Optimal scaling laws in ductile fracture

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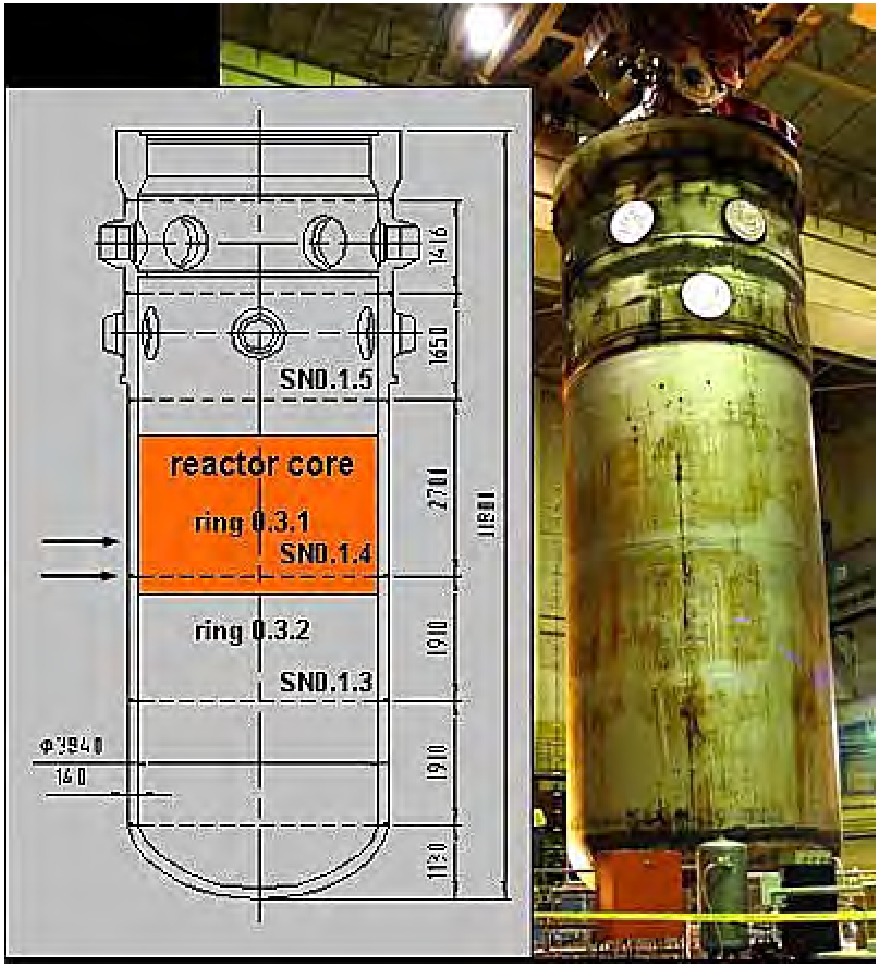
Symposium on
New Developments in Defects Mechanics
Chair: Xanthippi Markenscoff, UCSD
University of California, San Diego
La Jolla, CA, January 18-19, 2014

Outline

- Background/phenomenology of ductile fracture
- Metals: mathematical formulation
- Optimal scaling laws
- Numerical approximation
- Applications: Hypervelocity impact and explosively-driven caps
- Extension to polymers
- Application: Taylor anvil tests on polyurea

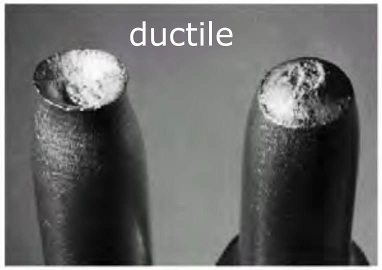
Background on ductile fracture



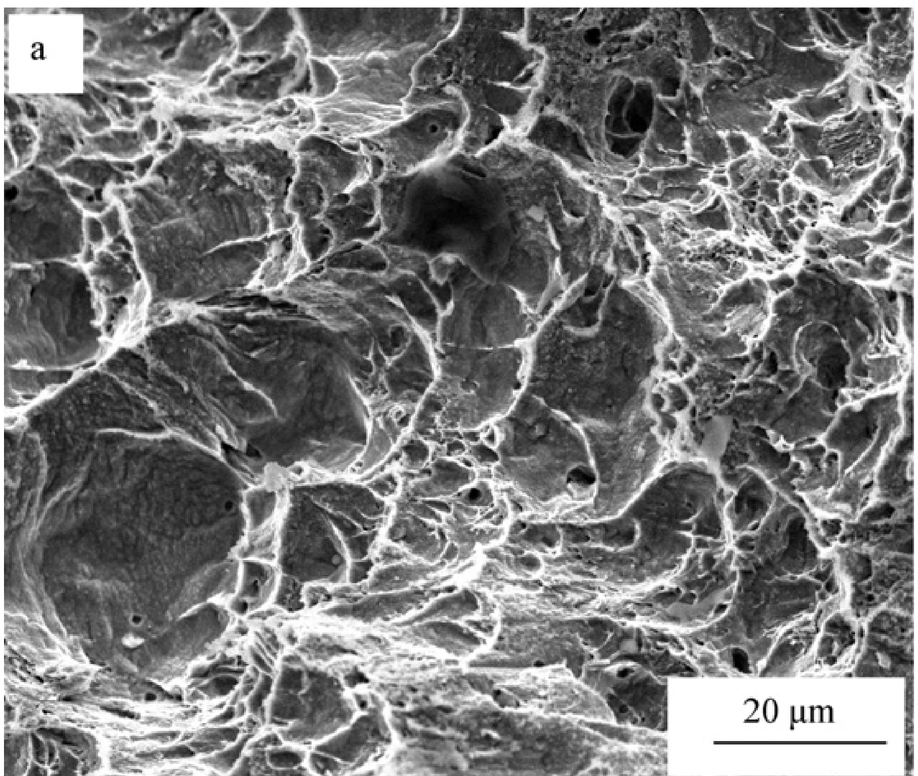
- Linear-elastic fracture mechanics proved inadequate for assessing, e.g., safety of mild-steel pressure vessels in nuclear power plants, which spurred the development of ***elastic-plastic fracture mechanics*** (with focus on Rice's J-integral formalism)

Reactor Pressure Vessel (RPV)
from Greifswald Nuclear Power Plant
(courtesy Viehrig, H.W. and Houska, M.,
Helmholtz Zentrum, Dresden-Rossendorf,
https://www.hzdr.de/db/Cms?pNid=2698)

Background on ductile fracture



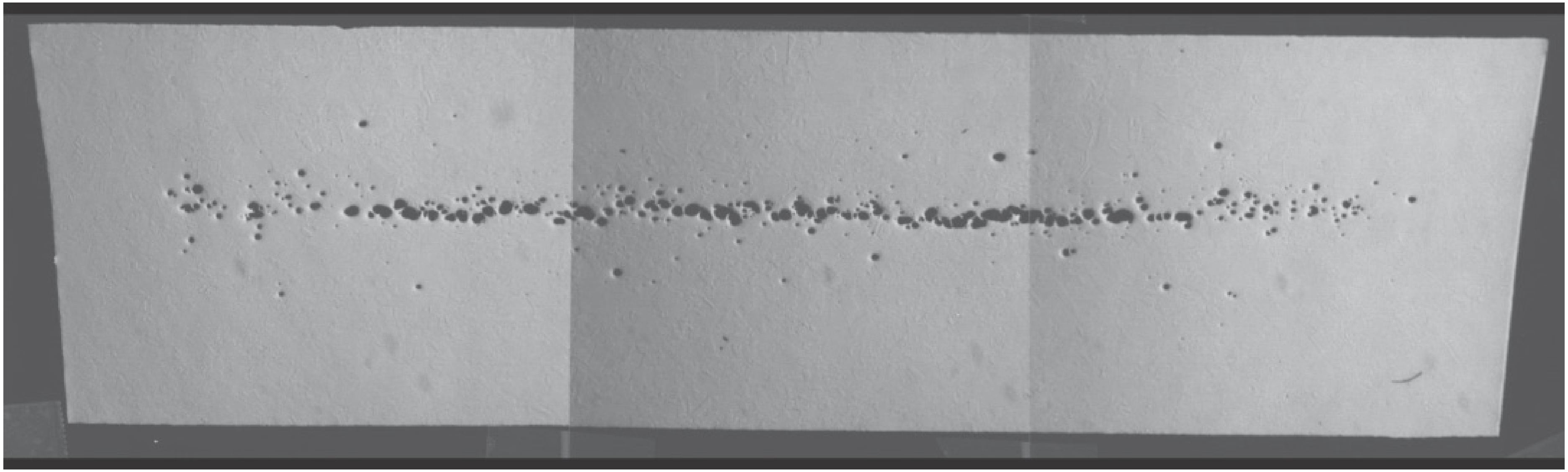
(Courtesy NSW HSC online)



- Ductile fracture in metals occurs by *void nucleation, growth and coalescence*
- Fractography of ductile-fracture surfaces exhibits profuse *dimpling*, vestige of microvoids
- Ductile fracture entails large amounts of *plastic deformation* (*vs*. surface energy) and dissipation.

Fracture surface in SA333 steel, room temp., $d\varepsilon/dt=3\times10^{-3}s^{-1}$ (S.V. Kamata, M. Srinivasa and P.R. Rao, Mater. Sci. Engr. A, ***528*** *(2011) 4141–4146)*

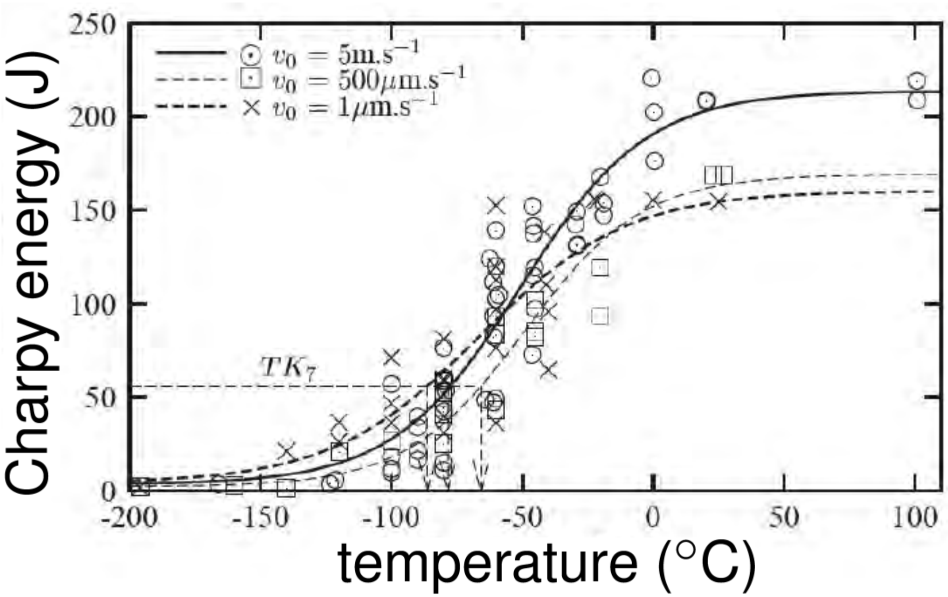
Background on ductile fracture



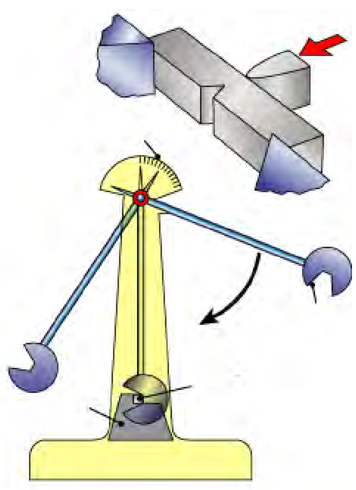
Photomicrograph of a copper disk tested in a gas-gun experiment showing the formation of voids and their coalescence into a fracture plane

Heller, A., How Metals Fail,
Science & Technology Review Magazine,
Lawrence Livermore National Laboratory,
pp. 13-20, July/August, 2002

Background on ductile fracture



Charpy energy of
A508 steel
(Tanguy *et al.*, *Eng. Frac. Mechanics*, 2005)



- A number of ASTM engineering standards are in place to characterize ductile fracture properties (J-testing, Charpy test)
- The Charpy test data reveals a brittle-to-ductile transition temperature
- In general, the specific fracture energy for ductile fracture is greatly in excess of that required for brittle fracture…

Outline/work plan

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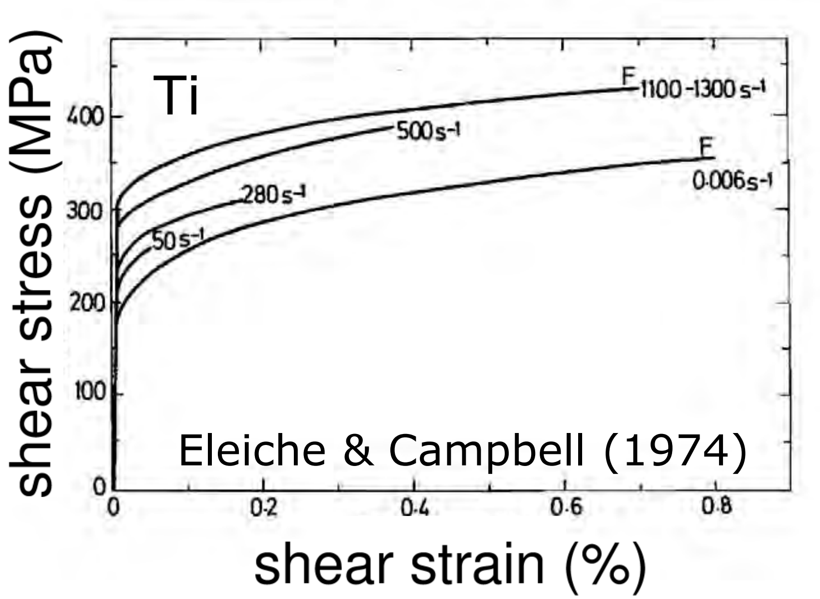
Naïve model: Local plasticity

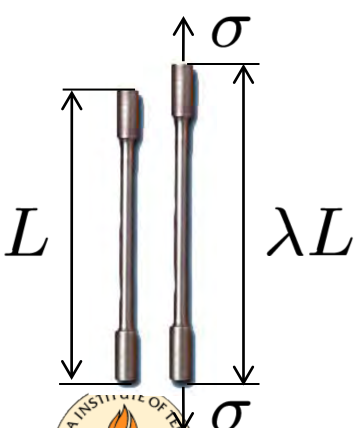


- Deformation theory: Minimize

$$E(y) = \int_\Omega W\Big(Dy(x)\Big)\,dx$$

- Growth of $W(F)$?
- Asume power-law hardening:

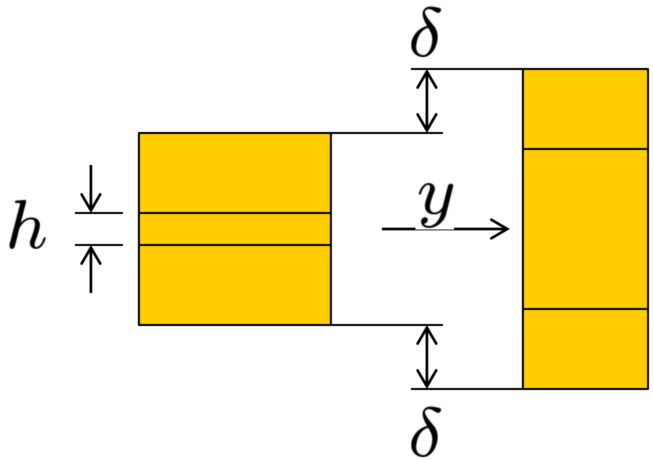
$$\sigma \sim K\epsilon^n = K(\lambda - 1)^n$$



- Nominal stress: $\partial_\lambda W = \sigma/\lambda = K(\lambda-1)^n/\lambda$
- For large λ: $\partial_\lambda W \sim K\lambda^{n-1} \Rightarrow W \sim K\lambda^n$
- In general: $W(F) \sim |F|^p,\ p = n \in (0,1)$

$\Rightarrow$ Sublinear growth!

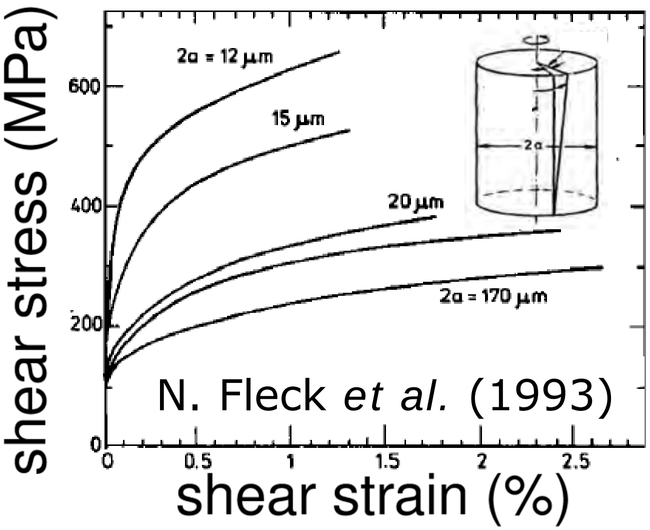
Naïve model: Local plasticity

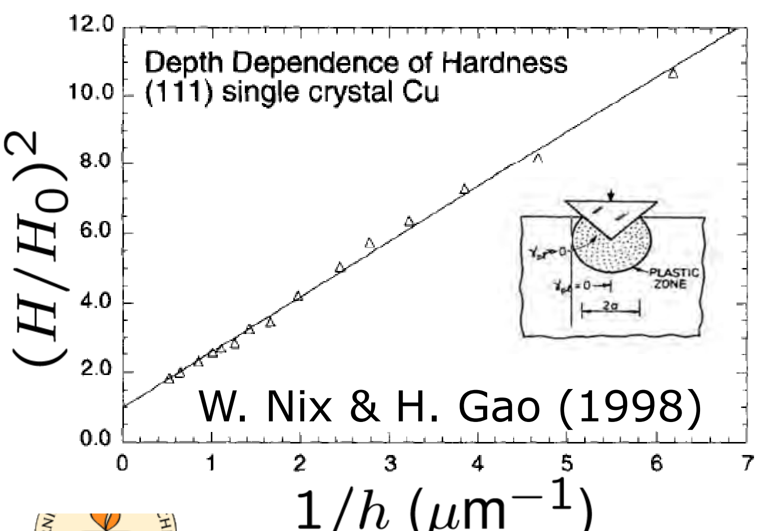


- Example: Uniaxial extension
- Energy: $E_h \sim h \left(\frac{2\delta}{h}\right)^p$
- For $p < 1$: $\lim_{h \to 0} E_h = 0$

- Energies with sublinear growth relax to 0.
- *For hardening exponents in the range of experimental observation, local plasticity yields no useful information regarding ductile fracture properties of materials!*
- Need additional physics, structure…

Strain-gradient plasticity





- The yield stress of metals is observed to increase in the presence of strain gradients
- Deformation theory of strain-gradient plasticity:

$$E(y) = \int_\Omega W\Big(Dy(x), D^2y(x)\Big)\,dx$$

$y : \Omega \to \mathbb{R}^n$, volume preserving

- Strain-gradient effects may be expected to oppose localization
- Growth of *W* with respect to the second deformation gradient?

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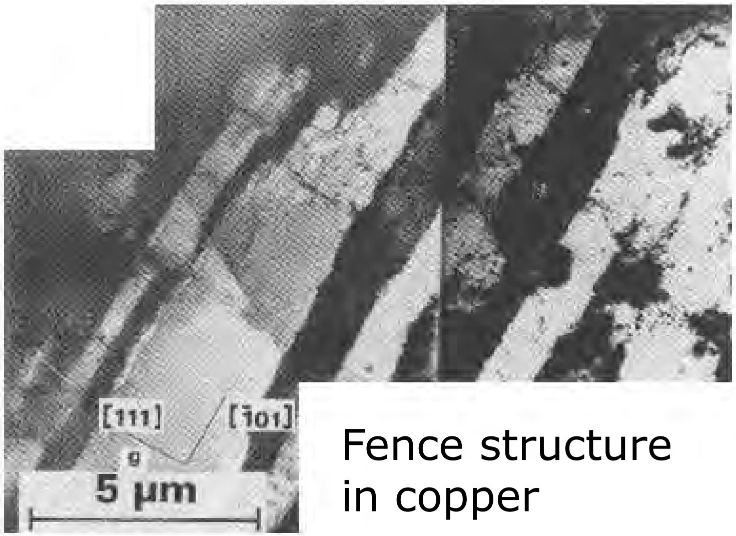
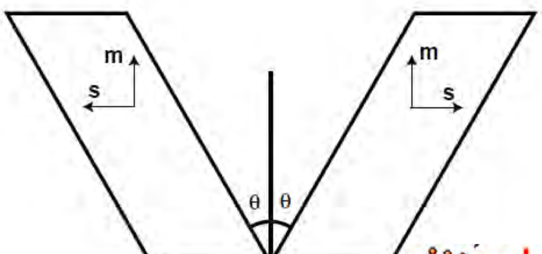
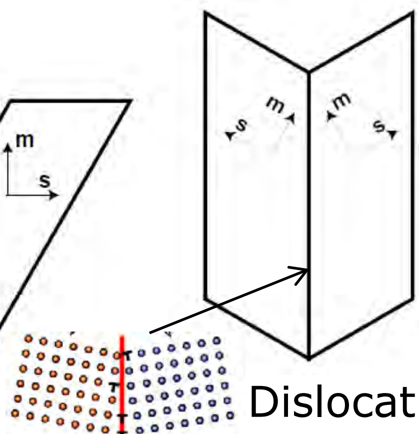
Strain-gradient plasticity



Fence structure
in copper

(J.W. Steeds, *Proc. Roy. Soc. London*, **A292**, 1966, p. 343)





Dislocation wall

- Growth of $W(F, \cdot)$?
- For fence structure:

$$F^{\pm} = R^{\pm}(I \pm \tan\theta \, s \otimes m)$$

- Across jump planes:

$$|[\![F]\!]| = 2\sin\theta$$

- Dislocation-wall energy:

$$E = \frac{T}{b} 2\sin\theta = \frac{T}{b} |[\![F]\!]|$$

$\Rightarrow W(F, \cdot)$ has linear growth!

Strain-gradient plasticity

- Mathematical model: Minimize

$$E(y) = \int_\Omega W\Big(Dy(x), D^2y(x)\Big)\, dx$$

$$y : \Omega \to \mathbb{R}^n, \text{ volume preserving}$$

- For metals, local plasticity exhibits sub-linear growth, which favors localization of deformations
- Strain-gradient plasticity may be expected to exhibit linear growth, which opposes localization
- *Question: Can ductile fracture be understood as the result of a competition between sublinear growth and strain-gradient plasticity?*

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Scaling laws in science

- A broad variety of physical phenomena obey power laws over wide ranges of parameters
- Scale invariance: If $y = C x^q$, then (x,y) iff $(\lambda x, \lambda^q y)$, law of corresponding states
- Universality: Systems displaying identical scaling behavior are likely to obey the same fundamental dynamics
- Experimental master curves, data collapse
- Examples:
  - *Critical phenomena (second-order transitions)*
  - *Materials science (Taylor, Hall-Petch, creep laws...)*
  - *Continuum mechanics (hydrodynamic, fracture...)*

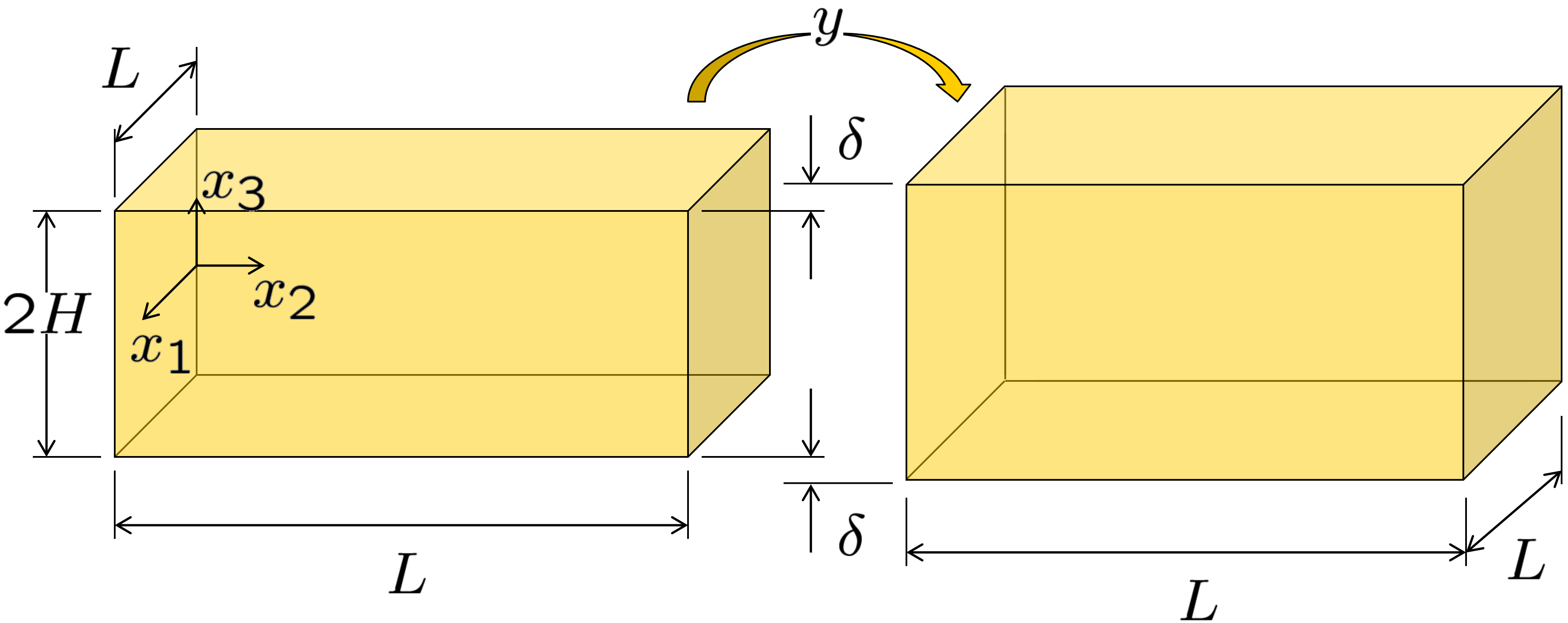
Optimal scaling

- Suppose: Energy $= E(u, \epsilon_1, \ldots, \epsilon_N)$
- Optimal (matching) upper and lower bounds:

$$C_L \epsilon_1^{\alpha_1} \ldots \epsilon_N^{\alpha_N} \leq \inf E(\cdot, \epsilon_1, \ldots, \epsilon_N) \leq C_U \epsilon_1^{\alpha_1} \ldots \epsilon_N^{\alpha_N}$$

- The exponents $\alpha_1, \ldots, \alpha_N$ are *sharp*, *universal*
- The constants C_L and C_U are often lax, imprecise...
- Upper bound by construction, *ansatz*-free lower bound
- Originally applied to branched microstructures in martensite (Kohn-Müller 92, 94; Conti 00)
- Applications to micromagnetics (Choksi-Kohn-Otto 99), thin films (Belgacem *et al* 00)...

Optimal scaling – Ductile fracture



- Approach: Optimal scaling
- Slab: $\Omega = [0, L]^2 \times [-H, H]$, periodic
- Uniaxial extension: $y_3(x_1, x_2, \pm H) = x_3 \pm \delta$

Optimal scaling – Ductile fracture

- $y : \Omega \to \mathbb{R}^3$, $[0, L]^2$-periodic, volume preserving
- $y \in W^{1,1}(\Omega; \mathbb{R}^3)$, $Dy \in BV(\Omega; \mathbb{R}^{3\times 3})$
- Growth: For $0 < K_L < K_U$, *intrinsic length* $\ell > 0$,

$$E(y) \geq K_L \left(\int_\Omega (|Dy|^p - 3^{p/2})\, dx + \ell \int_\Omega |D^2 y|\, dx \right)$$
$$E(y) \leq K_U \left(\int_\Omega (|Dy|^p - 3^{p/2})\, dx + \ell \int_\Omega |D^2 y|\, dx \right)$$

Theorem [Fokoua, Conti & MO, ARMA, 2013]. For ℓ sufficiently small, $p \in (0, 1)$, $0 < C_L(p) < C_U(p)$,

$$C_L(p) L^2 \ell^{\frac{1-p}{2-p}} \delta^{\frac{1}{2-p}} \leq \inf E \leq C_U(p) L^2 \ell^{\frac{1-p}{2-p}} \delta^{\frac{1}{2-p}}$$

Optimal scaling – Ductile fracture

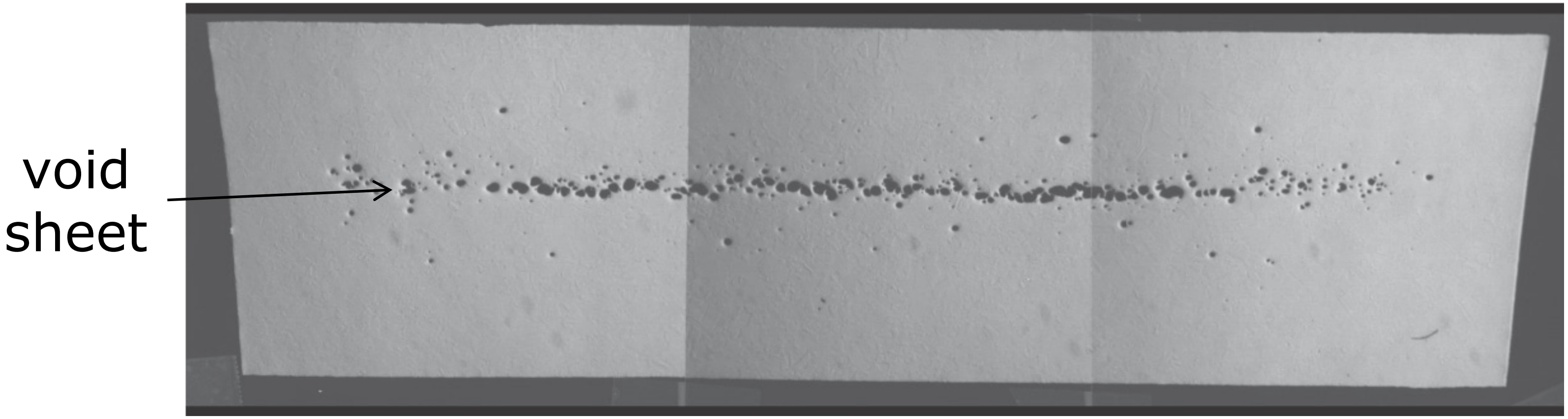
- Optimal (matching) upper and lower bounds:

$$C_L(p) L^2 \ell^{\frac{1-p}{2-p}} \delta^{\frac{1}{2-p}} \leq \inf E \leq C_U(p) L^2 \ell^{\frac{1-p}{2-p}} \delta^{\frac{1}{2-p}}$$

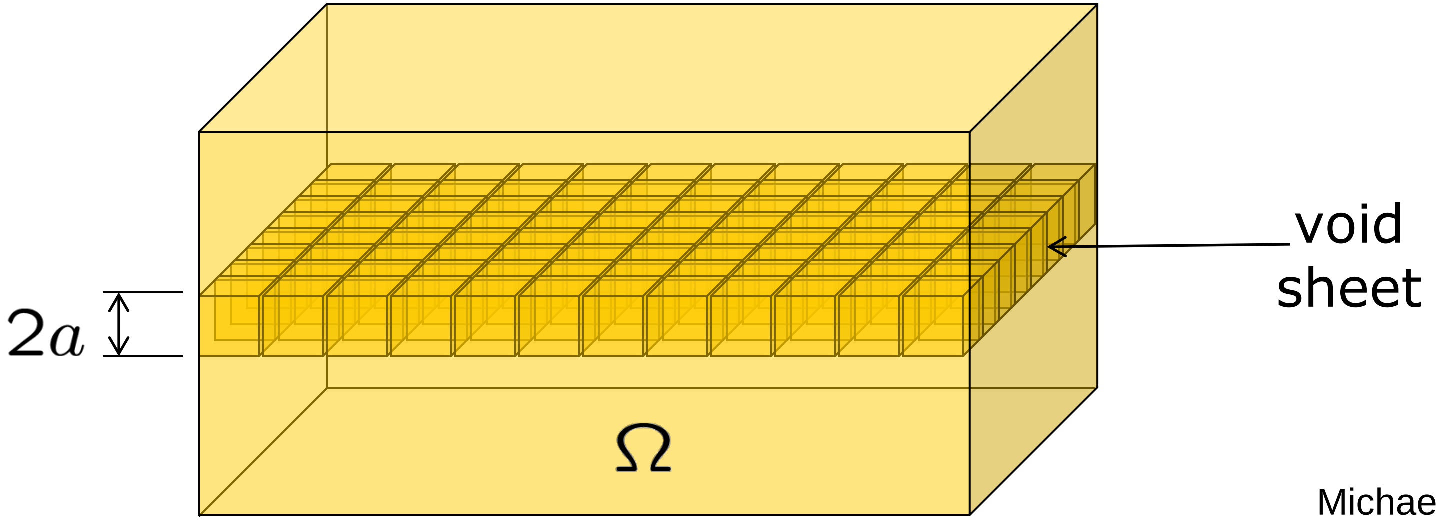
- Bounds apply to *classes of materials* having the same growth, specific model details immaterial
- Energy scales with *area* (L^2): Fracture scaling!
- Energy scales with power of *opening displacement* (δ): Cohesive behavior!
- Lower bound degenerates to 0 when the intrinsic length (ℓ) decreases to zero…
- Bounds on specific fracture energy:

$$C_L(p) \ell^{\frac{1-p}{2-p}} \delta_c^{\frac{1}{2-p}} \leq J_c \leq C_U(p) \ell^{\frac{1-p}{2-p}} \delta_c^{\frac{1}{2-p}}$$

Sketch of proof – Upper bound

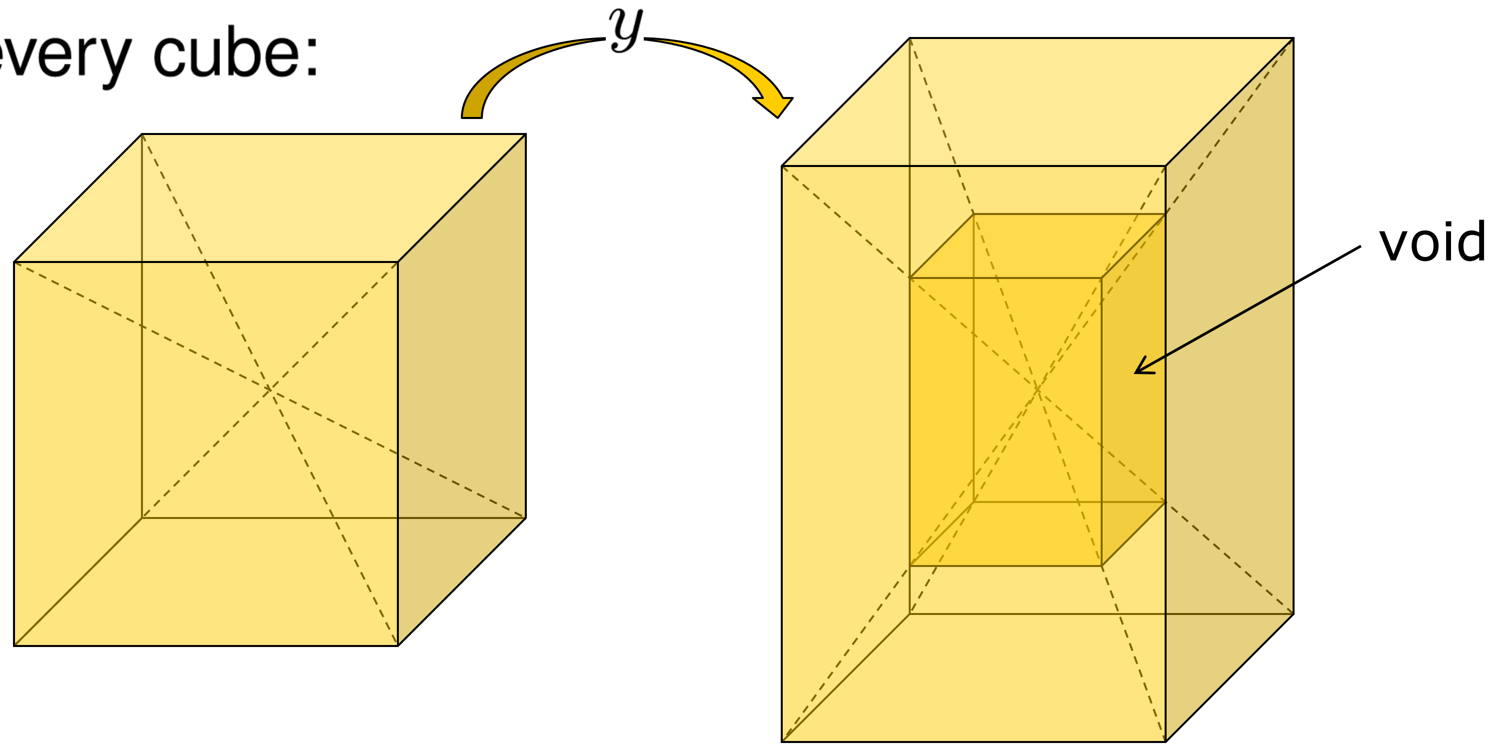


Heller, A., Science & Technology Review Magazine, LLNL, pp. 13-20, July/August, 2002



Sketch of proof – Upper bound

- In every cube:



- Calculate, estimate: $E \leq CL^2 \left(a^{1-p} \delta^p + \ell\delta / a \right)$
- Optimize: $a = (\ell\delta^{1-p})^{1/(2-p)} \Rightarrow E \leq C_U L^2 \ell^{\frac{1-p}{2-p}} \delta^{\frac{1}{2-p}}$

↑ void growth!

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Sketch of proof – Lower bound

- Lower bound: 1D arguments in x_3-direction
- For fixed (x_1, x_2): $f(x_3) \equiv |D_3 y(x_1, x_2, x_3)|$
- Then: $|D^2 y| \geq |D_3^2 y| \geq |D_3 |D_3 y|| = |Df|$

$$\bar{f} \equiv \frac{1}{2H} \int_{-H}^{H} f(x_3) dx_3 \geq \frac{1}{2H} \int_{-H}^{H} \frac{\partial u_3}{\partial x_3} dx_3 = 1 + \frac{\delta}{H}$$

- Define reduced energy density:

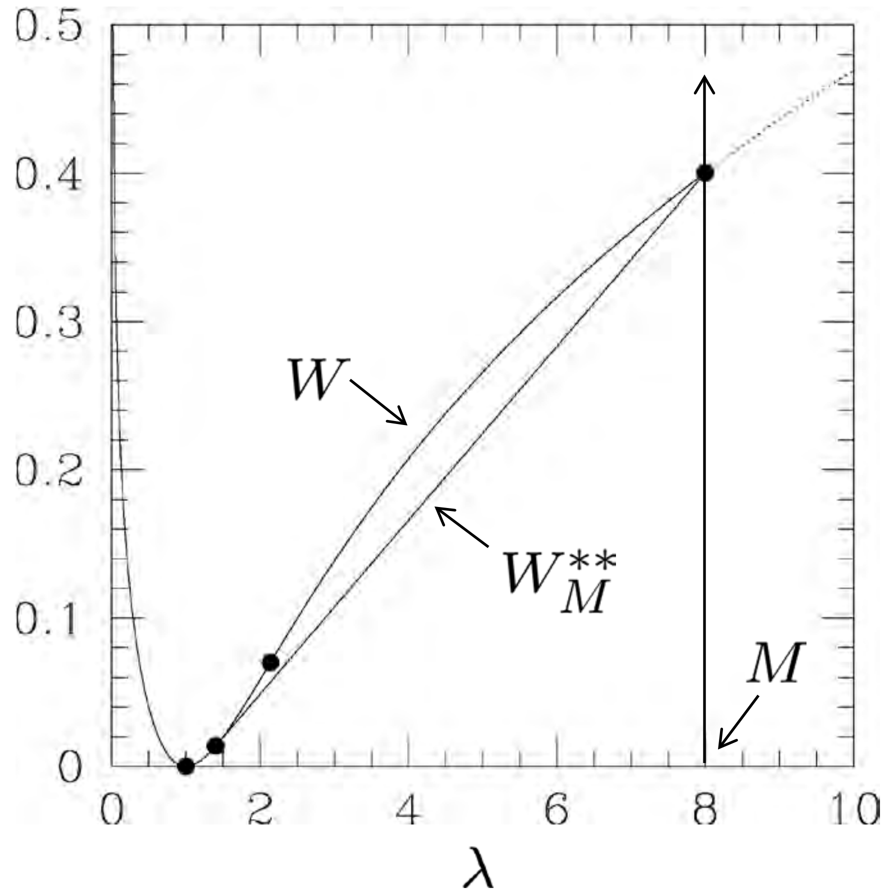
$$W(\lambda) = \min\{|F|^p - 3^{p/2},\ \det F = 1,\ |Fe_3| = \lambda\}$$

- Then: $\displaystyle\int_{-H}^{H} (|Dy|^p - 3^{p/2} + \ell |D^2 y|)\, dx_3 \geq$

$$\int_{-H}^{H} (W(f(x_3)) + \ell |Df(x_3)|)\, dx_3$$

Sketch of proof – Lower bound

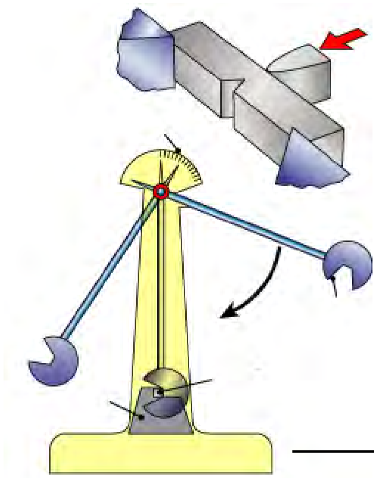
Graph of reduced energy density *W* and dual *W***



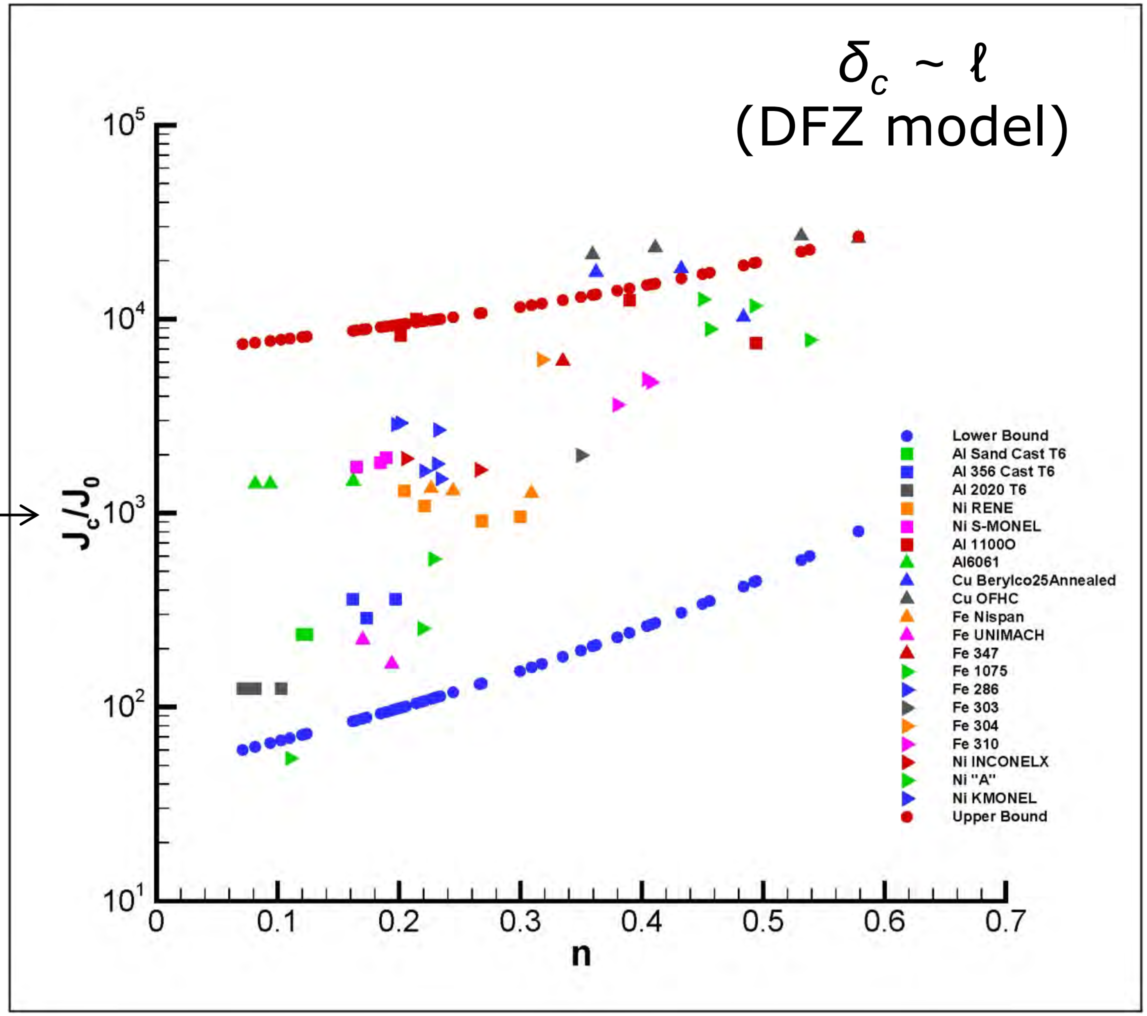
- Duality + Jensen + estimates: $E \geq C_L L^2 \ell^{\frac{1-p}{2-p}} \delta^{\frac{1}{2-p}}$

q. e. d.

Comparison with experiment



J_c from Charpy test data



L. Fokoua, S. Conti & MO,
J. Mech. Phys. Solids, **62** (2014) 295–311

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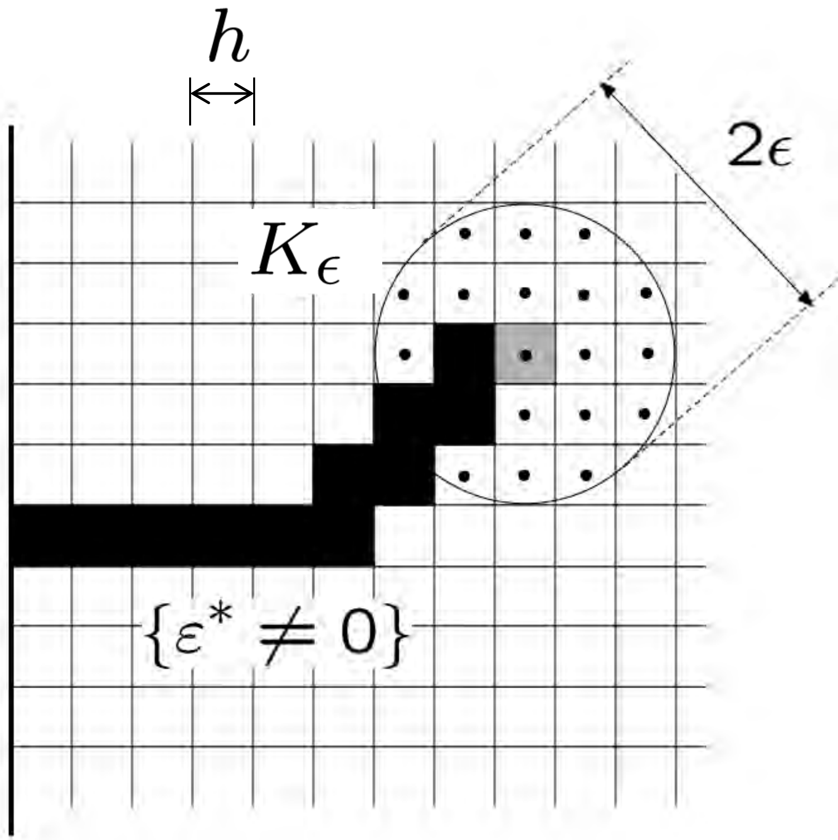
Numerical implementation
Material-point erosion



schematic of
ε-neighborhood
construction

- ε-neighborhood construction: Choose h << ε << L
- Erode material point if

$$\frac{h^2}{|K_\epsilon|}\int_{K_\epsilon} W(\nabla u)\,dx \geq J_c$$

- For linear elasticity, proof of Γ-convergence to Griffith fracture:
  – *Schmidt, B., Fraternali, F. & MO, SIAM J. Multiscale Model. Simul.*, **7**(3):1237-1366, 2009.

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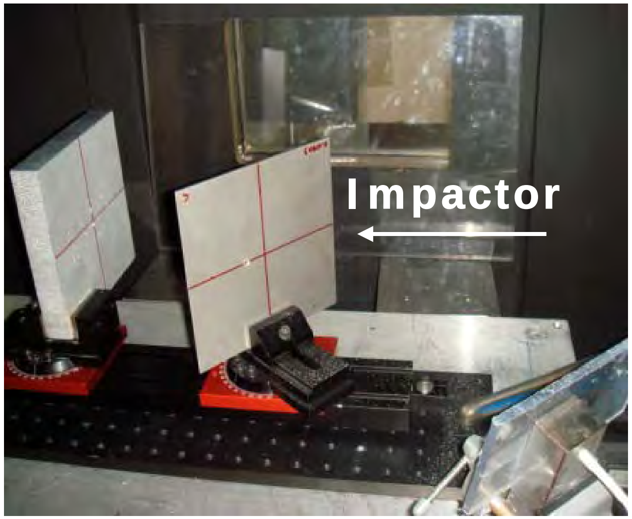
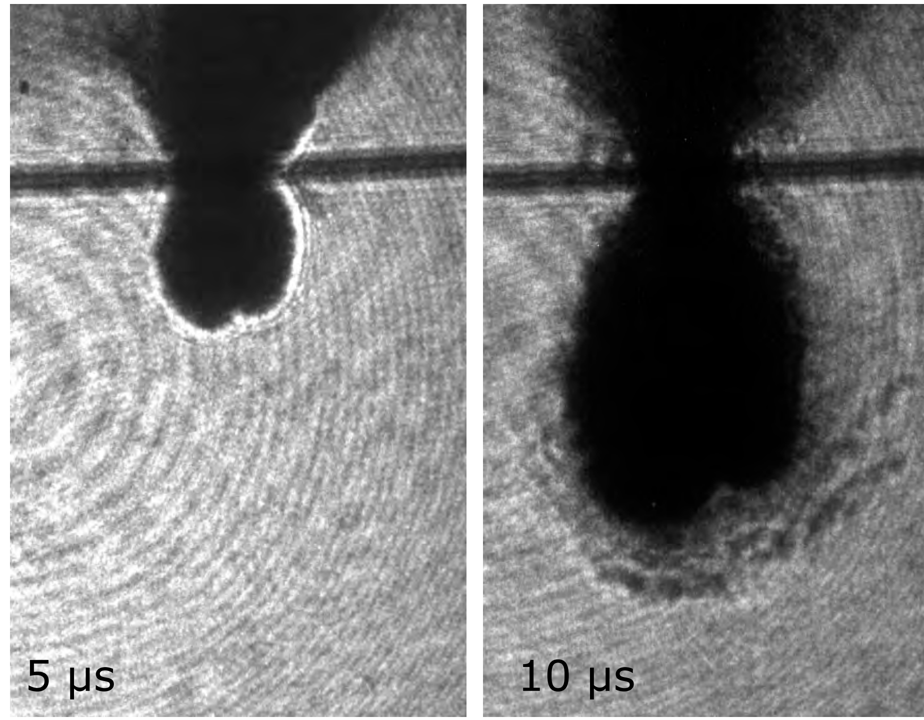
Application to hypervelocity impact

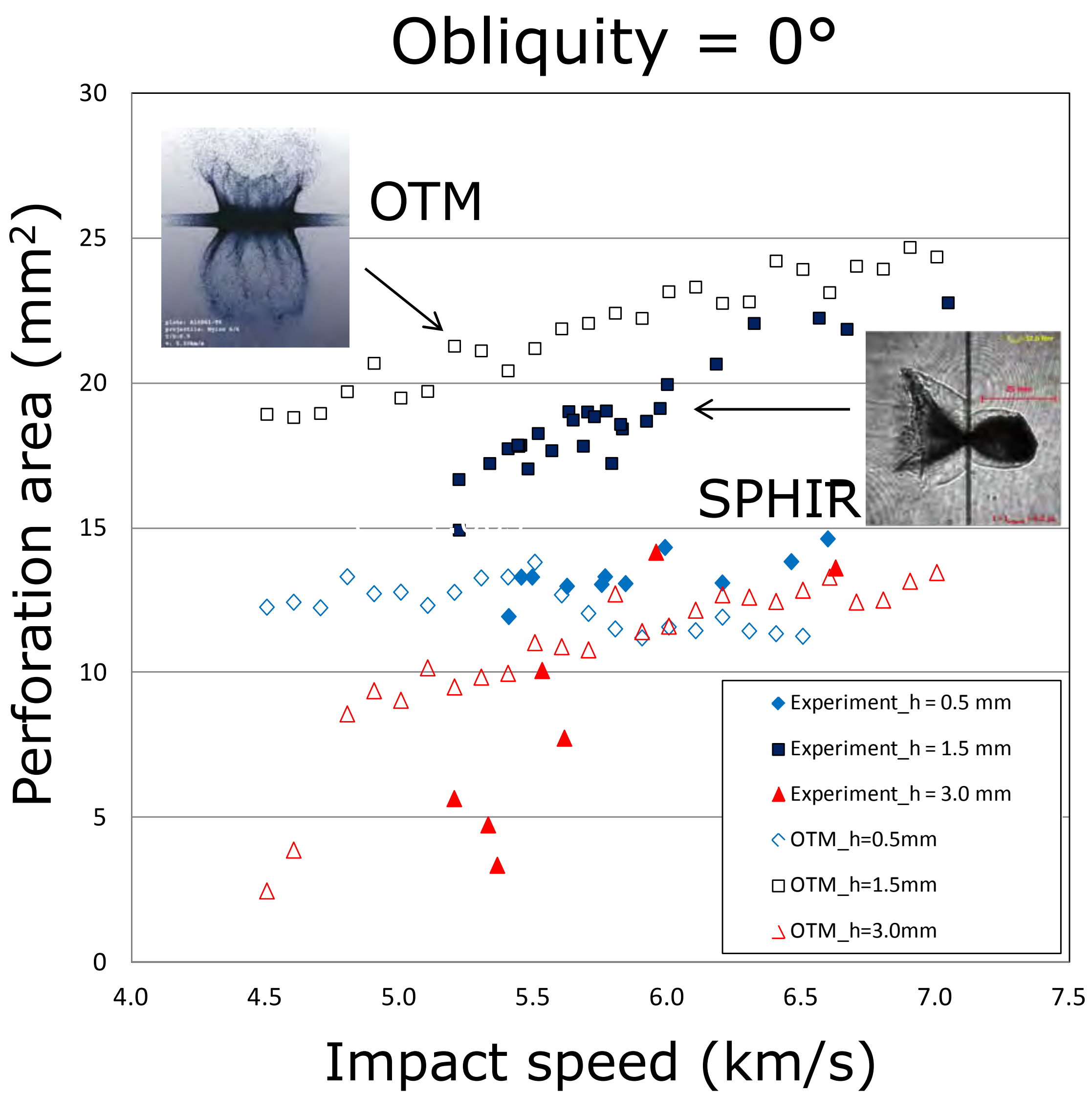


Caltech's hypervelocity
Impact facility



Hypervelocity impact (5.7 Km/s) of 0.96 mm thick aluminum plates by 5.5 mg nylon 6/6 cylinders (Caltech)

Application to hypervelocity impact



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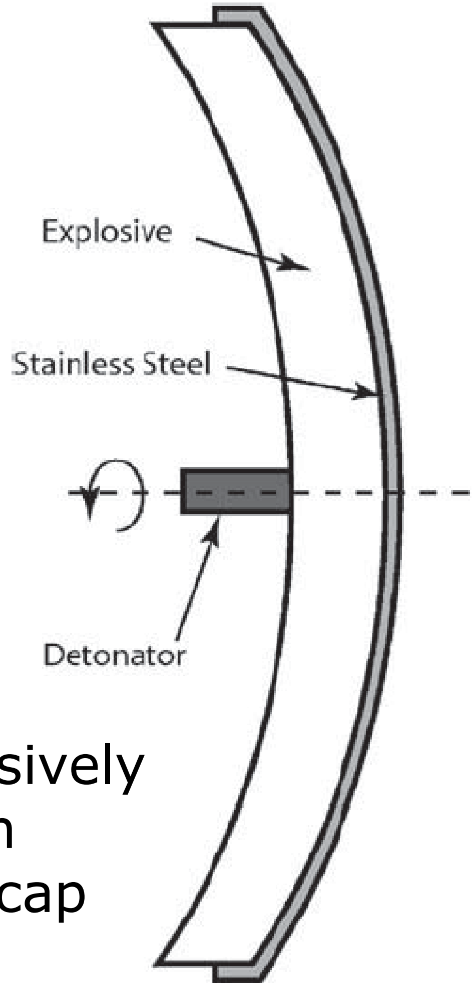
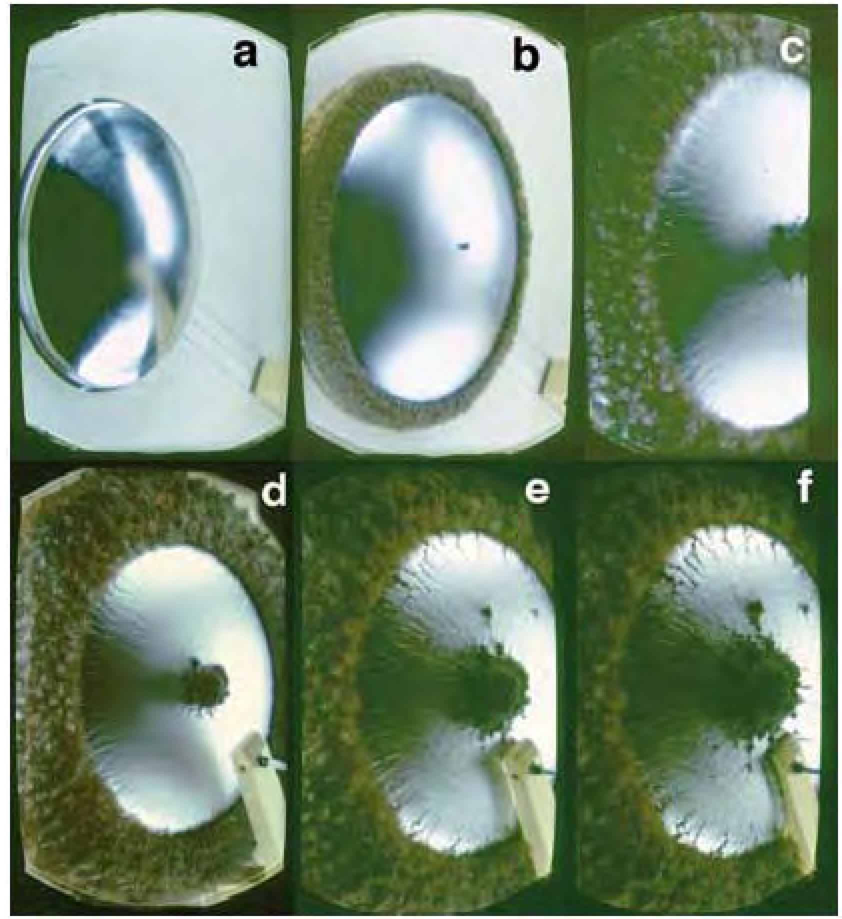
Application to explosively driven cap



Explosively driven steel cap

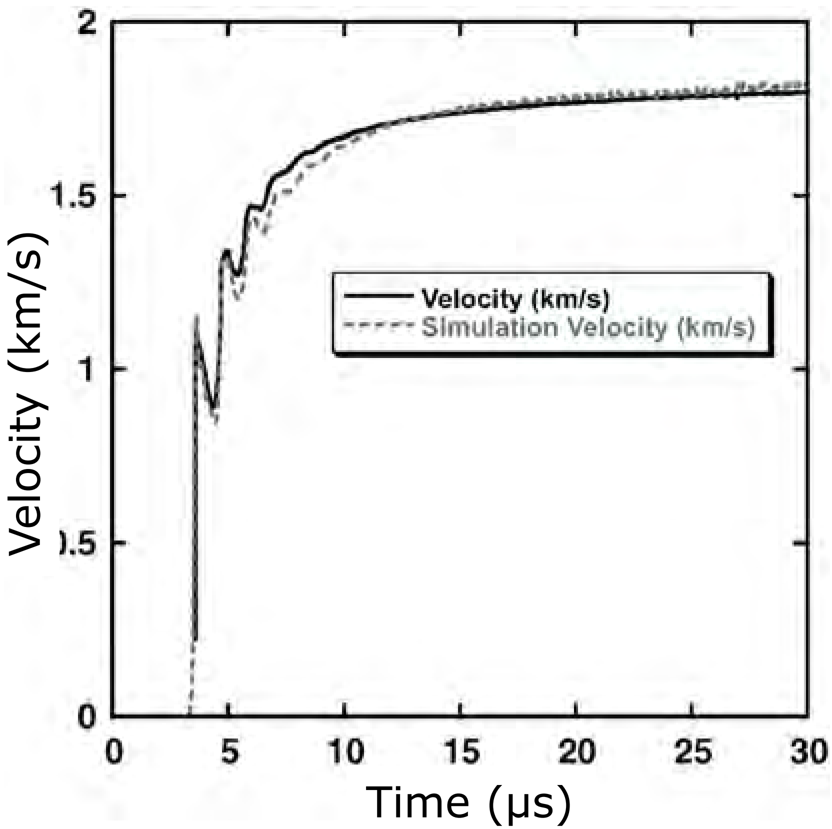


Optical framing camera records

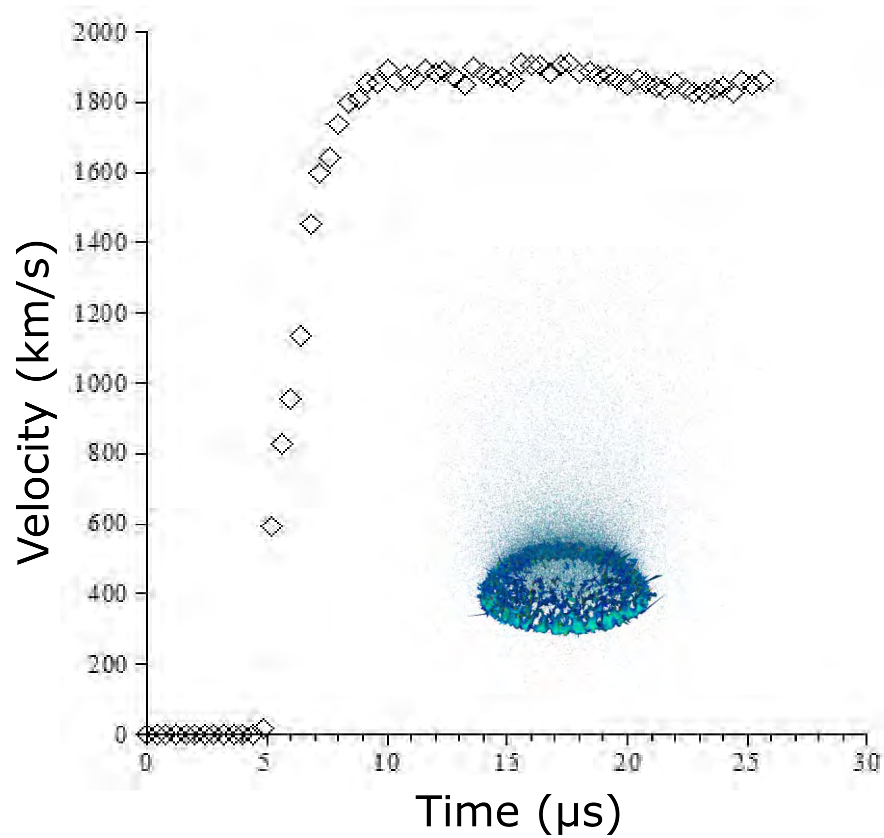
G.H. Campbell, G. C. Archbold, O. A. Hurricane and P. L. Miller, *JAP*, **101**:033540, 2007

Application to explosively driven cap

Experiment



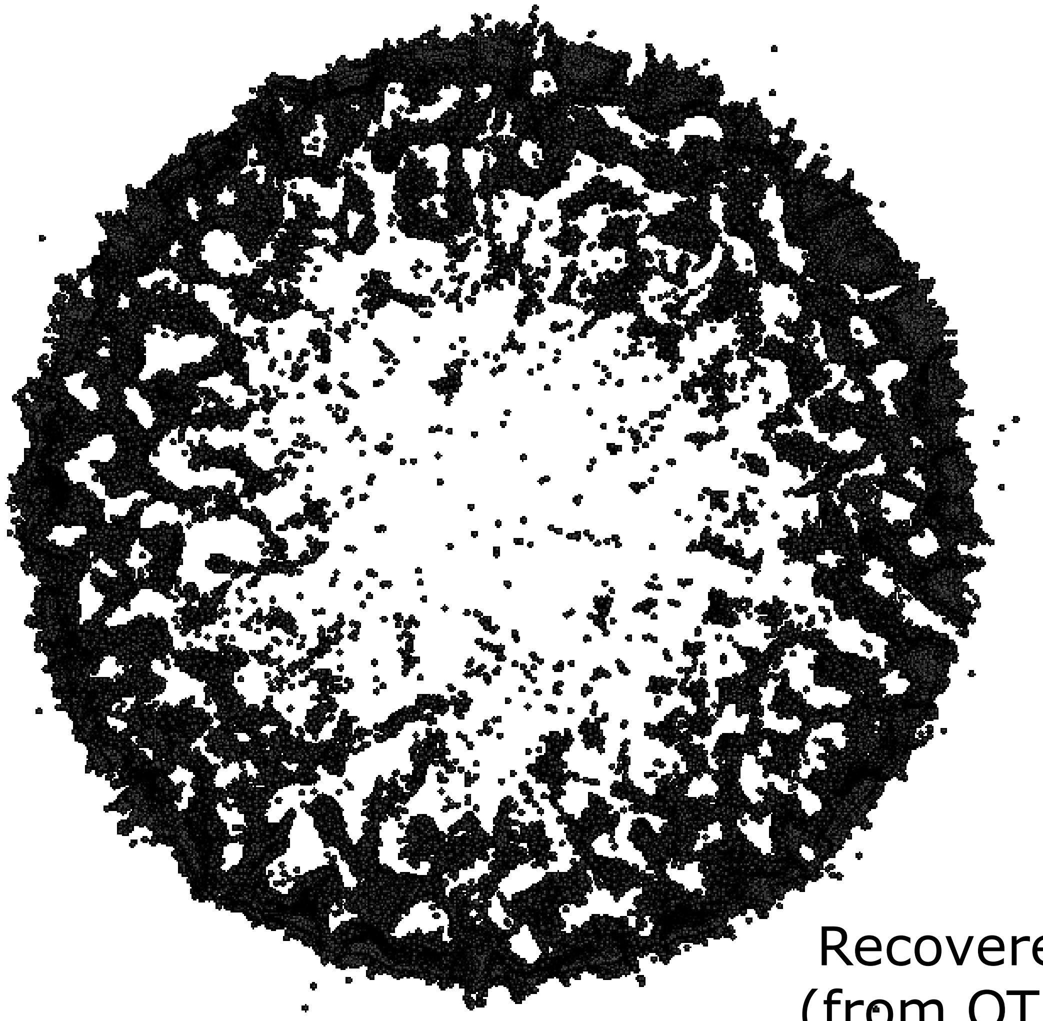
OTM simulation



Surface velocity for spot midway between pole and edge

G.H. Campbell, G. C. Archbold, O. A. Hurricane and P. L. Miller, *JAP*, **101**:033540, 2007

Application to explosively driven cap

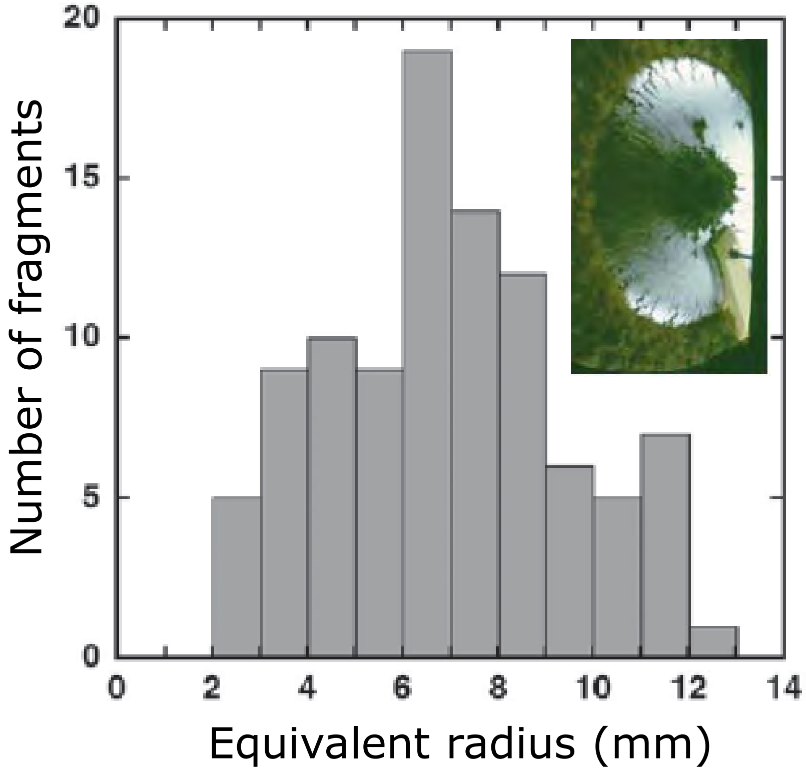


Recovered fragments
(from OTM simulation)

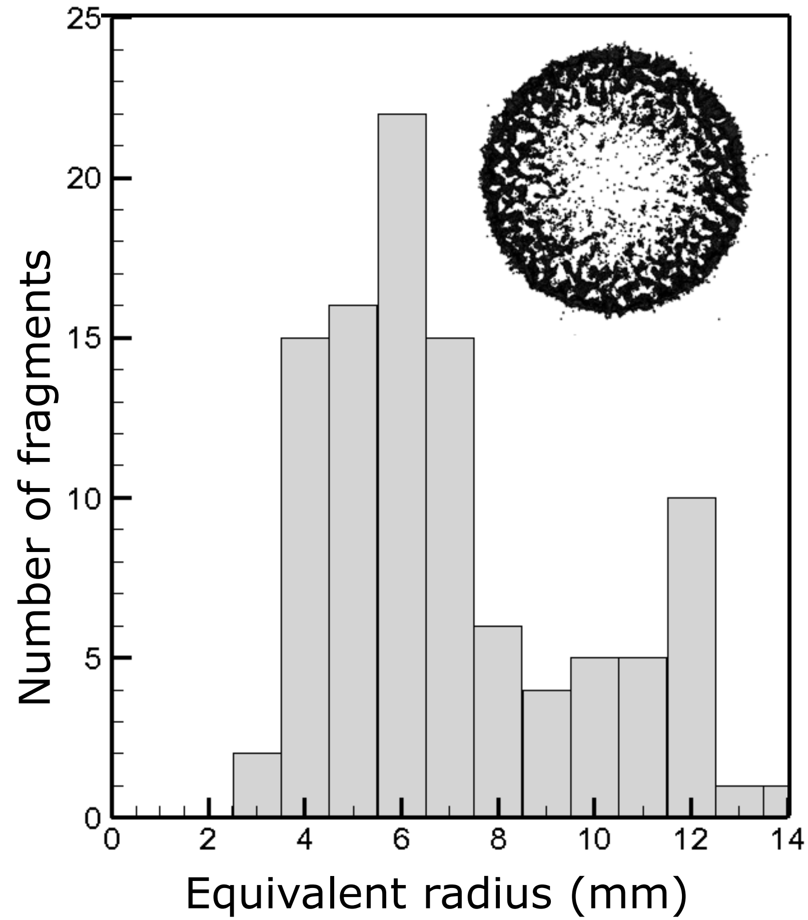
G.H. Campbell, G. C. Archbold, O. A. Hurricane and P. L. Miller, *JAP*, **101**:033540, 2007

Application to explosively driven cap

Experiment



OTM simulation



Histograms of equivalent fragment radii

G.H. Campbell, G. C. Archbold, O. A. Hurricane and P. L. Miller, *JAP*, **101**:033540, 2007

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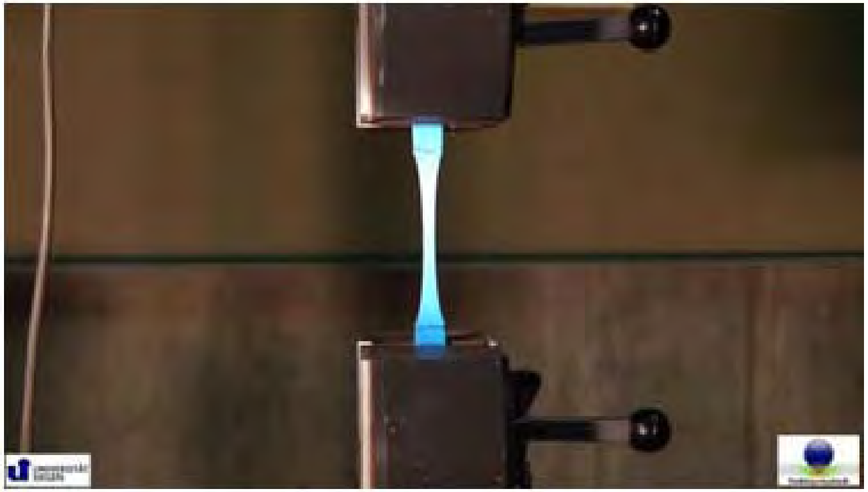
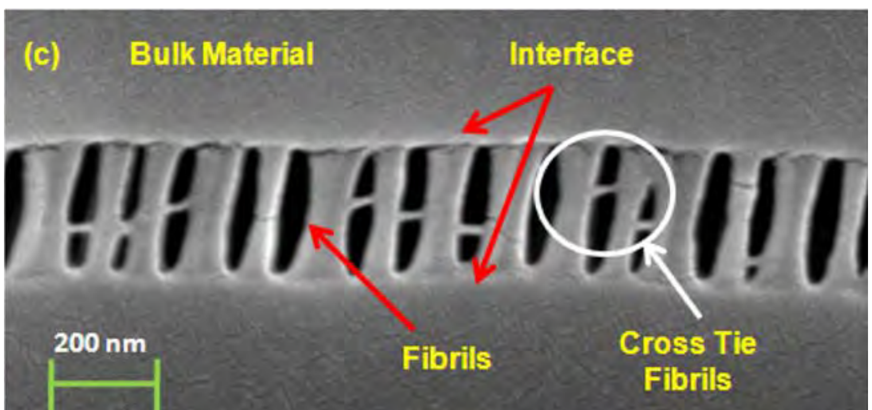
Fracture of polymers



T. Reppel, T. Dally, T. and K. Weinberg, Technische Mechanik, 33 (2012) 19-33.



Crazing in 800 nm polystyrene thin film (C. K. Desai *et al*., 2011)

- Polymers undergo entropic elasticity and damage due to chain stretching and failure
- Polymers fracture by means of the crazing mechanism consisting of fibril nucleation, stretching and failure
- The free energy density of polymers saturates in tension once the majority of chains are failed: p=0!
- Crazing mechanism is incompatible with strain-gradient elasticity…

Fracture of polymers

- Suppose: For $K_U > 0$, *intrinsic length* $\ell > 0$,

$$E(y) \leq K_U \left(\int_\Omega (|Dy|^p - 3^{p/2})\, dx + \ell \int_\Omega |D^2 y|\, dx \right)$$

- If $E(y) < +\infty \Rightarrow y$ continuous on a.e. plane!
- Crazing is precluded by the continuity of y!
- Instead suppose: For $\sigma \in (0,1)$,

$$E(y) \leq K_U \left(\int_\Omega (|Dy|^p - 3^{p/2})\, dx + \ell^\sigma |y|_{W^{1+\sigma,1}(\Omega)} \right)$$

Theorem [Conti, Heyden & MO]. For ℓ sufficiently small,

$p = 0,\ \sigma \in (0,1),\ 0 < C_L < C_U,$

$$C_L L^2 \ell^{\frac{\sigma}{1+\sigma}} \delta^{\frac{1}{1+\sigma}} \leq \inf E \leq C_U L^2 \ell^{\frac{\sigma}{1+\sigma}} \delta^{\frac{1}{1+\sigma}}$$

Optimal scaling – Crazing

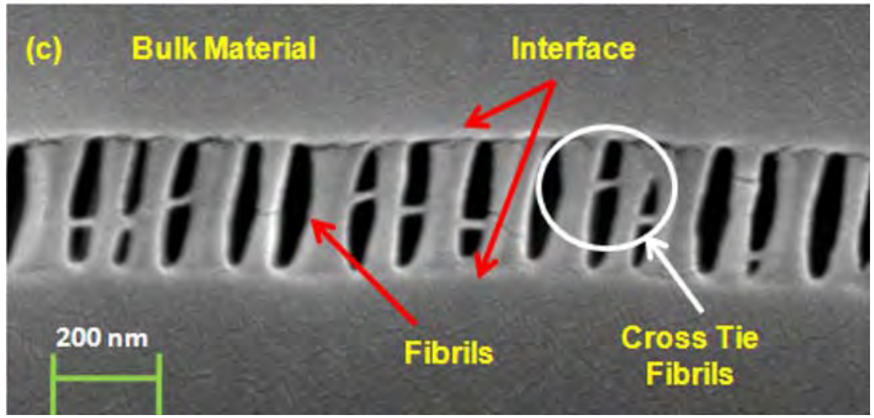
- Optimal (matching) upper and lower bounds:

$$C_L L^2 \ell^{\frac{\sigma}{1+\sigma}} \delta^{\frac{1}{1+\sigma}} \leq \inf E \leq C_U L^2 \ell^{\frac{\sigma}{1+\sigma}} \delta^{\frac{1}{1+\sigma}}$$

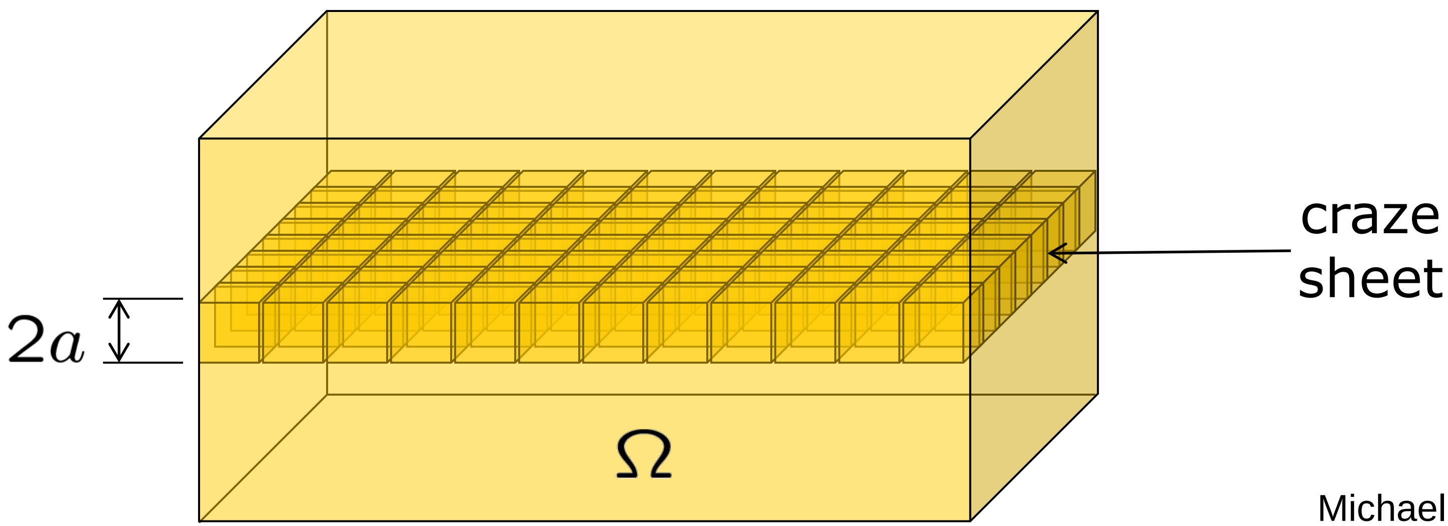
- *Fractional* strain-gradient elasticity supplies bounded energies for crazing mechanism
- Energy scales with *area* (L^2): *Fracture scaling*!
- Energy scales with power of opening displacement (δ): *Cohesive behavior*!
- Lower bound degenerates to 0 when the intrinsic length (ℓ) decreases to zero…
- Bounds on specific fracture energy:

$$C_L \ell^{\frac{\sigma}{1+\sigma}} \delta_c^{\frac{1}{1+\sigma}} \leq G_c \leq C_U \ell^{\frac{\sigma}{1+\sigma}} \delta_c^{\frac{1}{1+\sigma}}$$

Sketch of proof – Upper bound

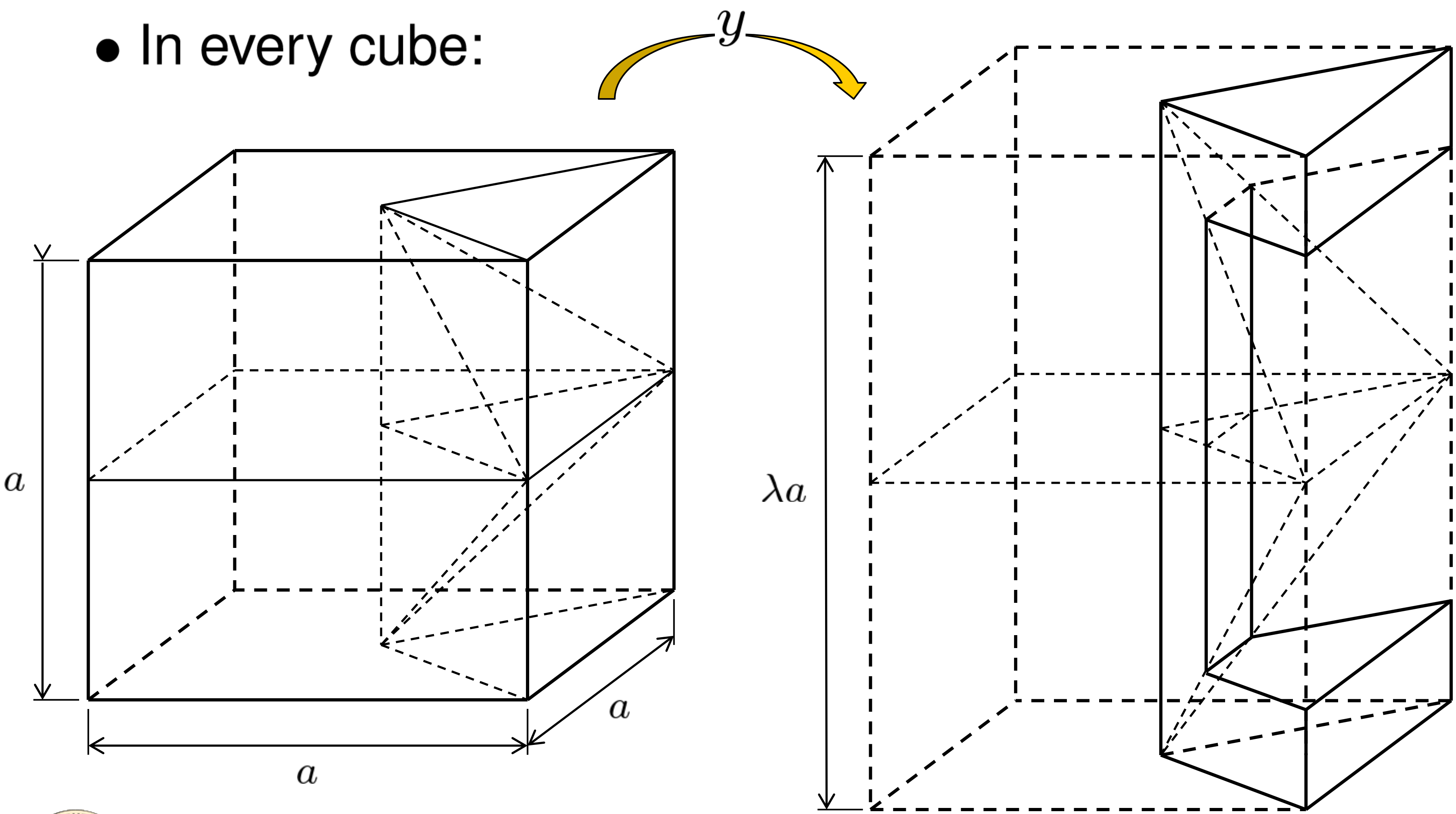


Crazing in 800 nm polystyrene
thin film (C. K. Desai *et al*., 2011)



Sketch of proof – Upper bound

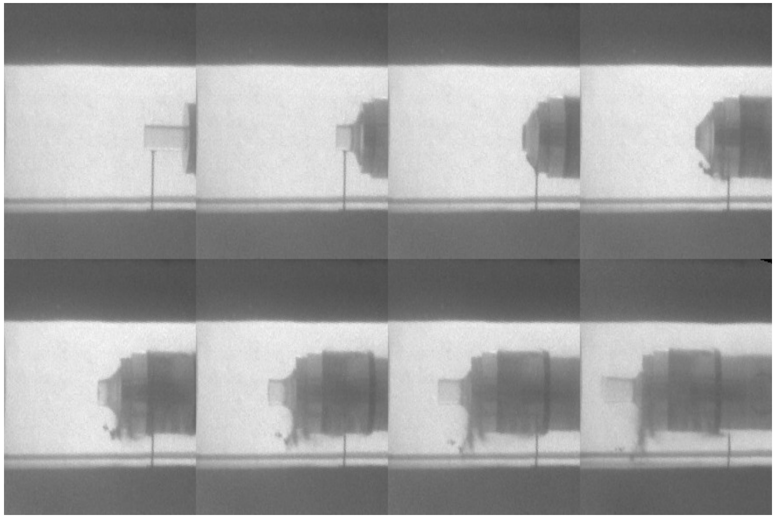
- In every cube:



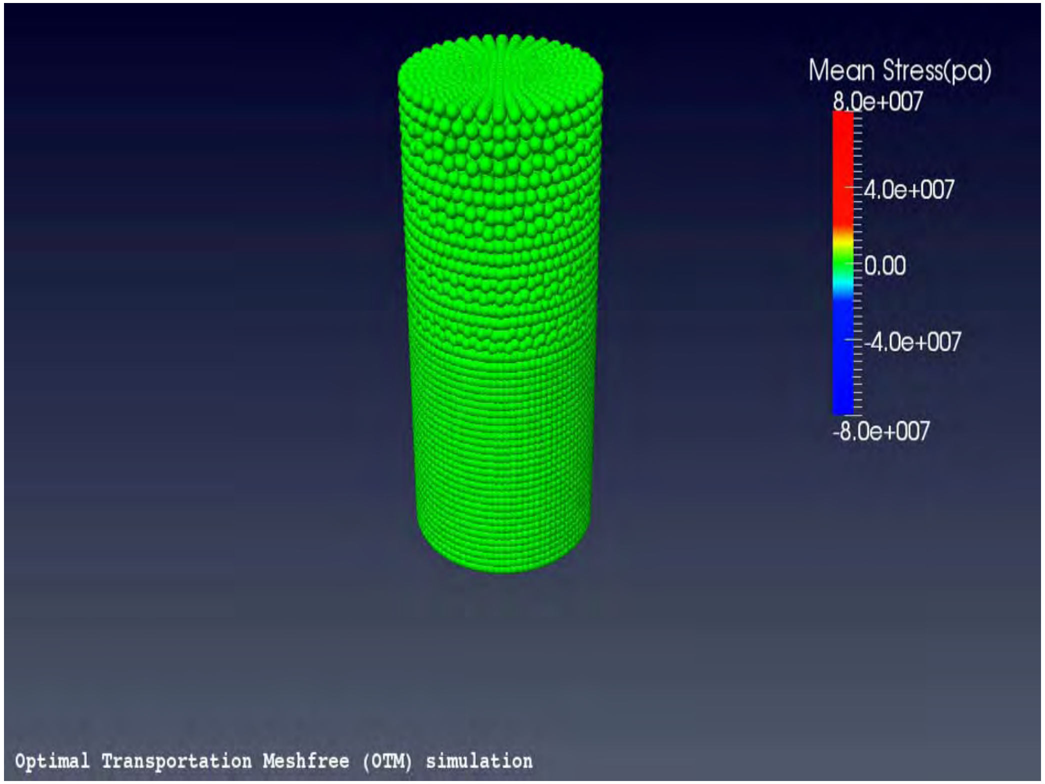
Outline

- Background/phenomenology of ductile fracture
- Metals: mathematical formulation
- Optimal scaling laws
- Numerical approximation
- Applications: Hypervelocity impact and explosively-driven caps
- Extension to polymers
- Application: Taylor anvil tests on polyurea

Taylor-anvil tests on polyurea



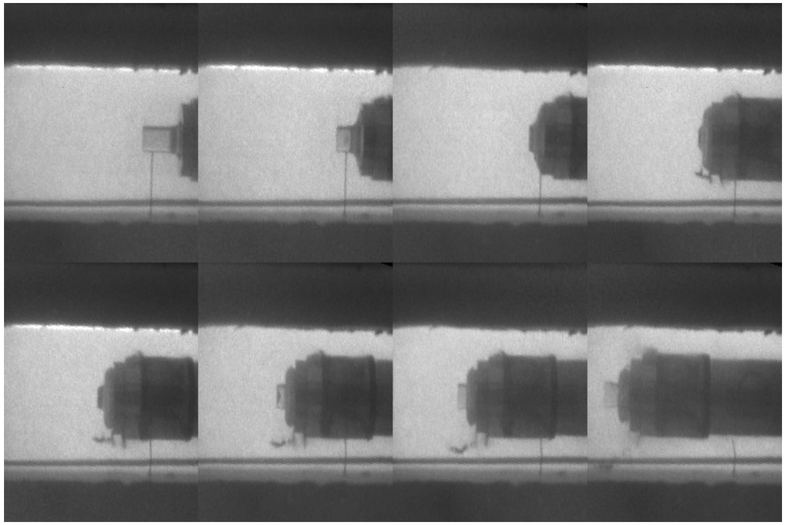
Shot #854:
R0 = 6.3075 mm,
L0 = 27.6897 mm,
v = 332 m/s



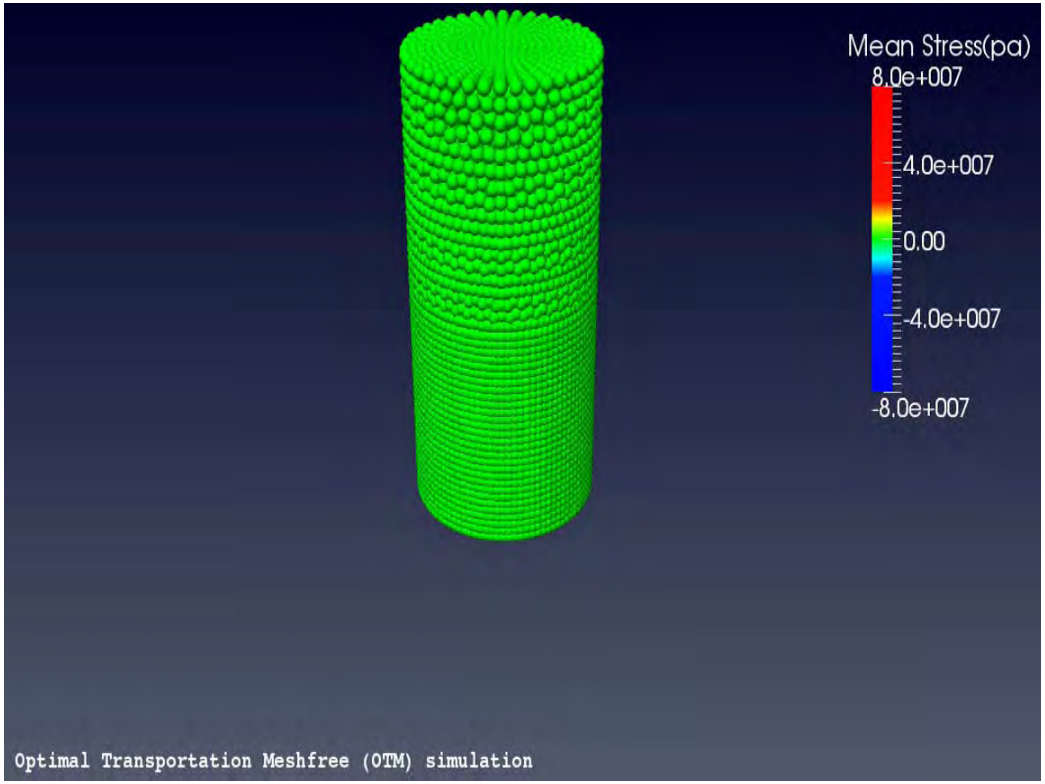
Experiments conducted by W. Mock, Jr. and J. Drotar,
at the Naval Surface Warfare Center (Dahlgren Division)
Research Gas Gun Facility, Dahlgren, VA 22448-5100, USA

S. Heyden *el al.*, *JMPS* (submitted)

Experiments and simulations



Shot #861:
R0 = 6.3039 mm,
L0 = 27.1698 mm,
v = 424 m/s



Experiments conducted by W. Mock, Jr. and J. Drotar,
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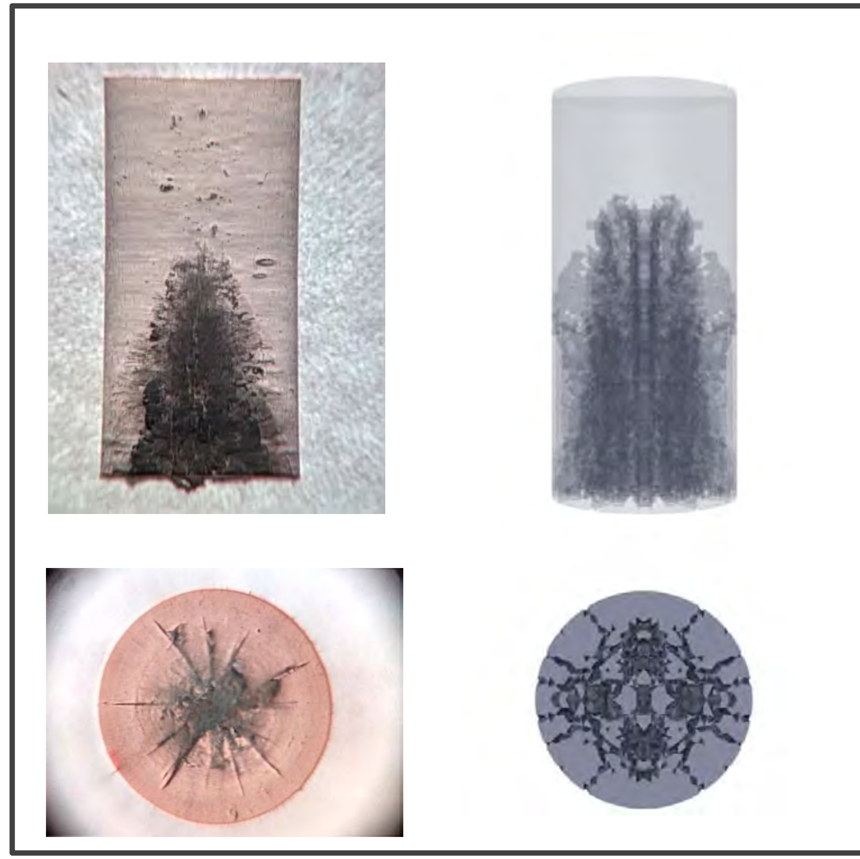
Research Gas Gun Facility, Dahlgren, VA 22448-5100, USA

S. Heyden *el al.*, *JMPS* (submitted)

Michael Ortiz
UCSD 01/14

Taylor-anvil tests on polyurea

Shot #854



Shot #861

Comparison of damage and fracture patterns
in recovered specimens and simulations

S. Heyden *el al.*, *JMPS* (submitted)

Michael Ortiz
UCSD 01/14

Concluding remarks

- Ductile fracture can indeed be understood as the result of the competition between sublinear growth and (possibly fractional) strain-gradient effects
- Optimal scaling laws are indicative of a well-defined specific fracture energy, cohesive behavior, and provide a (multiscale) link between macroscopic fracture properties and micromechanics (intrinsic micromechanical length scale, void-sheet and crazing mechanisms…)
- Ductile fracture can be efficiently implemented through material-point erosion schemes

Concluding remarks

Thanks!

(This program has been made possible by
generous contributions from family, mentors, alumni,
institutions, funding agencies and viewers like you!)