Measurement of a Sign-Changing Two-Gap Superconducting Phase in Electron-Doped Ba(Fe\textsubscript{1-x}Co\textsubscript{x})\textsubscript{2}As\textsubscript{2} Single Crystals using Scanning Tunneling Spectroscopy

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(Dated: January 27, 2011)

Scanning tunneling spectroscopic studies of Ba(Fe\textsubscript{1-x}Co\textsubscript{x})\textsubscript{2}As\textsubscript{2} (x = 0.06, 0.12) single crystals reveal direct evidence for predominantly two-gap superconductivity. These gaps decrease with increasing temperature and vanish above the superconducting transition $T_c$. The two-gap nature and the slightly doping- and energy-dependent quasiparticle scattering interferences near the wavevectors ($\pm \pi, 0$) and (0, $\pm \pi$) are consistent with sign-changing s-wave superconductivity. The excess zero-bias conductance and the large gap-to-$T_c$ ratios suggest dominant unitary impurity scattering.

PACS numbers: 74.55.+V, 74.70.Xa, 74.25.Jb

The recent discovery of iron-based superconductors [1–10] has renewed intense research activities in superconductivity. Comparison of the similarities and contrasts between the iron-based compounds and the cuprates can provide useful insights into the microscopic mechanism for high-temperature superconductivity [11]. In particular, antiferromagnetic spin fluctuations appear to influence the physical properties of both cuprate and iron-based superconductors [11–16]. On the other hand, theoretical calculations [8, 15–17] and angle resolved photoemission spectroscopy (ARPES) [18–20] suggest that multi-bands and inter-Fermi surface interactions are crucial to superconductivity in the iron-based compounds.

While substantial experimental results from ARPES studies [18–20] and phase sensitive measurements of various single crystalline iron-based compounds [21, 22] are supportive of the scenario of two-gap superconductivity with sign-changing s-wave (s\textsuperscript{h}) order parameters for the hole and electron Fermi pockets, tunneling and point-contact spectroscopic studies of the iron pnictides appear to be inconclusive [23–28]. For instance, reports of point-contact spectroscopy and scanning tunneling spectroscopy (STS) have suggested either BCS-like s-wave superconductivity [23] or nodes in the superconducting order parameter [24, 28] of the “1111” iron-arsenides LnFeAsO\textsubscript{1−x}F\textsubscript{x} (Ln: trivalent rare-earth elements). On the other hand, STS studies of the electron and hole-doped “122” iron arsenides Ba(Fe\textsubscript{1-x}Co\textsubscript{x})\textsubscript{2}As\textsubscript{2} and (Ba\textsubscript{1−x}K\textsubscript{x})Fe\textsubscript{2}As\textsubscript{2} have revealed findings ranging from strong spatial variations in the tunneling spectra with occasional observation of a large superconducting gap [25] to moderate spatial variations with predominantly a small superconducting gap [26]. In the non-superconducting limit, nematic surface reconstructions have been observed [27].

In this letter we report direct STS evidence for two-gap superconductivity in the electron-doped 122 system Ba(Fe\textsubscript{1-x}Co\textsubscript{x})\textsubscript{2}As\textsubscript{2} of two different doping levels. For each doping level, two different energy gaps can be clearly resolved at $T \ll T_c$. Both gaps decrease monotonically with increasing temperature and then completely vanish above $T_c$. The gap values agree favorably with those obtained from ARPES so that the larger gap $\Delta$ may be associated with the hole-like Fermi pockets and the smaller gap $\Delta_M$ with the electron-like Fermi pockets. Moreover, Fourier transformation (FT) of the tunneling conductance reveals energy and doping dependent quasiparticle scattering interferences (QPI) near the nesting wavevectors ($\pm \pi, 0$) and (0, $\pm \pi$) between the Fermi pockets at $\Gamma$ and $M$, which is consistent with the sign-changing order parameters for the hole and electron pockets [17]. Finally, excess zero-bias tunneling conductance and the large $2\Delta/k_BT_c$ ratios for both doping levels may be attributed to strong unitary impurity scattering [29].

The Ba(Fe\textsubscript{1−x}Co\textsubscript{x})\textsubscript{2}As\textsubscript{2} samples investigated in this work are single crystals with $x = 0.06$ (underdoped) and 0.12 (overdoped), and the corresponding superconducting transition temperatures are $T_c = 14$ and 20 K, respectively. The single crystals were grown from the flux method [30], and details of the synthesis and characterization of the samples have been described elsewhere [30–32]. Given the reactive nature of Ba(Fe\textsubscript{1−x}Co\textsubscript{x})\textsubscript{2}As\textsubscript{2}, freshly cleaved surfaces were essential for the STS studies. To date all reported STS studies were carried out on samples that were mechanically cleaved under ultrahigh vacuum (UHV) conditions and at cryogenic temperatures [25–27], and the resulting sample surfaces all exhibited significant (2 $\times$ 1) reconstructions of the Fe(Co)-layer. We chose to perform mechanical cleavage of the single crystals in pure argon atmosphere at room temperature, well above the tetragonal-to-orthorhombic structural phase transition. The cleaved samples were loaded in situ onto the cryogenic probe of our homemade scanning tunneling microscope (STM) in argon. The sealed STM assembly was subsequently evacuated and cooled to 6 K in UHV with a base pressure at $\sim 10^{-10}$ Torr.

Both spatially resolved topography and tunnelling conductance ($dI/dV$) versus energy ($E = eV$) spectroscopy were acquired pixel-by-pixel over an extended area of each sample simultaneously, with tunneling cur-
We subdivided the surface topography into $5 \times 5$ nm$^2$ squares. Detailed survey of the surface topography and tunneling conductance spectra was carried out over typically 5 meV. The symbols and solid lines indicate the width of the gap distributions obtained from the fitting using Eq. (1). The results shown in Figs. 1(a)-(b) are representative of the spectral characteristics of both samples, as manifest by the gap spatial maps and the corresponding histograms for both samples in Figs. 2(a)-(b), where the gap values are empirically determined as one-half of the peak-to-peak values. Overall, $\langle \Delta_T \rangle = 10.0$ meV and $\langle \Delta_M \rangle = 5.0$ meV for $x = 0.12$ and $\langle \Delta_T \rangle = 8.0$ meV and $\langle \Delta_M \rangle = 4.0$ meV for $x = 0.06$.

Our findings of the two-gap spectra differ from previous STS studies [25, 26], which may be the result of differences in the surface preparation. Specifically, we show in Fig. 3(a) an example of the atomically reconstructed surfaces, as exemplified by the height histogram in Fig. 3(b) for the same ($5 \times 5.4$) nm$^2$ area. In particular, we note that the topography in Fig. 3(a) exhibited no apparently reconstructed ($1 \times 2$) surfaces, and the overall height variations were limited to within one c-axis lattice constant $c_0 = 1.239$ nm. This finding is in contrast to the mostly flat and reconstructed surfaces reported by other groups for cold-cleaved samples [25–27].

Next, we employ a phenomenological fitting general-
FIG. 2. (color online) Superconducting gap maps and histograms at T = 6 K: (a) Left to right: The first two panels correspond to the $\Delta_M$ and $\Delta_T$ maps for the underdoped sample (x = 0.06), and the right two panels represent the corresponding histograms for both the quasiparticle (solid bars) and quasihole (shaded bars) branches, showing particle-hole symmetry and the mean values of $\langle |\Delta_M| \rangle = 4$ meV and $\langle |\Delta_T| \rangle = 8$ meV. (b) The left two panels are respectively the $\Delta_M$ and $\Delta_T$ maps for the sample with x = 0.12. The right two panels are histograms for both the quasiparticle (solid bars) and quasihole (shaded bars) branches, showing particle-hole symmetry and $\langle |\Delta_M| \rangle = 5$ meV, $\langle |\Delta_T| \rangle = 10$ meV.

FIG. 3. (color online) Correlation of atomically resolved surface topography with tunneling conductance spectra: (a) Surface topography of the overdoped sample (x = 0.12) over a (5.4 x 5.4)nm$^2$ area. (b) The height histogram of the area in (a), showing height variations within one lattice constant along the c-axis. (c) Spatial evolution of the normalized $(dI/dV)$-vs.-$V$ spectra across a horizontal line slightly below the middle of the topography image in (a), which reveals consistent two-gap features.

Despite the fragmented surfaces, consistent two-gap spectral features were found throughout the scanned area, as exemplified in Fig. 3(c) for the tunneling spectra taken near the middle of Fig. 3(a).

Next, we examine the spatial variations in the tunneling conductance at constant bias voltages. In Figs. 4(a)-(b), we illustrate the tunneling conductance maps of two samples and under three different bias voltages of $V = 0$, $\langle |\Delta_M| \rangle/e$ and $\langle |\Delta_T| \rangle/e$. The Fourier transformed (FT) tunneling conductance at $V = \langle |\Delta_M| \rangle/e$ and $\langle |\Delta_T| \rangle/e$ is shown in Fig. 4(c) for $x = 0.06$ and in Fig. 4(d) for $x = 0.12$. We find that dominant QPI occur at three wave-vectors [22]: $q_1$ between two electron pockets across the first Brillouin zone, $q_1 \sim (\pm 2\pi, 0)/(0, \pm 2\pi)$; $q_2$ between the hole- and electron-pockets at $\Gamma$ and M-points, $q_2 \sim (\pm \pi, 0)/(0, \pm \pi)$; and $q_3$ between two adjacent electron pockets, $q_3 \sim (\pm \pi, \pm \pi)$. For the underdoped sample, strong QPI at $q_2$ and absence of QPI at $q_1$ is consistent with the five-orbital theoretical calculations [15, 16, 34] for two possible scenarios: one is scalar-impurity QPI between sign-reversing order parameters associated with the hole- and electron-pockets, and the other is magnetic impurity QPI between order parameters of the same sign. Given that scalar scatterers are generally more common than magnetic impurities, we suggest that the prevailing QPI wave-vectors at $q_2$ are the result of sign-reversing order parameters. In this context, the appearance of $q_3$ wave-vectors in the overdoped sample might be due to a larger density of Co-atoms in the surface Fe/Co layers. Unlike those in the bulk, the charge transfer from the surface Co-atoms is incomplete so that they may behave like magnetic impurities [15, 16, 34, 35].

It is worth noting that for both samples the $q_2$ wave-vectors appear to exhibit one preferential direction. In addition, energy-dependent features at small wave-vectors ($|q| < \pi/2$) are observed along the same direction. These findings may be related to the nematic-order wave-vectors found in the parent state of these compounds [27], although our limited momentum resolution cannot provide detailed comparison. Overall, we attribute $q_1$, $q_2$, and $q_3$ to the QPI scattering wave-vectors because they are not only slightly energy dependent but also doping dependent, and so they cannot be simply attributed to Bragg diffractions of the lattices. The relatively weak energy dependence of the QPI wave-vectors is the result of small and nearly isotropic Fermi pockets in the iron arsenides, which differs from the highly energy-dependent QPI wave-vectors in the cuprate superconductors [36–39].

Finally, we note that the relatively high zero-bias conductance in all tunneling spectra at $T \ll T_c$ is suggestive of dominant unitary impurity scattering [29]. Further, unitary impurity effects on the suppression of $T_c$ for the sign-changing s-wave superconductors are found to be as significant as those on the d-wave superconductors [29]. On the other hand, unitary impurity effects on suppressing the sign-changing s-wave order parameters involve sign-dependent components that are partially cancelled, and are therefore weakened [29]. Hence, the large ratios of $(2\langle |\Delta_{T,M}| \rangle)/(k_B T_c)$ for the 122 system, with $(2\langle |\Delta_M| \rangle)/(k_B T_c) \sim 6.6$ (5.8) and $(2\langle |\Delta_T| \rangle)/(k_B T_c) \sim 13.2$ (11.6) for $x = 0.06$ (0.12), may be attributed to significant unitary impurity scattering in these sign-changing s-wave superconductors. A possible source for the unitary impurity scattering may be associated with disorder in the Co doping into the Fe-planes, which is in contrast to the situation in the 1111 iron arsenides where doping
takes place in the charge reservoir [5] so that the disorder effect is weaker and the $(2\Delta)/(k_B T_c)$ ratio is small ($\sim 3$) [23, 28], whereas the $T_c$ values are generally higher.

In summary, we have demonstrated direct STS evidence for two-gap superconductivity in electron-doped Ba(Fe$_{1-x}$Co$_x$)$_2$As$_2$ single crystals of two doping levels. The Fourier transformed tunneling conductance reveals strong quasiparticle scattering interferences near the nesting wave-vector between the hole Fermi pockets at $\Gamma$ and the electron Fermi pockets at M, consistent with sign-changing order parameters of the two Fermi pockets. The excess zero-bias conductance and the large $2\Delta_{T,M}/(k_B T_c)$ ratios for both doping levels may be attributed to significant unitary impurity scattering in a sign-changing s-wave superconducting system.

This work at Caltech was jointly supported by the NSF Grant DMR-0907251, and the Kavli and Moore Foundations. The work in China was supported by the NSFC, the Ministry of Science and Technology of China and Chinese Academy of Sciences within the knowledge innovation program. We thank Patrick A. Lee, Igor Mazin and Yunkyu Bang for valuable discussion.


FIG. 4. (color online) Tunneling conductance maps of Ba(Fe$_{1-x}$Co$_x$)$_2$As$_2$ samples at three constant bias voltages $V = 0$, $(\langle\Delta_M\rangle/e$, $(\langle\Delta_T\rangle/e$ and for $T = 6$ K: (a) $x = 0.06$ and (b) $x = 0.12$. (c) Fourier transformed (FT) tunneling conductance maps at $V = (\langle\Delta_M\rangle/e = 4$ meV (left panel) and $V = (\langle\Delta_T\rangle/e = 8$ meV (right panel) for $x = 0.06$. (d) FT conductance maps at $V = (\langle\Delta_M\rangle/e = 5$ meV (left panel) and $V = (\langle\Delta_T\rangle/e = 10$ meV (right panel) for $x = 0.12$. The thin white lines define the first and second Brillouin zones, and we have used the lattice constant $a = 0.53$ nm for the length of the first Brillouin zone (2$\pi/a$). All three QPI wave-vectors $q_1$, $q_2$ and $q_3$ are slightly energy and doping dependent.