

# Spectral Action Models of Gravity and Packed Swiss Cheese Cosmologies

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Based on:

- Adam Ball, Matilde Marcolli, *Spectral Action Models of Gravity on Packed Swiss Cheese Cosmology*, *Classical Quantum Gravity* 33 (2016), no. 11, 115018, 39 pp.



## Homogeneity versus Isotropy in Cosmology

- Homogeneous and isotropic: Friedmann universe  $\mathbb{R} \times S^3$

$$\pm dt^2 + a(t)^2 (\sigma_1^2 + \sigma_2^2 + \sigma_3^2)$$

with round metric on  $S^3$  with  $SU(2)$ -invariant 1-forms  $\{\sigma_i\}$  satisfying relations

$$d\sigma_i = \sigma_j \wedge \sigma_k$$

for all cyclic permutations  $(i, j, k)$  of  $(1, 2, 3)$

- Homogeneous **but not** isotropic:

Bianchi IX mixmaster models  $\mathbb{R} \times S^3$

$$F(t) \left( \pm dt^2 + \frac{\sigma_1^2}{W_1^2(t)} + \frac{\sigma_2^2}{W_2^2(t)} + \frac{\sigma_3^2}{W_3^2(t)} \right)$$

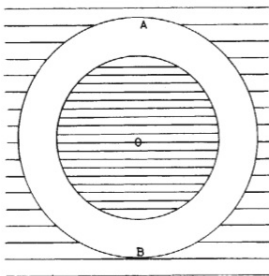
with a conformal factor  $F(t) \sim W_1(t)W_2(t)W_3(t)$

- Isotropic **but not** homogeneous?

$\Rightarrow$  Swiss Cheese Models

## Main Idea:

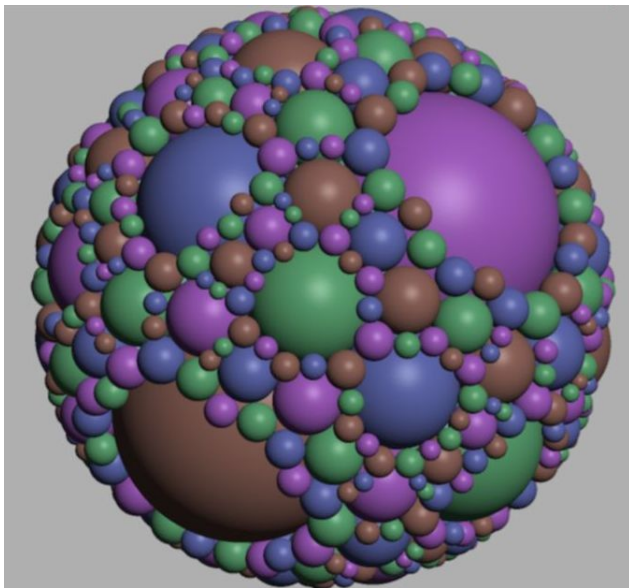
- M.J. Rees, D.W. Sciama, *Large-scale density inhomogeneities in the universe*, Nature, Vol.217 (1968) 511–516.



Cut off 4-balls from a FRW spacetime and replace with different density smaller region outside/inside patched across boundary with vanishing Weyl curvature tensor (isotropy preserved)

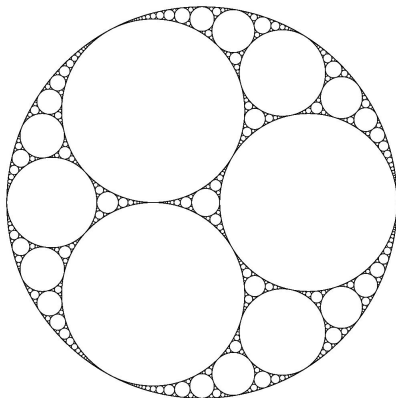
## Packed Swiss Cheese Cosmology

- Iterate construction removing more and more balls  $\Rightarrow$  Apollonian sphere packing of 3-dimensional spheres
- Residual set of sphere packing is fractal
- Proposed as explanation for possible fractal distribution of matter in galaxies, clusters, and superclusters
  - F. Sylos Labini, M. Montuori, L. Pietroneo, *Scale-invariance of galaxy clustering*, Phys. Rep. Vol. 293 (1998) N. 2-4, 61–226.
  - J.R. Mureika, C.C. Dyer, *Multifractal analysis of Packed Swiss Cheese Cosmologies*, General Relativity and Gravitation, Vol.36 (2004) N.1, 151–184.



## Apollonian sphere packings

- best known and understood case: Apollonian circle packing



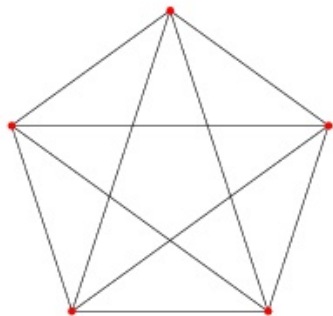
Configurations of mutually tangent circles in the plane, iterated on smaller scales filling a full volume region in the unit  $2D$  ball:  
residual set volume zero fractal of Hausdorff dimension  $1.30568\dots$



- Many results (geometric, arithmetic, analytic) known about Apollonian circle packings: see for example
  - R.L. Graham, J.C. Lagarias, C.L. Mallows, A.R. Wilks, C.H. Yan, *Apollonian circle packings: number theory*, J. Number Theory 100 (2003) 1–45
  - A. Kontorovich, H. Oh, *Apollonian circle packings and closed horospheres on hyperbolic 3-manifolds*, Journal of AMS, Vol 24 (2011) 603–648.
- **Higher dimensional** analogs of Apollonian packings: much more delicate and complicated geometry
  - R.L. Graham, J.C. Lagarias, C.L. Mallows, A.R. Wilks, C.H. Yan, *Apollonian Circle Packings: Geometry and Group Theory III. Higher Dimensions*, Discrete Comput. Geom. 35 (2006) 37–72.

## Some known facts on Apollonian sphere packings

- **Descartes configuration** in  $D$  dimensions:  $D + 2$  mutually tangent  $(D - 1)$ -dimensional spheres
- Example: start with  $D + 1$  equal size mutually tangent  $S^{D-1}$  centered at the vertices of  $D$ -simplex and one more smaller sphere in the center tangent to all



4-dimensional simplex

- **Quadratic Soddy–Gosset relation** between radii  $a_k$

$$\left( \sum_{k=1}^{D+2} \frac{1}{a_k} \right)^2 = D \sum_{k=1}^{D+2} \left( \frac{1}{a_k} \right)^2$$

- **curvature-center coordinates:**  $(D + 2)$ -vector

$$w = \left( \frac{\|x\|^2 - a^2}{a}, \frac{1}{a}, \frac{1}{a}x_1, \dots, \frac{1}{a}x_D \right)$$

(first coordinate curvature after inversion in the unit sphere)

- **Configuration space**  $\mathcal{M}_D$  of all Descartes configuration in  $D$  dimensions = all solutions  $\mathcal{W}$  to equation

$$\mathcal{W}^t Q_D \mathcal{W} = \begin{pmatrix} 0 & -4 & 0 \\ -4 & 0 & 0 \\ 0 & 0 & 2I_D \end{pmatrix}$$

with left and a right action of Lorentz group  $O(D+1, 1)$

- **Dual Apollonian group**  $\mathcal{G}_D^\perp$  generated by reflections: inversion with respect to the  $j$ -th sphere

$$S_j^\perp = I_{D+2} + 2 \mathbf{1}_{D+2} e_j^t - 4 e_j e_j^t$$

$e_j = j$ -th unit coordinate vector

- $D \neq 3$ : only relations in  $\mathcal{G}_D^\perp$  are  $(S_j^\perp)^2 = 1$
- $\mathcal{G}_D^\perp$  discrete subgroup of  $GL(D+2, \mathbb{R})$
- Apollonian packing  $\mathcal{P}_D =$  an orbit of  $\mathcal{G}_D^\perp$  on  $\mathcal{M}_D$

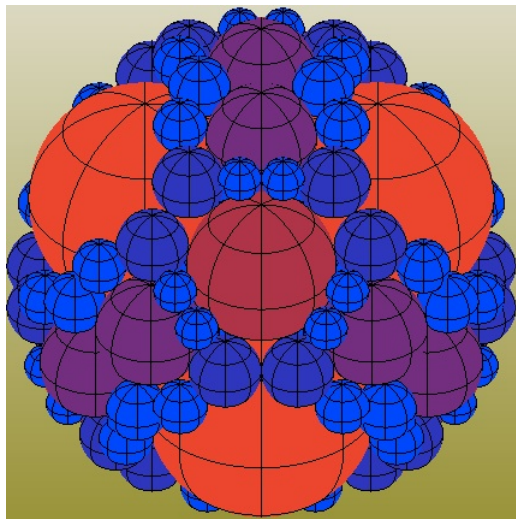
$\Rightarrow$  **iterative construction**: at  $n$ -th step add spheres obtained from initial Descartes configuration via all possible

$$S_{j_1}^\perp S_{j_2}^\perp \cdots S_{j_n}^\perp, \quad j_k \neq j_{k+1}, \quad \forall k$$

there are  $N_n$  spheres in the  $n$ -th level

$$N_n = (D+2)(D+1)^{n-1}$$

## iterative construction of sphere packings



- **Length spectrum**: radii of spheres in packing  $\mathcal{P}_D$

$$\mathcal{L} = \mathcal{L}(\mathcal{P}_D) = \{a_{n,k} : n \in \mathbb{N}, 1 \leq k \leq (D+2)(D+1)^{n-1}\}$$

radii of spheres  $S_{a_{n,k}}^{D-1}$

- **Melzak's packing constant**  $\sigma_D(\mathcal{P}_D)$  exponent of convergence of series

$$\zeta_{\mathcal{L}}(s) = \sum_{n=1}^{\infty} \sum_{k=1}^{(D+2)(D+1)^{n-1}} a_{n,k}^s$$

- **Residual set**:  $\mathcal{R}(\mathcal{P}_D) = B^D \setminus \cup_{n,k} B_{a_{n,k}}^D$  with

$$\partial B_{a_{n,k}}^D = S_{a_{n,k}}^{D-1} \in \mathcal{P}_D$$

- Packing  $\Rightarrow \text{Vol}_D(\mathcal{R}(\mathcal{P}_D)) = 0 \Rightarrow \sum_{\mathcal{L}} a_{n,k}^D < \infty \Rightarrow \sigma_D(\mathcal{P}_D) \leq D$
- **packing constant and Hausdorff dimension**:

$$\dim_H(\mathcal{R}(\mathcal{P}_D)) \leq \sigma_D(\mathcal{P}_D)$$

for Apollonian circles known to be same

- **Sphere counting function**: spheres with given curvature bound

$$\mathcal{N}_\alpha(\mathcal{P}_D) = \#\{S_{a_{n,k}}^{D-1} \in \mathcal{P}_D : a_{n,k} \geq \alpha\}$$

curvatures  $c_{n,k} = a_{n,k}^{-1} \leq \alpha^{-1}$

- for Apollonian circles power law (Kontorovich–Oh)

$$\mathcal{N}_\alpha(\mathcal{P}_2) \sim_{\alpha \rightarrow 0} \alpha^{-\dim_H(\mathcal{R}(\mathcal{P}_2))}$$

- for higher dimensions (Boyd): packing constant

$$\limsup_{\alpha \rightarrow 0} -\frac{\log \mathcal{N}_\alpha(\mathcal{P}_D)}{\log \alpha} = \sigma_D(\mathcal{P}_D)$$

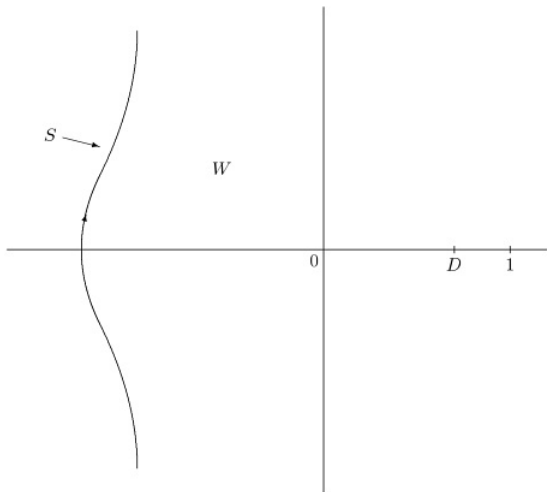
if limit exists  $\mathcal{N}_\alpha(\mathcal{P}_D) \sim_{\alpha \rightarrow 0} \alpha^{-(\sigma_D(\mathcal{P}_D)+o(1))}$

## Screens and Windows

- in general  $\zeta_{\mathcal{L}_D}(s)$  need have analytic continuation to meromorphic on whole  $\mathbb{C}$
  - $\exists$  *screen*  $\mathcal{S}$ : curve  $S(t) + it$  with  $S : \mathbb{R} \rightarrow (-\infty, \sigma_D(\mathcal{P}_D)]$
  - *window*  $\mathcal{W}$  = region to the right of screen  $\mathcal{S}$  where analytic continuation
- 
- M.L. Lapidus, M. van Frankenhuijsen, *Fractal geometry, complex dimensions and zeta functions. Geometry and spectra of fractal strings*, Second edition. Springer Monographs in Mathematics. Springer, 2013.



## Screens and windows



## Some additional assumptions

- **Definition:**

Apollonian packing  $\mathcal{P}_D$  of  $(D - 1)$ -spheres is *analytic* if

- 1  $\zeta_{\mathcal{L}}(s)$  has analytic to meromorphic function on a region  $\mathcal{W}$  containing  $\mathbb{R}_+$
  - 2  $\zeta_{\mathcal{L}}(s)$  has only one pole on  $\mathbb{R}_+$  at  $s = \sigma_D(\mathcal{P}_D)$ .
  - 3 pole at  $s = \sigma_D(\mathcal{P}_D)$  is simple
- **Also assume:**  $\exists \lim_{\alpha \rightarrow 0} -\frac{\log \mathcal{N}_{\alpha}(\mathcal{P}_D)}{\log \alpha} = \sigma_D(\mathcal{P}_D)$
  - **Question:** in general when are these satisfied for packings  $\mathcal{P}_D$ ?
  - focus on  $D = 4$  cases with these conditions

## Rough estimate of the packing constant

- $\mathcal{P} = \mathcal{P}_4$  Apollonian packing of 3-spheres  $S_{a_n,k}^3$
- at level  $n$ : average curvature

$$\frac{\gamma_n}{N_n} = \frac{1}{6 \cdot 5^{n-1}} \sum_{k=1}^{6 \cdot 5^{n-1}} \frac{1}{a_{n,k}}$$

- estimate  $\sigma_4(\mathcal{P}_4)$  with averaged version:  $\sum_n N_n \left(\frac{\gamma_n}{N_n}\right)^{-s}$

$$\sigma_{4,av}(\mathcal{P}) = \lim_{n \rightarrow \infty} \frac{\log(6 \cdot 5^{n-1})}{\log\left(\frac{\gamma_n}{6 \cdot 5^{n-1}}\right)}$$

- generating function of the  $\gamma_n$  known (Mallows)

$$G_{D=4} = \sum_{n=1}^{\infty} \gamma_n x^n = \frac{(1-x)(1-4x)u}{1 - \frac{22}{3}x - 5x^2}$$

$u$  = sum of the curvatures of initial Descartes configuration

- obtain explicitly ( $u = 1$  case)

$$\gamma_n = \frac{(11 + \sqrt{166})^n(-64 + 9\sqrt{166}) + (11 - \sqrt{166})^n(64 + 9\sqrt{166})}{3^n \cdot 10 \cdot \sqrt{166}}$$

- this gives a value

$$\sigma_{4,av}(\mathcal{P}) = 3.85193\dots$$

- in Apollonian circle case where  $\sigma(\mathcal{P})$  known this method gives larger value, so expect  $\sigma_4(\mathcal{P}) < \sigma_{4,av}(\mathcal{P})$
- constraints on the packing constant:

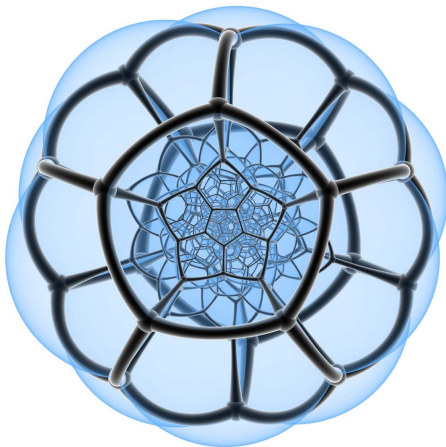
$$3 < \dim_H(\mathcal{R}(\mathcal{P})) \leq \sigma_4(\mathcal{P}) < \sigma_{4,av}(\mathcal{P}) = 3.85193\dots$$

## Models of (Euclidean, compactified) spacetimes

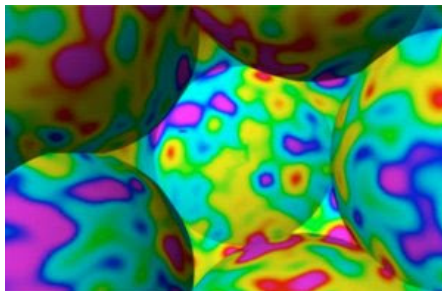
- 1 Homogeneous Isotropic cases:  $S^1_\beta \times S^3_a$
- 2 Cosmic Topology cases:  $S^1_\beta \times Y$  with  $Y$  a spherical space form  $S^3/\Gamma$  or a flat Bieberbach manifold  $T^3/\Gamma$  (modulo finite groups of isometries)
- 3 Packed Swiss Cheese:  $S^1_\beta \times \mathcal{P}$  with Apollonian packing of 3-spheres  $S^3_{a_{n,k}}$
- 4 Fractal arrangements with cosmic topology

## Fractal arrangements with cosmic topology

- Example: Poincaré homology sphere, dodecahedral space  $S^3/\mathcal{I}_{120}$ , fundamental domain dodecahedron

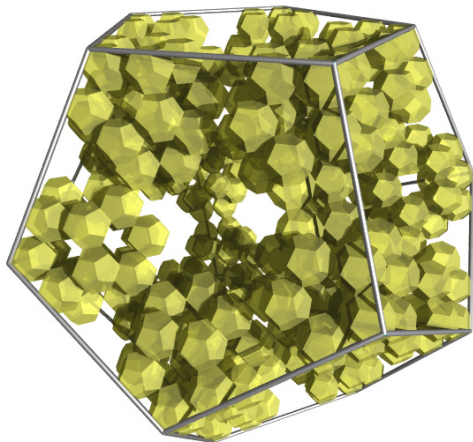


- considered a likely candidate for cosmic topology
  - S. Caillerie, M. Lachièze-Rey, J.P. Luminet, R. Lehoucq, A. Riazuelo, J. Weeks, *A new analysis of the Poincaré dodecahedral space model*, *Astron. and Astrophys.* 476 (2007) N.2, 691–696



- build a fractal model based on dodecahedral space

## Fractal configurations of dodecahedra (Sierpinski dodecahedra)





- spherical dodecahedron has  $\text{Vol}(Y) = \text{Vol}(S_a^3/\mathcal{I}_{120}) = \frac{\pi^2}{60} a^3$
- simpler than sphere packings because uniform scaling at each step:  $20^n$  new dodecahedra, each scaled by a factor of  $(2 + \phi)^{-n}$

$$\dim_H(\mathcal{P}_{\mathcal{I}_{120}}) = \frac{\log(20)}{\log(2 + \phi)} = 2.32958\dots$$

- close up all dodecahedra in the fractal identifying edges with  $\mathcal{I}_{120}$ : get fractal arrangement of Poincaré spheres  $Y_{a(2+\phi)^{-n}}$
- zeta function has analytic continuation to all  $\mathbb{C}$

$$\zeta_{\mathcal{L}}(s) = \sum_n 20^n (2 + \phi)^{-ns} = \frac{1}{1 - 20(2 + \phi)^{-s}}$$

exponent of convergence  $\sigma = \dim_H(\mathcal{P}_{\mathcal{I}_{120}}) = \frac{\log(20)}{\log(2+\phi)}$  and poles

$$\sigma + \frac{2\pi im}{\log(2 + \phi)}, \quad m \in \mathbb{Z}$$

## Spectral action models of gravity (modified gravity)

- **Spectral triple:**  $(\mathcal{A}, \mathcal{H}, D)$ 
  - ① unital associative algebra  $\mathcal{A}$
  - ② represented as bounded operators on a Hilbert space  $\mathcal{H}$
  - ③ Dirac operator: self-adjoint  $D^* = D$  with compact resolvent, with bounded commutators  $[D, a]$
- prototype:  $(C^\infty(M), L^2(M, S), \not{D}_M)$
- extends to non smooth objects (fractals) and noncommutative (NC tori, quantum groups, NC deformations, etc.)

## Action functional

- Suppose *finitely summable*  $ST = (\mathcal{A}, \mathcal{H}, D)$

$$\zeta_D(s) = \text{Tr}(|D|^{-s}) < \infty, \quad \Re(s) \gg 0$$

- **Spectral action** (Chamseddine–Connes)

$$\mathcal{S}_{ST}(\Lambda) = \text{Tr}(f(D/\Lambda)) = \sum_{\lambda \in \text{Spec}(D)} \text{Mult}(\lambda) f(\lambda/\Lambda)$$

$f$  = smooth approximation to (even) cutoff

**Asymptotic expansion** (Chamseddine–Connes) for  
(almost) commutative geometries:

$$\mathrm{Tr}(f(D/\Lambda)) \sim \sum_{\beta \in \Sigma_{ST}^+} f_\beta \Lambda^\beta \int |D|^{-\beta} + f(0) \zeta_D(0)$$

- Residues

$$\int |D|^{-\beta} = \frac{1}{2} \mathrm{Res}_{s=\beta} \zeta_D(s)$$

- Momenta  $f_\beta = \int_0^\infty f(v) v^{\beta-1} dv$
- **Dimension Spectrum**  $\Sigma_{ST}$  poles of zeta functions  
 $\zeta_{a,D}(s) = \mathrm{Tr}(a|D|^{-s})$
- positive dimension spectrum  $\Sigma_{ST}^+ = \Sigma_{ST} \cap \mathbb{R}_+^*$

**Warning:** for fractal spaces also oscillatory terms coming from part  
of  $\Sigma_{ST}$  off the real line

## Zeta function and heat kernel (manifolds)

- Mellin transform

$$|D|^{-s} = \frac{1}{\Gamma(s/2)} \int_0^\infty e^{-tD^2} t^{\frac{s}{2}-1} dt$$

- heat kernel expansion

$$\mathrm{Tr}(e^{-tD^2}) = \sum_{\alpha} t^{\alpha} c_{\alpha} \quad \text{for } t \rightarrow 0$$

- zeta function expansion

$$\zeta_D(s) = \mathrm{Tr}(|D|^{-s}) = \sum_{\alpha} \frac{c_{\alpha}}{\Gamma(s/2)(\alpha + s/2)} + \text{holomorphic}$$

- taking residues

$$\mathrm{Res}_{s=-2\alpha} \zeta_D(s) = \frac{2c_{\alpha}}{\Gamma(-\alpha)}$$

**Example** spectral action of the round 3-sphere  $S^3$

$$\mathcal{S}_{S^3}(\Lambda) = \text{Tr}(f(D_{S^3}/\Lambda)) = \sum_{n \in \mathbb{Z}} n(n+1) f\left(\left(n + \frac{1}{2}\right)/\Lambda\right)$$

- zeta function

$$\zeta_{D_{S^3}}(s) = 2\zeta\left(s - 2, \frac{3}{2}\right) - \frac{1}{2}\zeta\left(s, \frac{3}{2}\right)$$

$\zeta(s, q)$  = Hurwitz zeta function

- by asymptotic expansion

$$\mathcal{S}_{S^3}(\Lambda) \sim \Lambda^3 f_3 - \frac{1}{4}\Lambda f_1$$

- can also compute using Poisson summation formula (Chamseddine–Connes): estimate error term  $O(\Lambda^{-\infty})$

**Example:** round 3-sphere  $S_a^3$  radius  $a$

$$\zeta_{D_{S_a^3}}(s) = a^s \left( 2\zeta\left(s-2, \frac{3}{2}\right) - \frac{1}{2}\zeta\left(s, \frac{3}{2}\right) \right)$$

$$\mathcal{S}_{S_a^3}(\Lambda) \sim (\Lambda a)^3 f_3 - \frac{1}{4}(\Lambda a) f_1$$

**Example:** spherical space form  $Y = S_a^3/\Gamma$  (Ćačić, Marcolli, Teh)

$$\mathcal{S}_Y(\Lambda) \sim \frac{1}{\#\Gamma} \mathcal{S}_{S_a^3}(\Lambda)$$

## Why a model of (Euclidean) Gravity?

- $M$  compact Riemannian 4-manifold

$$\text{Tr}(f(D/\Lambda)) \sim 2\Lambda^4 f_4 a_0 + 2\Lambda^2 f_2 a_2 + f_0 a_4$$

coefficients  $a_0$ ,  $a_2$  and  $a_4$ :

- cosmological term

$$f_4 \Lambda^4 \int |D|^{-4} = \frac{48 f_4 \Lambda^4}{\pi^2} \int \sqrt{g} d^4 x$$

- Einstein–Hilbert term

$$f_2 \Lambda^2 \int |D|^{-2} = \frac{96 f_2 \Lambda^2}{24 \pi^2} \int R \sqrt{g} d^4 x$$

- modified gravity terms (Weyl curvature and Gauss–Bonnet)

$$f(0) \zeta_D(0) = \frac{f_0}{10 \pi^2} \int \left( \frac{11}{6} R^* R^* - 3 C_{\mu\nu\rho\sigma} C^{\mu\nu\rho\sigma} \right) \sqrt{g} d^4 x$$

$$C^{\mu\nu\rho\sigma} = \text{Weyl curvature and } R^* R^* = \frac{1}{4} \epsilon^{\mu\nu\rho\sigma} \epsilon_{\alpha\beta\gamma\delta} R_{\mu\nu}^{\alpha\beta} R_{\rho\sigma}^{\gamma\delta}$$

momenta: (effective) gravitational and cosmological constant



## Spectral action on a fractal spacetime:

- $S_{\beta}^1 \times \mathcal{P}$ : Apollonian packing
  - $S_{\beta}^1 \times \mathcal{P}_{\Upsilon}$ : fractal dodecahedral space
- 1 Construct a spectral triple for the geometries  $\mathcal{P}$  and  $\mathcal{P}_{\Upsilon}$
  - 2 Compute the zeta function
  - 3 Compute the asymptotic form of the spectral action
  - 4 Effect of product with  $S_{\beta}^1$

⇒ look for **new terms** in the spectral action (in addition to usual gravitational terms) that detect **presence of fractality**

## The spectral triple of a fractal geometry

- case of Sierpinski gasket: Christensen, Ivan, Lapidus
- similar case for  $\mathcal{P}$  and  $\mathcal{P}_Y$
- for  $D$ -dim packing

$$\mathcal{P}_D = \{S_{a_{n,k}}^{D-1} : n \in \mathbb{N}, 1 \leq k \leq (D+2)(D+1)^{n-1}\}$$

$$(\mathcal{A}_{\mathcal{P}_D}, \mathcal{H}_{\mathcal{P}_D}, \mathcal{D}_{\mathcal{P}_D}) = \bigoplus_{n,k} (\mathcal{A}_{\mathcal{P}_D}, \mathcal{H}_{S_{a_{n,k}}^{D-1}}, \mathcal{D}_{S_{a_{n,k}}^{D-1}})$$

- for  $\mathcal{P}_Y$  with  $Y_a = S^3/\mathcal{I}_{120}$ :

$$(\mathcal{A}_{\mathcal{P}_Y}, \mathcal{H}_{\mathcal{P}_Y}, \mathcal{D}_{\mathcal{P}_Y}) = (\mathcal{A}_{\mathcal{P}_Y}, \bigoplus_n \mathcal{H}_{Y_{a_n}}, \bigoplus_n \mathcal{D}_{Y_{a_n}})$$

with  $a_n = a(2 + \phi)^{-n}$

## Zeta functions for Apollonian packing of 3-spheres:

- Lengths zeta function (fractal string)

$$\zeta_{\mathcal{L}}(s) := \sum_{n \in \mathbb{N}} \sum_{k=1}^{6 \cdot 5^{n-1}} a_{n,k}^s$$

with  $\mathcal{L} = \mathcal{L}_4 = \{a_{n,k} \mid n \in \mathbb{N}, k \in \{1, \dots, 6 \cdot 5^{n-1}\}\}$

- zeta function of Dirac operator of the spectral triple

$$\mathrm{Tr}(|\mathcal{D}_{\mathcal{P}}|^{-s}) = \sum_{n=1}^{\infty} \sum_{k=1}^{6 \cdot 5^{n-1}} \mathrm{Tr}(|D_{S_{a_{n,k}}^3}|^{-s})$$

each term  $\mathrm{Tr}(|D_{S_{a_{n,k}}^3}|^{-s}) = a_{n,k}^s (2\zeta(s-2, \frac{3}{2}) - \frac{1}{2}\zeta(s, \frac{3}{2}))$  gives

$$\begin{aligned} \mathrm{Tr}(|\mathcal{D}_{\mathcal{P}}|^{-s}) &= \left( 2\zeta\left(s-2, \frac{3}{2}\right) - \frac{1}{2}\zeta\left(s, \frac{3}{2}\right) \right) \sum_{n,k} a_{n,k}^s \\ &= \left( 2\zeta\left(s-2, \frac{3}{2}\right) - \frac{1}{2}\zeta\left(s, \frac{3}{2}\right) \right) \zeta_{\mathcal{L}}(s) \end{aligned}$$

**Spectral action** for Apollonian packing of 3-spheres:  
(under good conditions on  $\zeta_{\mathcal{L}}(s)$ )

- Positive Dimension Spectrum:  $\Sigma_{STPSC}^+ = \{1, 3, \sigma_4(\mathcal{P})\}$
- asymptotic spectral action

$$\begin{aligned} \text{Tr}(f(\mathcal{D}_{\mathcal{P}}/\Lambda)) &\sim \Lambda^3 \zeta_{\mathcal{L}}(3) f_3 - \Lambda \frac{1}{4} \zeta_{\mathcal{L}}(1) f_1 \\ &+ \Lambda^\sigma \left( \zeta\left(\sigma - 2, \frac{3}{2}\right) - \frac{1}{4} \zeta\left(\sigma, \frac{3}{2}\right) \right) \mathcal{R}_\sigma f_\sigma + \mathcal{S}_\Lambda^{\text{osc}} \end{aligned}$$

$\sigma = \sigma_4(\mathcal{P})$  packing constant; residue  $\mathcal{R}_\sigma = \text{Res}_{s=\sigma} \zeta_{\mathcal{L}}(s)$ , and momenta  $f_\beta = \int_0^\infty v^{\beta-1} f(v) dv$

- additional term  $\mathcal{S}_\Lambda^{\text{osc}}$  coming from series of contributions of poles of zeta function off the real line: **oscillatory terms**

## Oscillatory terms (fractals)

- zeta function  $\zeta_{\mathcal{L}}(s)$  on fractals in general has additional poles off the real line (position depends on Hausdorff and spectral dimension: depending on how homogeneous the fractal)
- best case exact self-similarity:  $s = \sigma + \frac{2\pi im}{\log \ell}$ ,  $m \in \mathbb{Z}$
- heat kernel on fractals has additional log-oscillatory terms in expansion

$$\frac{C}{t^\sigma} \left( 1 + A \cos\left(\frac{2\pi}{\log \ell} \log t + \phi\right) \right) + \dots$$

for constants  $C, A, \phi$ : series of terms for each complex pole

**Log-oscillatory terms** in expansion of the spectral action:

- G.V. Dunne, *Heat kernels and zeta functions on fractals*, J. Phys. A 45 (2012) 374016 [22p]
- M. Eckstein, B. Iochum, A. Sitarz, *Heat kernel and spectral action on the standard Podleś sphere*, Comm. Math. Phys. 332 (2014) 627–668
- M. Eckstein, A. Zającz, *Asymptotic and exact expansion of heat traces*, arXiv:1412.5100

effect of product with  $S_\beta^1$  (leading term without oscillations)

- case of  $S_\beta^1 \times S_a^3$  (Chamseddine–Connes)

$$D_{S_\beta^1 \times S_a^3} = \begin{pmatrix} 0 & D_{S_a^3} \otimes 1 + i \otimes D_{S_\beta^1} \\ D_{S_a^3} \otimes 1 - i \otimes D_{S_\beta^1} & 0 \end{pmatrix}$$

Spectral action

$$\mathrm{Tr}(h(D_{S_\beta^1 \times S_a^3}^2/\Lambda)) \sim 2\beta\Lambda \mathrm{Tr}(\kappa(D_{S_a^3}^2/\Lambda)),$$

test function  $h(x)$ , and test function

$$\kappa(x^2) = \int_{\mathbb{R}} h(x^2 + y^2) dy$$

- Case of  $S_{\beta}^1 \times \mathcal{P}$ :

$$\begin{aligned} \mathcal{S}_{S_{\beta}^1 \times \mathcal{P}}(\Lambda) &\sim 2\beta \left( \Lambda^4 \zeta_{\mathcal{L}}(3) \mathfrak{h}_3 - \Lambda^2 \frac{1}{4} \zeta_{\mathcal{L}}(1) \mathfrak{h}_1 \right) \\ &+ 2\beta \Lambda^{\sigma+1} \left( \zeta\left(\sigma - 2, \frac{3}{2}\right) - \frac{1}{4} \zeta\left(\sigma, \frac{3}{2}\right) \right) \mathcal{R}_{\sigma} \mathfrak{h}_{\sigma} \end{aligned}$$

with momenta

$$\mathfrak{h}_3 := \pi \int_0^{\infty} h(\rho^2) \rho^3 d\rho, \quad \mathfrak{h}_1 := 2\pi \int_0^{\infty} h(\rho^2) \rho d\rho$$

$$\mathfrak{h}_{\sigma} = 2 \int_0^{\infty} h(\rho^2) \rho^{\sigma} d\rho$$



## Interpretation:

- Term  $2\Lambda^4\beta a^3\mathfrak{h}_3 - \frac{1}{2}\Lambda^2\beta a\mathfrak{h}_1$ , cosmological and Einstein–Hilbert terms, replaced by

$$2\Lambda^4\beta\zeta_{\mathcal{L}}(3)\mathfrak{h}_3 - \frac{1}{2}\Lambda^2\beta\zeta_{\mathcal{L}}(1)\mathfrak{h}_1$$

zeta regularization of divergent series of spectral actions of 3-spheres of packing

- Additional term in gravity action functional: corrections to gravity from fractality

$$2\beta\Lambda^{\sigma+1}\left(\zeta\left(\sigma-2, \frac{3}{2}\right) - \frac{1}{4}\zeta\left(\sigma, \frac{3}{2}\right)\right)\mathcal{R}_{\sigma}\mathfrak{h}_{\sigma}$$

## Case of fractal dodecahedral space $\mathcal{P}_Y$

- Zeta functions

$$\zeta_{\mathcal{L}(\mathcal{P}_Y)}(s) = \sum_{n \geq 0} 20^n (2 + \phi)^{-ns}$$

$$\zeta_{\mathcal{D}_{\mathcal{P}_Y}}(s) = \frac{a^s}{120} \left( 2\zeta\left(s-2, \frac{3}{2}\right) - \frac{1}{2}\zeta\left(s, \frac{3}{2}\right) \right) \zeta_{\mathcal{L}(\mathcal{P}_Y)}(s)$$

- Spectral action:

$$\begin{aligned} \mathrm{Tr}(f(\mathcal{D}_{\mathcal{P}_Y}/\Lambda)) &\sim (\Lambda a)^3 \frac{\zeta_{\mathcal{L}(\mathcal{P}_Y)}(3)}{120} f_3 - \Lambda a \frac{\zeta_{\mathcal{L}(\mathcal{P}_Y)}(1)}{120} f_1 \\ &+ (\Lambda a)^\sigma \frac{\zeta\left(\sigma-2, \frac{3}{2}\right) - \frac{1}{4}\zeta\left(\sigma, \frac{3}{2}\right)}{120 \log(2+\phi)} f_\sigma + \mathcal{S}_{Y,\Lambda}^{\mathrm{osc}} \end{aligned}$$

$$\sigma = \dim_H(\mathcal{P}_Y) = \frac{\log(20)}{\log(2+\phi)} = 2.3296\dots$$

- on product geometry  $S_{\beta}^1 \times \mathcal{P}_Y$

$$\mathcal{S}_{S_{\beta}^1 \times \mathcal{P}_Y}(\Lambda) \sim 2\beta \left( \Lambda^4 \frac{a^3 \zeta_{\mathcal{L}(\mathcal{P}_Y)}(3)}{120} \mathfrak{h}_3 - \Lambda^2 \frac{a \zeta_{\mathcal{L}(\mathcal{P}_Y)}(1)}{120} \mathfrak{h}_1 \right) \\ + 2\beta \Lambda^{\sigma+1} \frac{a^{\sigma} (\zeta(\sigma-2, \frac{3}{2}) - \frac{1}{4} \zeta(\sigma, \frac{3}{2}))}{120 \log(2+\phi)} \mathfrak{h}_{\sigma} + \mathcal{S}_{S_{\beta}^1 \times Y, \Lambda}^{\text{osc}}$$

- **Note:** correction term now at different  $\sigma$  than Apollonian  $\mathcal{P}$
- oscillatory terms  $\mathcal{S}_{Y, \Lambda}^{\text{osc}}$  more explicit than in the Apollonian case

## Oscillatory terms: dodecahedral case

- zeros of zeta function  $\zeta_{\mathcal{L}}(s)$

$$s_m = \sigma + \frac{2\pi im}{\log(2 + \phi)}, \quad m \in \mathbb{Z}$$

with  $\sigma = \log(20)/\log(2 + \phi)$

- contribution to heat kernel expansion of non-real zeros:

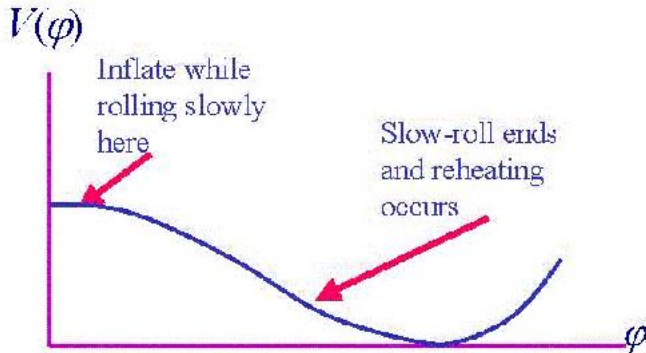
$$\frac{C}{t^\sigma} (a_0 + 2\Re(a_1 t^{-2\pi i/\log(2+\phi)}) + \dots)$$

with coefficients  $a_m$  proportional to  $\Gamma(s_m)$ : for fixed real part  $\sigma$  decays exponentially fast along vertical line

- oscillatory terms are small

## Slow-roll inflation potential from the spectral action

- perturb the Dirac operator by a scalar field  $D^2 + \phi^2 \Rightarrow$  spectral action gives potential  $V(\phi)$



- shape of  $V(\phi)$  distinguishes most cosmic topologies: spherical forms and Bieberbach manifolds (Marcolli, Pierpaoli, Teh)

## Fractality corrections to potential $V(\phi)$

- additional term in potential

$$\mathcal{U}_\sigma(x) = \int_0^\infty u^{(\sigma-1)/2} (h(u+x) - h(u)) du$$

depends on  $\sigma$  fractal dimension

- size of correction depends on (leading term)

$$\left(\zeta\left(\sigma - 2, \frac{3}{2}\right) - \frac{1}{4}\zeta\left(\sigma, \frac{3}{2}\right)\right)\mathcal{R}_\sigma$$

- further corrections to  $\mathcal{U}_\sigma$  come from the oscillatory terms

$\Rightarrow$  presence of fractality (in this spectral action model of gravity)  
can be read off the slow-roll potential (hence the slow-roll coefficients, which depend on  $V$ ,  $V'$ ,  $V''$ )