zW$$

$$\left(Y(t, u, P + (z+h)W) - Y(t, u, P + zW) \right)$$

$$= \left(\phi(t, u, P + (z+h)W) - \phi(t, u, P + zW) \right) \cdot C$$

$$+ \int_u^t \left(\phi(t, r, P + (z+h)W) - \phi(t, r, P + zW) \right) F(r) dr$$

$$S(z) = \phi^{-1}(\cdot, u, P + zW) W(\cdot) \phi(\cdot, u, P + zW)$$

$$\geq \phi(t, u, P + (z+h)W) \left(\phi(t, u, hS(z)) - \mathbb{I} \right)$$

$$= \left(h \int_u^t S(z) + o(h) \right)$$

$$= h \phi(t, u, P + (z+h)W) \int_u^t S(z) + o(h)$$

$$\Rightarrow = h \left(\phi(t, u, P + zW) \left(\int_u^t S(z) \right) \right) \cdot C$$

$$+ \int_u^t \phi(t, r, P + zW) \left(\int_r^t S(z) \right) F(r) dr$$

[17] So eigenvalues λ_k are zeros of

$$\delta(\lambda) = \det(A + B\phi(a, b, \lambda))$$

i.e. $\phi(t, u, \lambda)$ soln of problem SL w/ λ

and $\delta(\lambda) = 0 \Rightarrow$ boundary conditions are satisfied

zeros of an entire function $\Rightarrow$ discrete set (no accumulation pt.)

- The unboundedness below/above when p has $p_- \neq 0$ (resp. $p_+ \neq 0$) is a result of Møller (paper added to course website)

- for # of zeros of eigenfunctions
Rong, Wu, Zettl "Dependence of n -th Sturm-Liouville eigenvalue on the problem" 2019

- eigenvalues simple or double: at most two lin indep sol'n
(∂ -cond. can fix only one sol'n)

- simple eigenvalues for separated boundary condition $\Leftrightarrow$ uniqueness of solution w/ those boundary cond. for fixed λ (if existence)

- for complex coupled ∂ -cond. w/ L self adjoint

$$\text{have } \phi(a, b, \lambda_n) = e^{i\gamma} K$$

(∂ -condition)
again then fix unique sol'n

$\gamma = 0$ or $\gamma = \pi$

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Spectral thm for unbounded self adjoint operators on a Hilbert space

- spectral thm for bounded self adjoint op's

$$T \in B(H) \text{ lin op. } \|T\| = \sup \frac{\|Tx\|}{\|x\|} < \infty$$

$$\text{and } T = T^*$$

$$\sigma(T) = \underbrace{\text{Spec}(T)}_{\text{complement of reg set}}$$

discrete sp.

essential sp. {

residual spec if this fails

Cont. sp. if this fails (R_λ exists on dense dom. not bounded)

Reg. value if

i) R_λ(T) exists

ii) R_λ(T) def on dense dom.

iii) R_λ(T) bounded

P(T) set reg. values

if T ∈ B(H) ⇒ no residual spectrum

Spec(T) closed contained in disk ||T|| radius

∃ family of projections in H
(spectral projections) E_λ λ ∈ ℝ

$$E_\lambda E_\mu = E_\mu E_\lambda = E_\lambda \text{ for } \lambda < \mu$$

$$\lim_{\lambda \rightarrow -\infty} E_\lambda x = 0 \quad \lim_{\lambda \rightarrow +\infty} E_\lambda x = Ix = x$$

$$E_\lambda x = \lim_{\mu \rightarrow \lambda^+} E_\mu x$$

$$\left[\begin{array}{l} \text{on an interval} \\ \text{if } E_\lambda = 0 \quad \lambda < a \\ E_\lambda = I \quad \lambda > b \end{array} \right]$$

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$$u_{xy}(\lambda) := \langle E_\lambda x, y \rangle$$

Spectral thm for bounded self adjoint T : $\exists E_\lambda$ s.t. $\forall x, y \in \mathcal{H}$

$$\langle Tx, y \rangle = \int_a^b \lambda du_{xy}(\lambda)$$

$$\langle f(T)x, y \rangle = \int_a^b f(\lambda) du_{xy}(\lambda) \quad \left(\begin{array}{l} \text{poly.} \\ \text{more gen. (class} \\ \text{of cont. fun's)} \end{array} \right)$$

Spectral thm for unitary operators

$$U \in B(\mathcal{H}) \quad UU^* = U^*U = I$$

$$\sigma(U) \text{ closed subset of } S^1 = \{|\lambda| = 1\} \\ \lambda = e^{2\pi i \theta}$$

$\exists$ spectral family E_θ on $[-\pi, \pi]$ s.t. $\forall x, y \in \mathcal{H}$

$$\langle f(U)x, y \rangle = \int_{-\pi}^{\pi} f(e^{i\theta}) du_{xy}(\theta)$$

poly. funct.

$$u_{xy}(\theta) = \langle E_\theta x, y \rangle$$

by $U = e^{iS}$ for S bounded self-adjoint w/ $\sigma(S) \subset [-\pi, \pi]$

Spectral thm for unbounded self-adj. op's

Cayley transform

$$U = (T - iI)(T + iI)^{-1}$$

$T: D(T) \rightarrow \mathcal{H}$
dense domain
 $T^* = T$

$\Rightarrow U$ unitary operator

(Möbius transf. that maps real axis to unit circle)

(20) can invert $T = i(I+U)(I-U)^{-1}$ ($\Rightarrow 1$ not in $\sigma(U)$ as inverse exists)

$\exists$ spectral family E_θ (for $-U$)

$$\forall x \in \text{Dom}(T) \quad u_{xx}(\theta) = \langle E_\theta x, x \rangle$$

$$\langle Tx, x \rangle = \int_{-\pi}^{\pi} \tan\left(\frac{\theta}{2}\right) du_{xx}(\theta)$$

$$= \int_{\mathbb{R}} \lambda d\nu(\lambda) \quad \nu(\lambda) = \langle F_\lambda x, x \rangle$$

$$F_\lambda = E_{2\arctan(\lambda)}$$

$\mathcal{L}$ self-adjoint op. on $L^2(J, w)$

(unbounded)

$\leadsto$ spectral theorem

but from spectral theorem
can have discrete and continuous spectrum

cases w/ only discrete spectrum:

can ensure only discrete spectrum for T

if resolvent $R_\lambda(T)$ is a compact operator

[21]

• compact operators :

• $\forall B \in \mathcal{H}$ bounded set
 $\overline{T(B)}$ in $\mathcal{H}$ compact set
(e.g. image of unit ball)

• closure in $B(\mathcal{H})$ of
finite rank op's (matrices)

• $\{\psi_k\}$ bounded seq in $\mathcal{H}$
 $\{T\psi_k\}$ has
conv. subseq.

spectrum of compact operators :

- no essential spectrum
- $0 \in \text{Spec}(T)$ always (otherwise would get I comp. impossible on $\dim \mathcal{H} = \infty$)
- 0 not nec. eigenvalue
can be limit pt.
can have ∞ -multpl. if eigenvalue
- $\lambda \neq 0 \in \text{Spec}(T)$ isolated pt
finite multiplicity

$\text{Spec}(T) \setminus \{0\}$ either finite or sequence $\lambda_k \rightarrow 0$

• if resolvent $R_z(T)$ is compact it is so for all $z \in \rho(T)$

$$(T - zI)^{-1} = (T - z_0I)^{-1} - (z - z_0) \underbrace{(T - z_0I)^{-1}(T - zI)^{-1}}_{\text{comp op } \mathcal{K} \subset B(\mathcal{H}) \text{ ideal so compact if } (T - z_0I)^{-1} \text{ is}}$$

if $0 \notin \sigma(T)$ can just take $R_{z=0}(T) = T^{-1}$ compact

then $\sigma(T^{-1})$ is sequence w/ only accum. pt at 0

$\Rightarrow \{\lambda_k^{-1} \mid \lambda_k \in \sigma(T)\} \Rightarrow \sigma(T)$ discrete seq
no accum. pt.
except at ∞
unbounded
w/ $(\lambda_k) \rightarrow \infty$