

Lecture 13: Oscillatory Motion II

Energy of Oscillation

Like all physical systems, oscillators have energy.

harmonic
oscillator:

$$\ddot{x} = -\frac{k}{m}x$$

↓ multiply by \dot{x} and integrate

$$\int (\dot{x} \ddot{x} = -\frac{k}{m} x \dot{x}) dt$$

$$\frac{1}{2} \dot{x}^2 = -\frac{k}{m} \frac{x^2}{2} + \text{const}$$

energy
conservation!

$$\underbrace{\frac{m \dot{x}^2}{2}} + \underbrace{\frac{1}{2} k x^2} = \underbrace{\text{const}} = E$$

kinetic
energy

potential
energy

total
energy

Plugging the harmonic oscillator, $X(t) = C \cos(\omega_0 t + \delta)$:

$$E_{HO} = \frac{m}{2} (-C\omega_0 \sin(\dots))^2 + \frac{1}{2} k (C \cos(\dots))^2$$

$$= C^2 \left(\frac{m\omega_0^2}{2} \sin^2(\dots) + \frac{k}{2} \cos^2(\dots) \right)$$

$$= \frac{kC^2}{2} \left(\sin^2(\dots) + \cos^2(\dots) \right)$$

$$\Rightarrow \boxed{E_{HO} = \frac{kC^2}{2}}$$

increasing amplitude increases energy

For the damped harmonic oscillator, the amplitude of oscillation falls exponentially, $\sim Ce^{-\beta t/2}$

$$\Rightarrow \boxed{E_{DHO} \sim \frac{kC^2}{2} e^{-\beta t}}$$

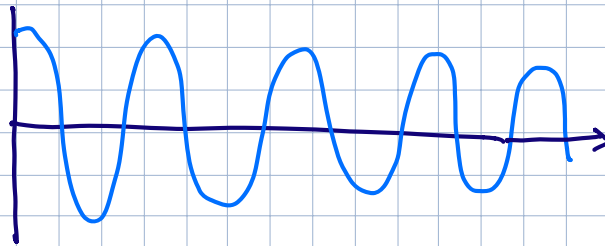
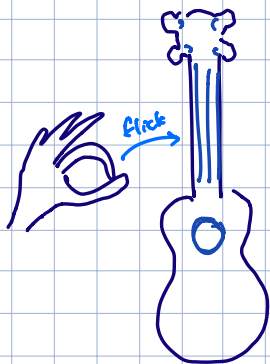
friction dissipates away energy and sets arrow of time

The "Q" Factor

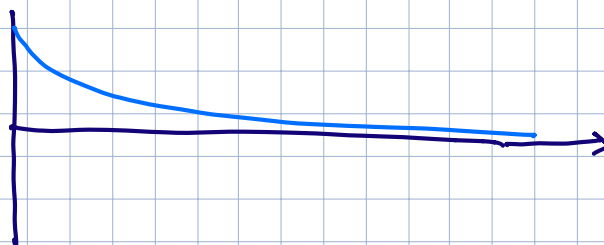
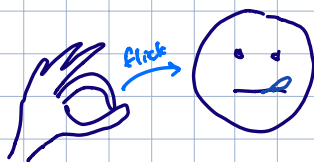
Q = "quality factor"

where $\frac{1}{Q}$ = fraction of energy lost
in one radian of oscillation

high Q (guitar string)



low Q (your face)



From this definition we find

$$Q = \frac{E(t)}{E(t) - E(t + \frac{1}{\omega_0})}$$

one radian later

$$= \frac{e^{-\beta t}}{e^{-\beta t} - e^{-\beta(t + \frac{1}{\omega_0})}}$$

For high quality oscillators, $\beta \ll \omega_0$, so $e^{-\beta/\omega_0} \sim 1 - \beta/\omega_0$

$$\rightarrow Q \sim \frac{e^{-\beta t}}{e^{-\beta t} - e^{-\beta t} (1 - \beta/\omega_0)}$$

$$Q \sim \frac{\omega_0}{\beta}$$

Driven Damped Harmonic Oscillator

$$m \ddot{x} = F = -kx - \underbrace{\gamma \dot{x}}_{\text{friction}} + \underbrace{F_0 \sin \omega t}_{\text{external drive}}$$

$$\Rightarrow \ddot{x} + \beta \dot{x} + \omega_0^2 x = a_0 \sin \omega t$$

$a_0 = F_0/m$

As usual, behold the power of guess and check:

$$x(t) \stackrel{??}{=} A \sin(\omega t - \alpha)$$

frequency of drive

$$\dot{x} = \omega A \cos(\omega t - \alpha)$$

$$\ddot{x} = -\omega^2 A \sin(\omega t - \alpha)$$

Rewrite $\sin \omega t = \sin(\omega t - \alpha + \alpha)$

$$= \sin(\omega t - \alpha) \cos \alpha + \cos(\omega t - \alpha) \sin \alpha$$

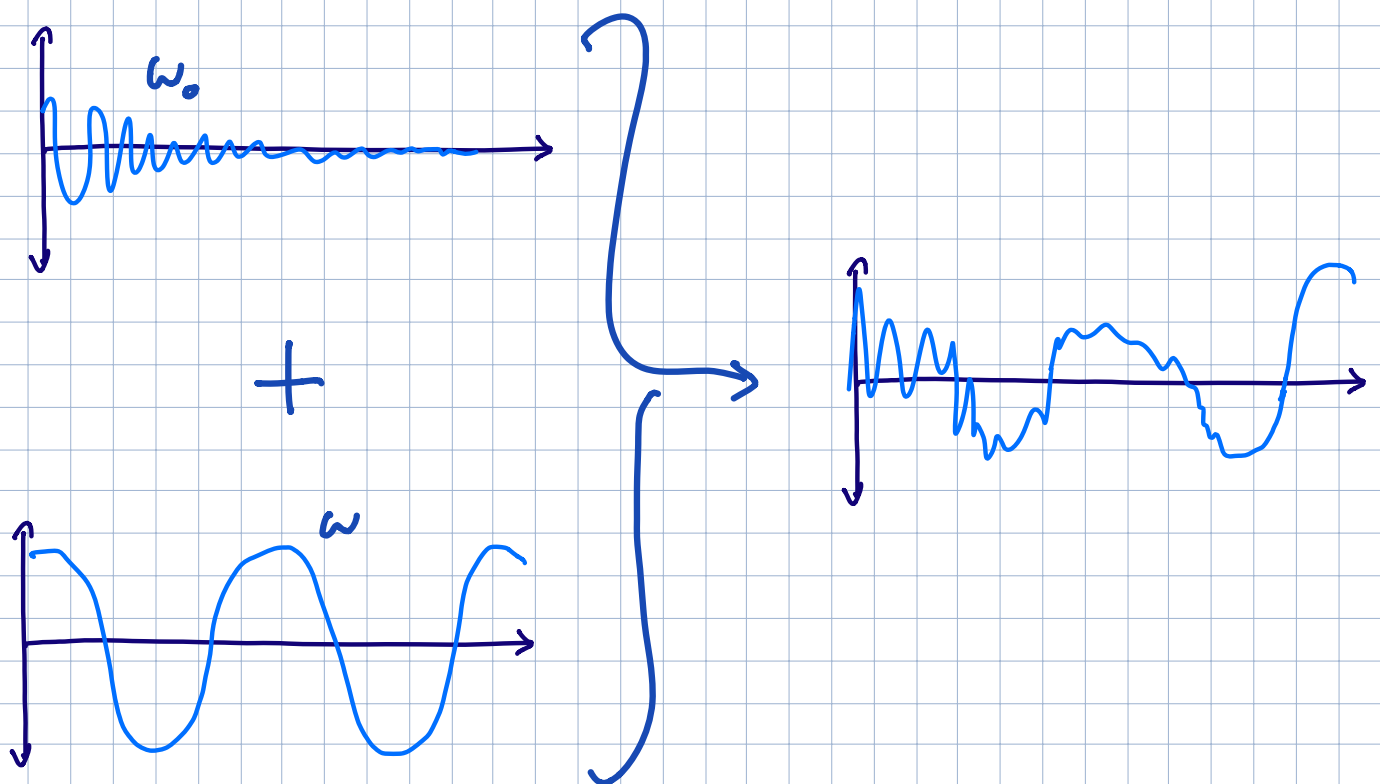
We can add this "particular" solution to the "homogeneous" solution we computed before.

$$X_{\text{full}}(t) = X_{\text{hom}}(t) + X_{\text{part}}(t)$$

$$C e^{-\beta t/2} \cos(\omega_1 t + \phi) + A \sin(\omega t - \alpha)$$

fixed by initial cond fixed by system

At late times, the $e^{-\beta t/2}$ squashes one solution.



We saw before that the homogeneous solution dies off from damping while the particular solution persists.

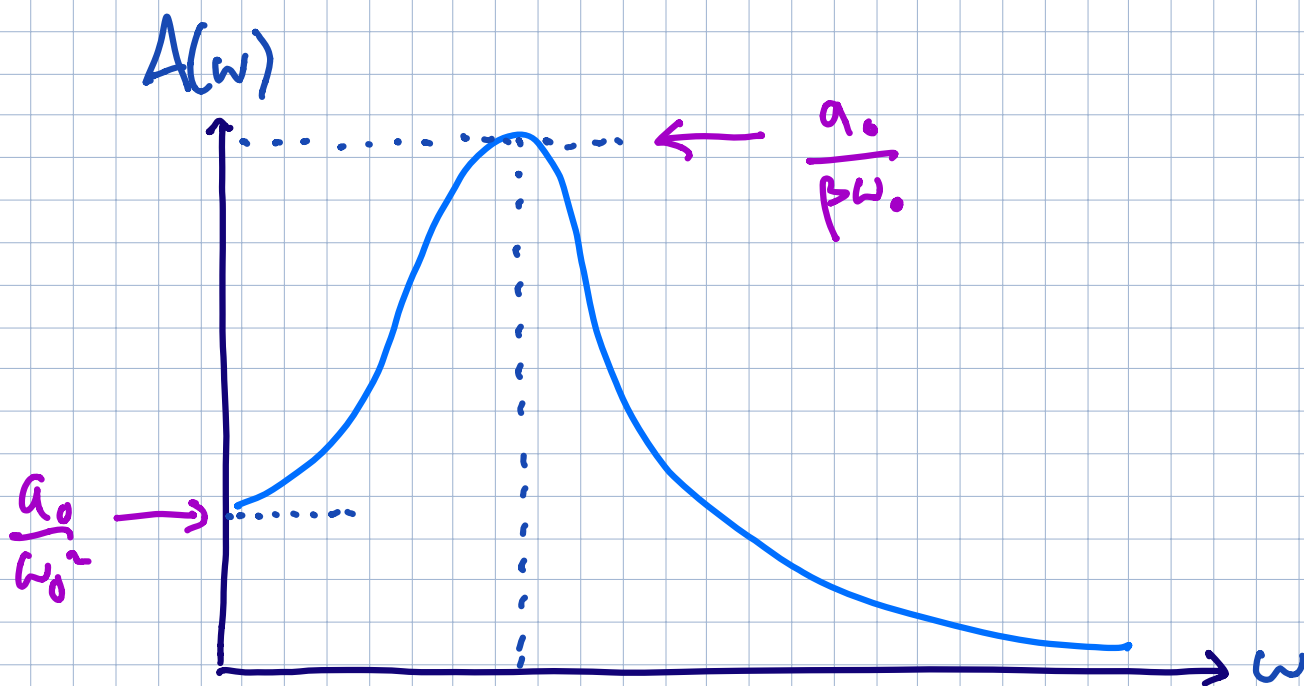
⇒
particular
or
"steady state"
solution

$$X(t) = A \sin(\omega t - \alpha)$$

$$A = \frac{a_0}{\sqrt{(\beta\omega)^2 + (\omega_0^2 - \omega^2)^2}}$$

$$\tan \alpha = \frac{\beta\omega}{\omega_0^2 - \omega^2}$$

- Consider amplitude $A(\omega)$ as a function of ω



demo: "driven pendulum"

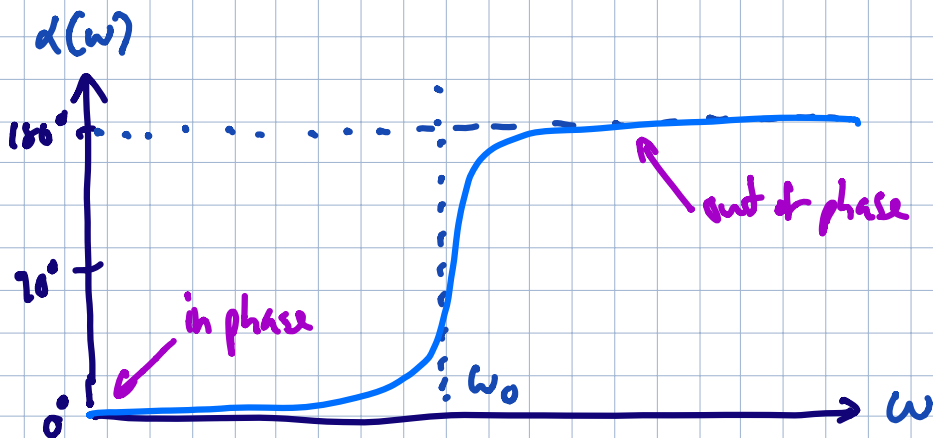
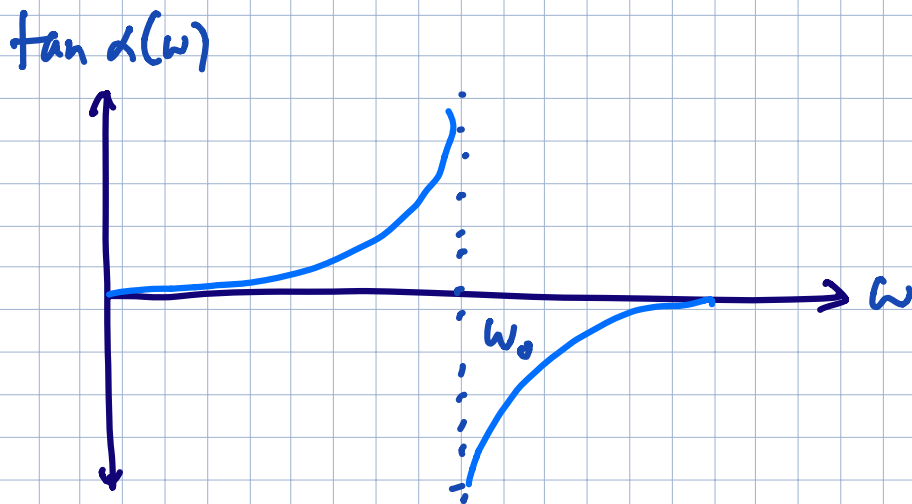
For $\beta \ll \omega_0$, the maximum is at $\omega = \omega_0$.

Note also that:

$$\frac{A(\omega_0)}{A(0)} = \frac{a_0 / \beta \omega_0}{a_0 / \omega_0} = \frac{\omega_0}{\beta} = Q$$

$\Rightarrow Q \sim$ height of resonance peak

- Consider phase lag $\alpha(\omega)$ as a function of ω



demo: "shattering wine glass"