

Tuesday May 4

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$H$  scaling hamiltonian in quantum system  
is infinitesimal generator of scaling action of  
 $\mathbb{R}_+^*$  on  $\mathcal{H}$

$N_E$  = spectral projection on  $[-E, E]$  in  $\mathbb{R}$  (dual to  $\mathbb{R}_+^*$   
under  $\lambda^{it}$  pairing)

$$N_E = \mathcal{D}_a(h_E) \quad \text{for} \quad h_E(u) = |u|^{-\frac{1}{2}} \frac{1}{2\pi} \int_{-E}^E |u|^{is} ds$$

using  $\lambda^{it} \mathcal{D}_a(\lambda)$  unitary action

from  $\chi_{[-E, E]}(t) = \int \lambda^{it} k(\lambda) d\lambda$  where the function  
 $k(\lambda)$  is characteristic function

$$k(\lambda) = \frac{1}{2\pi} \int_{-E}^E \lambda^{is} ds \quad \lambda \in \mathbb{R}_+^*$$

Then counting # quantum states of  $H \leq E$   
is done by computing dimension of int. of  $Q_\lambda$  &  $N_E$

where  $Q_\lambda$  = projection onto  $B_\lambda \subset \mathcal{H}$  span of prolate  
spheroidal wave functions  $2n \leq 4\lambda^2$

i.e. compute  $\text{Tr}(Q_\lambda N_E)$

$\sim \text{Tr}(R_\lambda \mathcal{D}_a(h_E))$

$R_\lambda = \hat{P}_\lambda P_\lambda$  error of order  $\Lambda^{-\alpha} \log \Lambda$  some  $\alpha > 0$

with respect to  $Q_\lambda$

Connes: if  $h \in \mathcal{F}(\mathbb{R}^*)$   $h(-u) = h(u)$  even, compact support  
Then  $\text{Tr}(R_\lambda \mathcal{D}_a(h)) = 2h(1) \log \Lambda + \int' \frac{h(u')}{|1-u|} du + o(1)$

$\int'$  is a principal value

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$$\lim_{\varepsilon \rightarrow 0} \left( \int_{|x| \geq \varepsilon} f(x) dx + f(0) \log \varepsilon \right) \quad \text{when } f \text{ loc. const. near zero}$$

More precisely:  $K$  local field ( $\mathbb{R}, \mathbb{C}, \mathbb{Q}_p$ , fin ext.)

& character  $F_\alpha = \text{Fourier transf. w/ char. } \alpha$

$\exists!$  distribution  $\rho_\alpha$  extending  $d\hat{u}$   $\int_K f(x) \alpha(xg) dx$

(normalized Haar measure of  $K^*$ ) to  $u=0$

with  $F_\alpha(\rho_\alpha)(1) = 0$

$$\int'_{(K, \alpha)} \frac{f(u^{-1})}{|1-u|} du = \langle \rho_\alpha, g \rangle \quad g(\lambda) = \frac{f((\lambda+1)^{-1})}{|\lambda+1|}$$

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But ...  $h_E$  is not of compact support  $\Leftrightarrow$

$$\text{cannot apply directly } \text{Tr}(R_\Lambda h_E(h)) = 2h(1) \log 1 + \int' \frac{h(u^{-1})}{|1-u|} du + o(1)$$

can still get same result unitary equiv. of

$N_E P_\Lambda P_\Lambda^\dagger$  with other operator

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Assume then can apply (3)

$$\text{Tr}(R_A \vartheta_{\alpha}(h)) = 2 h(1) \log \Lambda + \int' \frac{h(u^{-1})}{|1-u|} d^*u + o(1)$$

also to  $h=h_E$  even though it does not have compact support (use Connes' quantized calculus argument)

then obtain

$$2 h_E(1) \log \Lambda = \frac{1}{2\pi} 2E \cdot 2 \log \Lambda$$

while other term gives

$$\begin{aligned} & \int' \frac{h_E(u^{-1})}{|1-u|} d^*u \quad (d^*u = \frac{1}{2} \frac{du}{|u|}) \\ &= -2 (\langle N(E) \rangle - 1) \end{aligned}$$

to see this: have  $\langle N(E) \rangle = N(E) - N_{\text{osc}}(E)$

$$\text{with } N_{\text{osc}}(E) = \frac{1}{\pi} \text{Im}(\log \zeta(\frac{1}{2} + iE))$$

use an explicit formula for principal value, for even  $f(-u)=f(u)$   
with  $f(u)=f(u^{-1})$

$$\int' f(u) \frac{|u|^{1/2}}{|1-u|} d^*u = (\log \pi + \gamma) f(1) + \lim_{\varepsilon \rightarrow 0} \left( \int_0^1 f(u) \frac{2u^{1/2}}{(1-u^2)^{1-\varepsilon}} \frac{du}{u} - \frac{1}{\varepsilon} f'(1) \right)$$

$$\text{apply to } f(u) = |u|^{is} + |u|^{-is}$$

and compute  $\int_0^1 u^{is} \frac{2u^{1/2}}{(1-u^2)^{1-\varepsilon}} \frac{du}{u}$  using

$$\int_0^1 x^{(\frac{1}{4}+i\frac{\varepsilon}{2})} (1-x)^{-1+\varepsilon} \frac{dx}{x} = \frac{\Gamma(\frac{1}{4}+i\frac{\varepsilon}{2}) \Gamma(\varepsilon)}{\Gamma(\frac{1}{4}+i\frac{\varepsilon}{2}+\varepsilon)}$$

At  $\varepsilon=0$  residue = finite part =

$$-\frac{\Gamma'(\frac{1}{4}+i\frac{\varepsilon}{2})}{\Gamma(\frac{1}{4}+i\frac{\varepsilon}{2})} - \gamma$$

$$\Rightarrow \int' f(u) \frac{|u|^{1/2}}{|1-u|} d^*u = 2(\log \pi + \gamma) - 2\gamma - \frac{\Gamma'(\frac{1}{4}+i\frac{\varepsilon}{2})}{\Gamma(\frac{1}{4}+i\frac{\varepsilon}{2})} - \frac{\Gamma'(\frac{1}{4}-i\frac{\varepsilon}{2})}{\Gamma(\frac{1}{4}-i\frac{\varepsilon}{2})}$$

$$\Rightarrow \int' \frac{h_E(u^{-1})}{|1-u|} d^*u = \frac{1}{2\pi} \int_0^E (2\log \pi - \frac{\Gamma'(\frac{1}{4}+i\frac{s}{2})}{\Gamma(\frac{1}{4}+i\frac{s}{2})} - \frac{\Gamma'(\frac{1}{4}-i\frac{s}{2})}{\Gamma(\frac{1}{4}-i\frac{s}{2})}) ds$$

but in fact

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$$\langle N(E) \rangle = 1 + \frac{1}{2\pi} \int_0^E (-\log \pi + \operatorname{Re}(\frac{\Gamma'}{\Gamma}(\frac{1}{4} + i\frac{E}{2}))) ds$$

from

$$d \operatorname{Im} \log \Gamma(\frac{1}{4} + i\frac{E}{2}) = \frac{1}{2i} d (\log \Gamma(\frac{1}{4} + i\frac{E}{2}) - \log \Gamma(\frac{1}{4} - i\frac{E}{2})) \\ = \frac{1}{2} \operatorname{Re}(\frac{\Gamma'}{\Gamma}(\frac{1}{4} + i\frac{E}{2})) ds$$

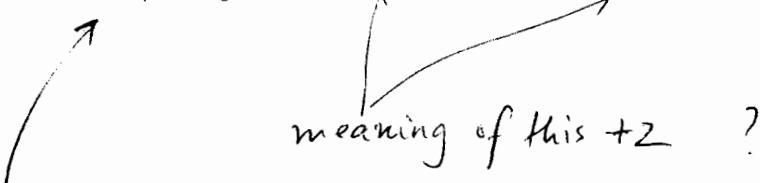
and from

$$\langle N(E) \rangle = 1 - \frac{E}{2\pi} \log \pi + \frac{1}{\pi} \operatorname{Im} \log \Gamma(\frac{1}{4} + i\frac{E}{2})$$

which follows from  $\langle N(E) \rangle = 1 + \frac{\theta(E)}{\pi}$

$$\text{with } \theta(E) = -\frac{E}{2} \log \pi + \operatorname{Im} \log \Gamma(\frac{1}{4} + i\frac{E}{2})$$

So obtain

$$\operatorname{Tr}(Q_N E) \sim \frac{4E}{2\pi} \log \Lambda - 2(\langle N(E) \rangle - 1) + o(1)$$


This is the symplectic volume of  $W(E, \Lambda)$  as in the classical Hamiltonian system

$$W(E, \Lambda) = \{(\lambda, s) \in \mathbb{R}_+^* \times \mathbb{R} : |\log \lambda| \leq \log \Lambda \text{ and } |s| \leq E\}$$

So want to match the cutoffs of original classical system to those implemented in quantum system

$$|q| \leq \Lambda \rightsquigarrow |\lambda| \leq \Lambda$$

$$|p| \leq \Lambda \rightsquigarrow |\lambda|^2 \leq \Lambda$$

$$|\hbar(p, q)| \leq E \rightsquigarrow |s| \leq E$$

Construct a map from even functions of  $u \in \mathbb{R}$  to functions of  $\lambda \in \mathbb{R}_+^*$  so that the cutoffs match

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$$E(f)(\lambda) ?$$

\* if map  $f(u) \mapsto \lambda^{1/2} f(\lambda)$  first cutoff ok

\* if map  $E$  also satisfies  $E(f)(\lambda) = E(\hat{f})(\lambda^\dagger)$  with  $\hat{f}$  Fourier transform  $\int e^{ixu} f(u) dx$   
then second cutoff also ok

but restriction  $f(u) \mapsto \lambda^{1/2} f(\lambda)$  does not satisfy this  
correct to give

$$E(f)(\lambda) = \lambda^{1/2} \sum_{n \in \mathbb{Z}} f(n\lambda) \quad \lambda \in \mathbb{R}_+^*$$

if impose conditions  $f(0) = \hat{f}(0) = 0$  these two condition  
⇒ extra 2 in  $\text{Tr}(Q_\lambda^{N_E})$

then Poisson summation formula gives second property

Note functions  $f$  even so just

$$\lambda^{1/2} \sum_{n=1}^{\infty} f(n\lambda) \quad (\text{up to factor 2})$$

then third cutoff also matches:

$$(D_a(\lambda) f)(x) = f(\lambda^\dagger x) \quad \text{gives}$$

$$E \circ D_a(\lambda) = |\lambda|^{1/2} D_m(\lambda) \circ E$$

action of  $\mathbb{R}_+^*$  on  
 $L^2(\mathbb{R})$

regular rep. of  $\mathbb{R}_+^*$  on  $L^2(\mathbb{R}^*)$

warning:  
 $E$  not an  
isometry

(6)

Refine this quantum system:

instead of just the scaling Hamiltonian  $h(q, p) = 2\pi p \cdot q$  .  
 view this as a component of a system that also has  
 a contribution at the archimedean place  
 for each non-archimedean place in the adeles & ideles,

$$A_Q = A_{Q,f} \times \mathbb{R} \quad A_Q^* = A_{Q,f}^* \times \mathbb{R}^*$$

finitely many degrees of freedom approximation

$$A_{Q,S} = \prod_{p \in S} Q_p \times \mathbb{R}$$

~~ideles~~

$S = \text{finite sets of primes } v \neq \infty$

$$C_{Q,S} = GL_1(A_{Q,S}) / Q_S^*$$

$$Q_S^* = \{ \gamma \in GL_1(Q) : |\gamma|_v = 1 \quad \forall v \notin S \}$$

$$Q_S = \{ q \in Q : |q|_v \leq 1 \quad \forall v \notin S \}$$

(rational numbers whose denominator  
only involves the primes  $p \in S$ )

$$Q_S^* = \{ \pm p_1^{n_1} \cdots p_k^{n_k} : p_j \in S \setminus \{\infty\}, n_j \in \mathbb{Z} \}$$

invertible elements  
of  $Q_S$   
inverse also has  
denom's in  $S$  only

Then have again scaling action

$$(\mathcal{D}_a(\lambda) \xi)(x) = \xi(\lambda^{-1}x) \quad \forall x \in A_{Q,S}$$

$$\forall \xi \in \mathcal{S}(A_{Q,S})$$

$$\lambda \in GL_1(A_{Q,S})$$

The map  $\mathcal{E} : \mathcal{P}(R) \rightarrow L^2(\mathbb{R}^*)$   
 extends to a map  $\mathcal{E} : \mathcal{P}(A_{Q,S}) \rightarrow L^2(C_{Q,S})$   
 with dense range given by

$$\mathcal{E}(f)(x) = |x|^{\frac{1}{2}} \Sigma(f)(x) \quad \text{where}$$

$$\Sigma(f)(x) = \sum_{q \in Q_S^*} f(qx)$$

Then one obtains

$$\text{Tr}(\mathcal{D}_a(h) R_\lambda) = 2 h(1) \log 1 + \sum_{v \in S} \int_{Q_v^*} \frac{h(u)}{|1-u|} d^* u + o(1)$$

$$\text{for } \mathcal{D}_a(h) = \int_{C_{Q,S}} h(g) \mathcal{D}_a(g) dg$$

trace formula computed on Hilbert space  $L^2(X_{Q,S})$

where  $X_{Q,S} = \frac{\text{adelic class space}}{\text{semi-local}} A_{Q,S}/Q_S^*$

and  $L^2(X_{Q,S}) = \text{completion of } \mathcal{P}(A_{Q,S}) \text{ in inner product}$

$$\langle f_1, f_2 \rangle_S := \langle \mathcal{E}(f_1), \mathcal{E}(f_2) \rangle$$

also obtained as convergent series

$$\langle f_1, f_2 \rangle_S = \sum_{\lambda \in Q_S^*} \langle f_1, \mathcal{D}_a(\lambda) f_2 \rangle \quad \begin{matrix} \text{inner prod in} \\ L^2(A_{Q,S}) \end{matrix}$$

Relation of the trace formula to the Weil explicit formula

(8)

Weil's distributional form of the explicit formulae

$\mathbb{K}$  = number field

$$\hat{f}(x, \rho) = \int_{C_K} f(u) \chi(u)/|u|^{\rho} du^*$$

$\chi$  character of  $C_K$  equal 1 on id. comp.

$h \in \mathcal{F}(C_K)$  w. compact support

$$\hat{h}(0) + \hat{h}(1) - \sum_{\chi} \sum_{\rho \in Z_{\tilde{\chi}}} \hat{h}(x, \rho) = \sum_v \int_{K_v} \frac{h(u^{-1})}{|1-u|} du$$

$Z_{\tilde{\chi}} = \text{set of zeros of } L(\tilde{\chi}, s) \text{ in strip } 0 < \operatorname{Re}(s) < 1$

(analogous for function fields)

So if  $\mathcal{H} = \operatorname{Coker}$  of map  $\mathcal{E}$

from  $L^2(X_{A_{K_0}}) \rightarrow L^2(C_{K_0})$  ~~closed~~

~~Weighted norms~~  $\cup$  Subspace w. conditions  
~~closed~~  $f(0) = \hat{f}(0) = 0$

$$\Rightarrow \sum_{\substack{\text{zeros} \\ \text{on the line} \\ \text{only}}} \hat{h}(\tilde{\chi}, \rho) = \operatorname{Tr}(\mathcal{D}_m(h)) \quad (\text{Connes '97})$$

↑  
Scaling, induced on  $\mathcal{H}$   
action

does not prove RH

Return to the endomotive formulation

(9)

find  $h \in \mathcal{P}(C_K)$  (Schwartz space w/ additional cond. "strong Schwartz space")

$$\mathcal{D}(h) = \int_{C_K} h(g) \mathcal{D}(g) dg \quad \text{acting on } H_0(C(D(A, \varphi)))$$

$$\text{Tr}(\mathcal{D}(h)) = \hat{h}(0) + \hat{h}(1) - \Delta \cdot \Delta h(1) - \sum_v \int'_{K_v^*} \frac{h(u^{-1})}{|1-u|} du$$

$$\Delta \cdot \Delta = -\log |D|$$

$\uparrow$  discriminant of the number field, "size of the ring of integers" (analog of  $\chi(C)$  for function field: self-inters of diag.)  $= \begin{cases} \text{Volume of fundam. domain} \\ \text{(measures how many primes are ramified)} \end{cases}$

Now: get counting of all zeros, not only those on critical line  
gets Weil expl. formula as trace formula

RH equivalent to

$$\text{Tr}(\mathcal{D}(h * h^\#)) \geq 0 \quad \forall h \in \mathcal{P}(C_\mathbb{Q})$$

where  $(h_1 * h_2)(g) = \int h_1(k) h_2(k^{-1}g) dk$   $\leftarrow$  multip. Haar meas

and  $h^\#(g) = |g|^\gamma \overline{h(g^{-1})}$

A positivity statement for the trace of a correspondence

$\Rightarrow$  this looks somewhat like Weil proof for function fields...

$$\mathbb{K} = \mathbb{F}_q(\mathbf{C}) \quad \Sigma_K = \text{places} = \text{pts of } \mathbf{C}(\bar{\mathbb{F}}_q) \\ \text{degree } n_v = \# \text{ orbit of Fr on } \mathbf{C}(\bar{\mathbb{F}}_q) \rightarrow \Sigma_K$$

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$$\sum_K(s) = \prod_{v \in \Sigma_K} (1 - q^{-n_v s})^{-1} = \frac{P(q^{-s})}{(1 - q^{-s})(1 - q^{1-s})}$$

$$P(T) = \prod (1 - \lambda_n T) \quad \begin{matrix} \text{eigenvalues of Frobenius acting on} \\ H^1_{\text{et}}(\bar{C}, \mathbb{Q}_\ell) \end{matrix}$$

$$RH \Leftrightarrow |\lambda_n| = q^{1/2}$$

Correspondences: divisors (lin. combinations of codim one subvarieties)  $Z \subset C \times C$   
degree, codegree, trace

$$d(Z) = Z \bullet (\{\text{pt}\} \times C)$$

$$d'(Z) = Z \bullet (C \times \{\text{pt}\}) \quad \text{diagonal} \swarrow$$

$$\text{Tr}(Z) = d(Z) + d'(Z) - Z \bullet \Delta$$

Weil:  $RH \Leftrightarrow \text{Tr}(Z * Z') > 0$   
applied to corresp: Frobenius action

$$Z_1 * Z_2 = (\pi_{13})_* (\pi_{12}^* Z_1 \circ \pi_{23}^* Z_2) \quad \text{and } Z' - \text{transpose of } Z$$

In NCG setting: correspondences  $\Rightarrow$  graph of scaling actions

$$g \in G_K \quad Z_g = \{(x, g^{-1}x)\} \subset A_{K/K^\star} \times A_{K/K^\star}$$

$$Z(f) = \int_{G_K} f(g) Z_g \, dg \quad f \in \mathcal{F}(C_K)$$

$$\text{degree: } d(Z(f)) = \hat{f}(1) = \int f(u)|u| \, du$$

$$\text{codegree: } d'(Z(f)) = \hat{f}(0) = \int f(u) \, du$$

(ii)

In Weil proof first adjust deg & codeg.  
by adding trivial correspondences  $\{pt\} \times C$  and  $C \times \{pt\}$

here can do same:

$$d(Z(f)) = f'(1) \text{ add } h \text{ with}$$

$$h(u, \lambda) = \sum_{n \in \mathbb{Z}, n \neq 0} \gamma(n\lambda) \quad \lambda \in R_+^*, u \in \mathbb{Z}^* \\ Q = \mathbb{Z}^* \times R_+^*$$

can find a function  $h$  as above  
which has  $h'(1) \neq 0$

because  $\int_R \sum_n \gamma(n\lambda) d\lambda \neq \sum_n \int_R \gamma(n\lambda) d\lambda = 0$   
(Fubini's theorem not nec. hold)

$$h \in V$$

$Z \cap P \times C$   
divisor  $Z(P)$  on  $C$

Next step in Weil's proof

Riemann-Roch's theorem applied to

$Z(P)$  of degree  $g \Rightarrow$  linearly equivalent (mod principal divisor)  
to an effective divisor (positivity)

$$d'(Z * Z') = d(Z)d(Z') = g d'(Z) = d(Z + \Delta)$$

$$\begin{aligned} \text{Tr}(Z * Z') &= 2g d'(Z) + (2g-2) d'(\Delta) - Y \cdot \Delta \\ &\geq (4g-2) d'(Z) - (4g-4) d'(Z) = 2 d'(Z) \geq 0 \end{aligned}$$

$$\text{where } Z * Z' = d'(Z) \Delta + Y$$

$Y$  divisor on  $C$        $1 Q_i \in C : Q_i(P) = Q_j(P) \quad (i \neq j) \}$   
where  $Z(P) = Q_1 + \dots + Q_g$   
effective divisor

$$\text{estimate } Y \cdot \Delta \leq (4g-4) d'(Z)$$

$$\Delta \cdot \Delta = 2-2g = \chi(C) \text{ Eul. char.}$$

Need in NCG: principal divisors & Riemann-Roch

