

# Non-Equilibrium Thermodynamics and Multiscale Modeling of Dynamic Void Growth in Crystals

M. Ortiz, G. Venturini (Caltech)

M.P. Ariza, M. Ponga (U.P. Seville)

I. Romero (U.P. Madrid)

Intl. Symp. on Current Problems in Solid  
Mechanics in honor of

**Professor Rodney J. Clifton**

June 24th – 29th, 2012, Symi, Greece



Michael Ortiz  
RJC 06/12

# To Professor Rodney J. Clifton



Scientist, teacher, mentor



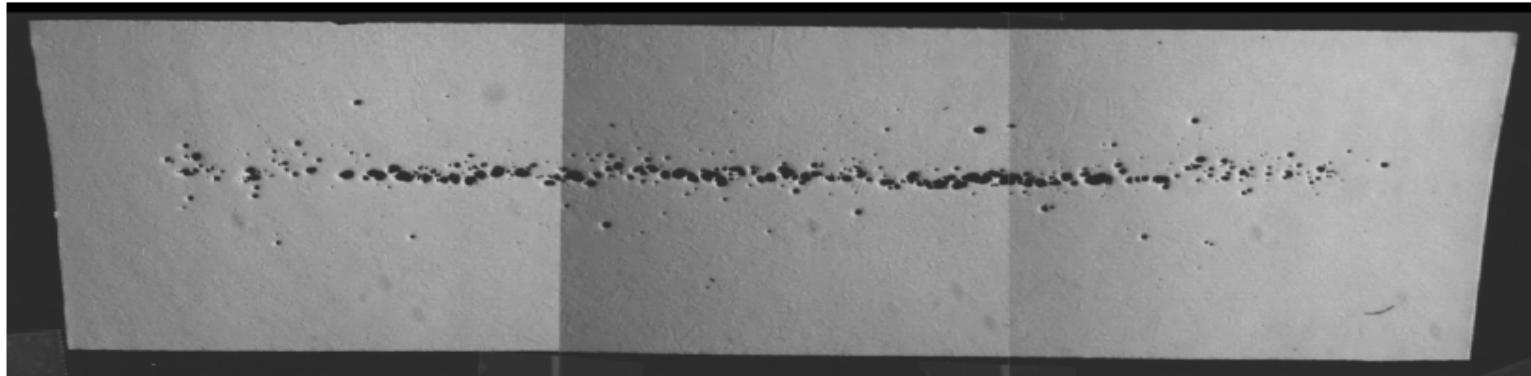
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# Atomistic-to-continuum - Outline

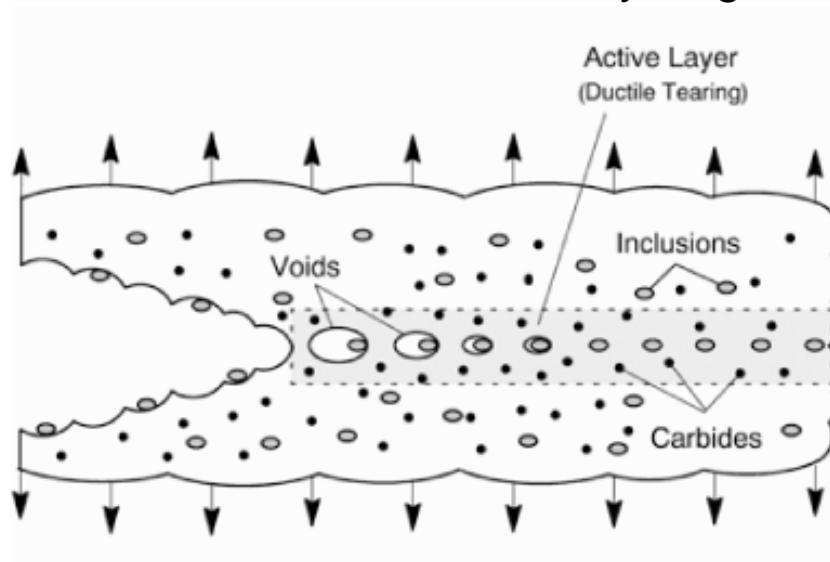
- The essential difficulty: Multiple scales,
  - *Atomic level rate-limiting processes: Thermal vibrations, lattice defects, transport (thermal, mass...)*
  - *Macroscopic processes of interest: Ductile fracture, GB embrittlement, irradiation damage, aging...*
- Time-scale gap: From molecular dynamics (MD) (femtosecond) to macroscopic (seconds-years)
- Spatial-scale gap: From lattice defects (Angstroms) to macroscopic (mm-m)
- Application to nanovoid plastic cavitation in Cu at intermediate strain rates (e.g.,  $p_c(a,T,d\varepsilon/dt)$ ):
  - *Dislocation emission mediates cavitation*
  - *MD limited to extreme strain rates ( $10^8$ - $10^{12}$  1/s)*
  - **Question:** *Intermediate strain rates? (<  $10^8$  1/s)*



# Motivation – Ductile fracture in metals



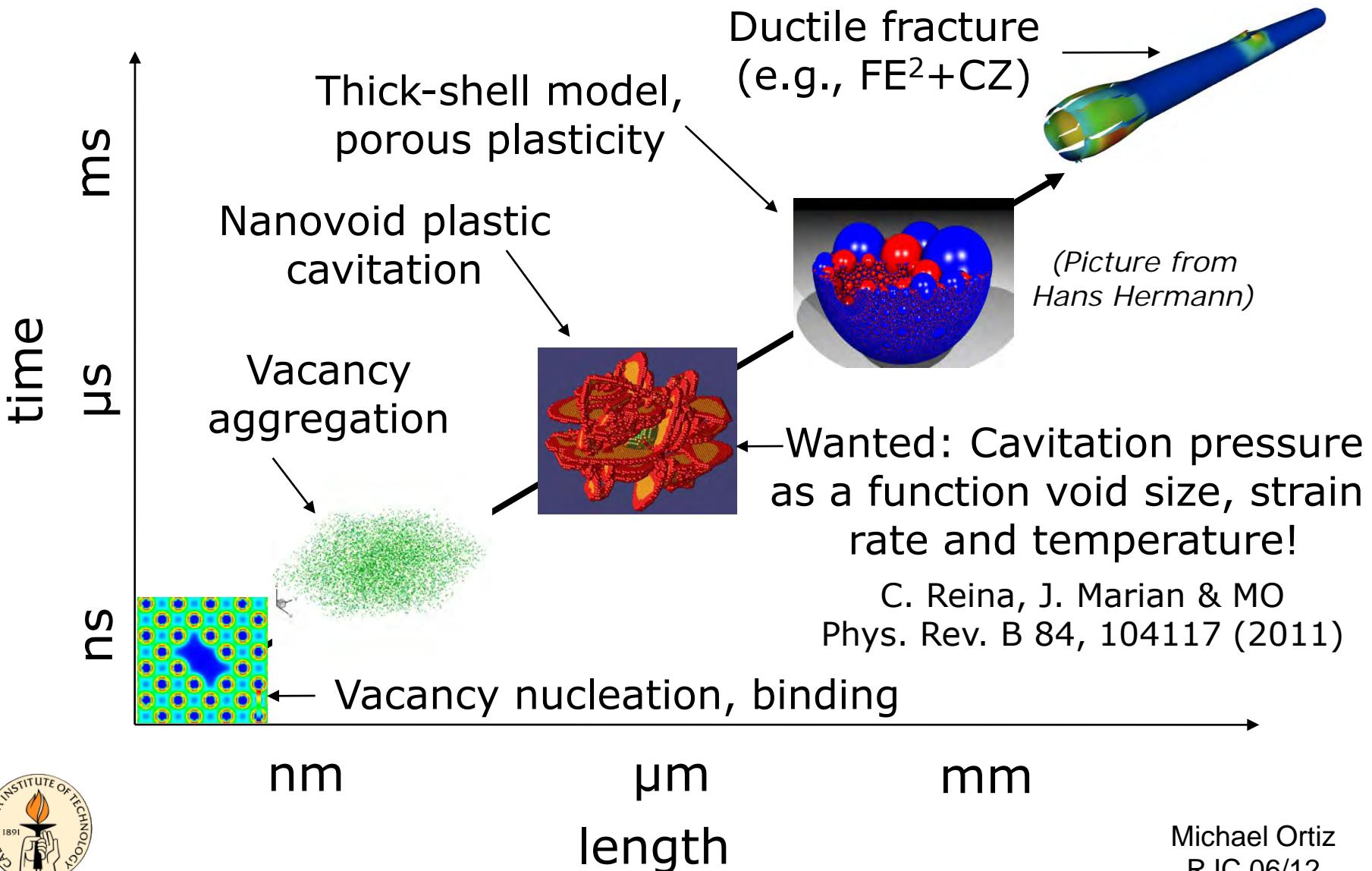
R. Becker "How Metals Fail", *Science and Technology Review*, LLNL,  
July/August 2002



C. Ruggieri, J. of the Braz. Soc. of Mech. Sci. & Eng., Vol. XXVI, No. 2 (2004) 190-198.

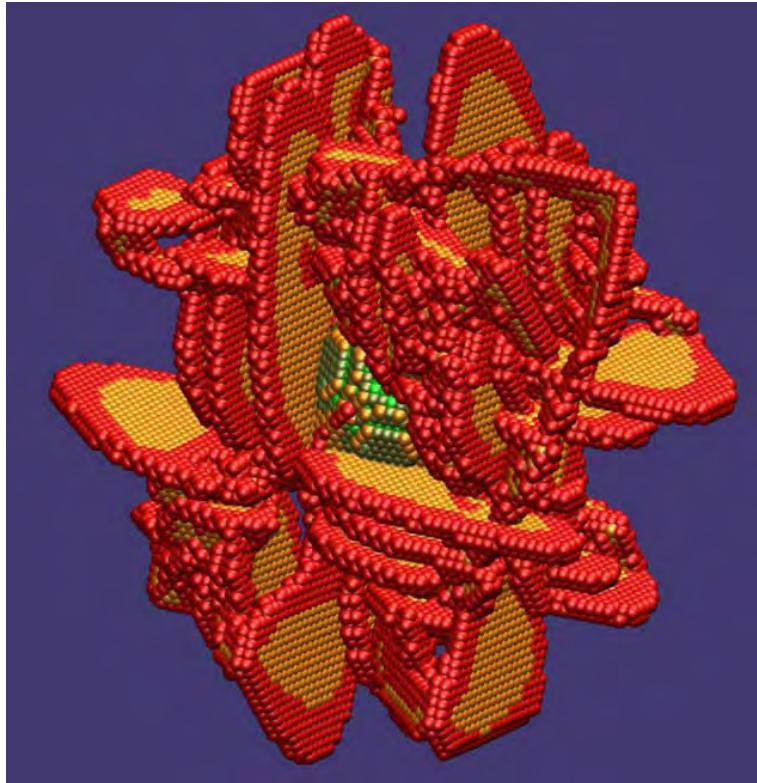


# Motivation – Ductile fracture in metals

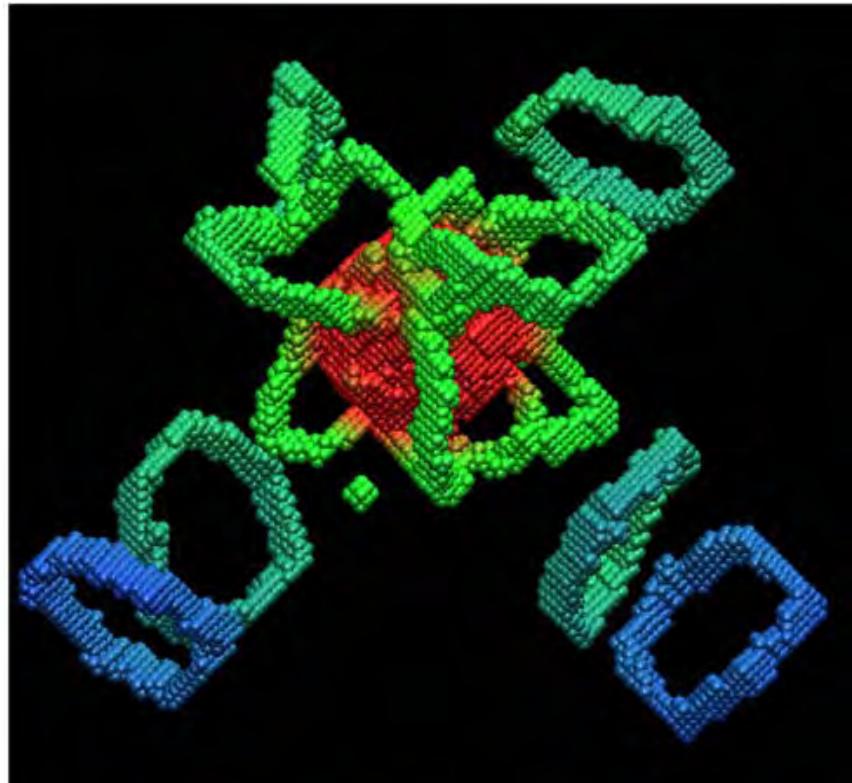


# Nanovoid cavitation – MD work

$T=0, d\varepsilon/dt < 10^8 \text{ 1/s}$



$T>0, d\varepsilon/dt < 10^8-10^{12} \text{ 1/s}$

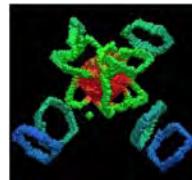


Marian J., Knap J. and MO, "Nanovoid deformation in aluminum under simple shear", *Acta Materialia*, **53**:2893, 2005

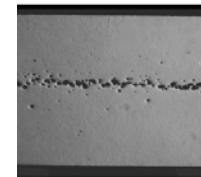


Tang, Y., Bringa, E.M., Remington, B.A., and Meyers, M.A., "Growth and collapse of nanovoids in Ta monocrystals", *Acta Materialia*, **59**:1354, 2011 Michael Ortiz  
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# Atomistic-to-continuum



Tang, Y. et al.,  
*Acta Mater.*,  
**59**:1354, 2011.



R. Becker,  
*S&T. Rev.*, (LLNL)  
July/Aug 2002

	Atomistic	Coarse-grained
Configuration space	Phase space (micro)	<ul style="list-style-type: none"><li>Phase space (meso)</li><li>Temperature</li><li>Concentrations</li></ul>
Governing equations	$\Sigma F = ma$	<ul style="list-style-type: none"><li><math>\Sigma F = ma</math> (averaged)</li><li>Diffusive transport</li></ul>
Spatial resolution	Atomic lattice	<ul style="list-style-type: none"><li>Lattice defects</li><li>Temperature grads.</li><li>Concentration grads.</li></ul>
Temporal resolution	<ul style="list-style-type: none"><li>Thermal vibrations</li><li>Transition states</li></ul>	<ul style="list-style-type: none"><li>Elastic waves (meso)</li><li>Diffusional transients</li></ul>
Time-scale bridging	Non-equilibrium statistical mechanics	
Spatial-scale bridging	Quasicontinuum method	



# Equilibrium SM – Max-ent formulation

- Gran-canonical ensemble,  $N$  atoms,  $M$  species:
  - State:  $(\{\mathbf{q}\}, \{\mathbf{p}\}, \{\mathbf{n}\}) \in \mathbb{R}^{3N} \times \mathbb{R}^{3N} \times \mathcal{O}_{NM}$
  - Atomic positions:  $\{\mathbf{q}\} = \{q_1, \dots, q_N\}$
  - Atomic momenta:  $\{\mathbf{p}\} = \{p_1, \dots, p_N\}$
  - Occupancy:  $n_{ik} = \begin{cases} 1, & \text{site } i \text{ occupied by species } k, \\ 0, & \text{otherwise.} \end{cases}$
- Ensemble average of observable  $A(\{\mathbf{q}\}, \{\mathbf{p}\}, \{\mathbf{n}\})$ :

$$\langle A \rangle = \sum_{\{\mathbf{n}\} \in \mathcal{O}_{NM}} \int A(\{\mathbf{q}\}, \{\mathbf{p}\}, \{\mathbf{n}\}) p(\{\mathbf{q}\}, \{\mathbf{p}\}, \{\mathbf{n}\}) dq dp$$



•  $p(\{\mathbf{q}\}, \{\mathbf{p}\}, \{\mathbf{n}\}) \equiv$  Gran-canonical pdf

# Equilibrium SM – Max-ent formulation

- Information-theoretical entropy:  $\mathcal{S}[p] = -k_B \langle \log p \rangle$

- Principle of maximum entropy:  $\mathcal{S}[p] \rightarrow \max!$

subject to:  $\langle H \rangle = E, \quad \left\langle \sum_{i=1}^N n_{ik} \right\rangle = N_k$

- Lagrangian: reciprocal temperature

$$\mathcal{L}[p, \beta, \Gamma] = \mathcal{S}[p] - k_B \beta \downarrow \langle H \rangle - k_B \sum_{k=1}^M \Gamma_k \uparrow \left\langle \sum_{i=1}^N n_{ik} \right\rangle.$$

- Gran-canonical distribution:

$$p = \frac{1}{\Xi} e^{-\beta H - \Gamma^T \sum_{i=1}^N \mathbf{n}_i},$$

$$\Xi = \sum_{\{\mathbf{n}\} \in \mathcal{O}_{NM}} \int e^{-\beta H - \Gamma^T \sum_{i=1}^N \mathbf{n}_i} dq dp$$

chemical potentials



# Equilibrium SM – Macroscopic equilibrium

- Macroscopic variables  $\equiv \mathbf{X}$  (e.g., volume, strain...)
- Parametrized Hamiltonians:  $H(\{q\}, \{p\}, \{n\}, \mathbf{X})$
- Gran-canonical thermodynamic potentials:
  - Internal energy:  $U(S, \Gamma, \mathbf{X}) = -\frac{\partial \log \Xi}{\partial \beta}$
  - Free energy:  $F(\beta, \Gamma, \mathbf{X}) = -\frac{1}{\beta} \log \Xi$
  - Free entropy:  $\Phi(\beta, \Gamma, \mathbf{X}) = k_B \log \Xi$
- Equivalent statements of macroscopic equilibrium:
  - $U(S, \Gamma, \cdot) \rightarrow \min!$  (internal energy principle)
  - $F(\beta, \Gamma, \cdot) \rightarrow \min!$  (free energy principle)
  - $\Phi(\beta, \Gamma, \cdot) \rightarrow \max!$  (free entropy principle)



# Equilibrium SM – Meanfield theory

- Space of trial Hamiltonians:  $\mathcal{H}_0$
- Free-entropy inequality: For all  $H_0 \in \mathcal{H}_0$ ,
  - i)  $\Phi \geq \Phi_0 - k_B\beta\langle H - H_0 \rangle_0 \equiv -\mathcal{F}[H_0, \beta, \Gamma]$
  - ii)  $\Phi = -\mathcal{F}[H_0, \beta, \Gamma] \Leftrightarrow H_0 = H$
- Best approximation:  $\mathcal{F}[\cdot, \beta, \Gamma] \rightarrow \min!$  over  $\mathcal{H}_0$
- Example:  $\mathcal{H}_0 \equiv$  uncoupled harmonic oscillators,

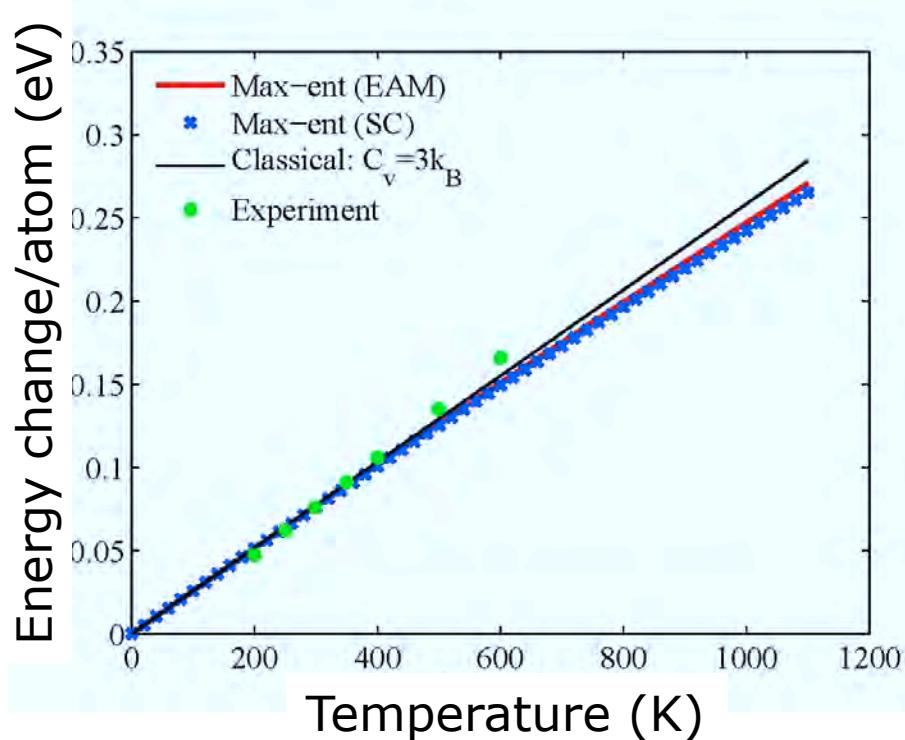
$$H_0 = \sum_{i=1}^N \left( \frac{1}{2m(n_i)} |\mathbf{p}_i - \bar{\mathbf{p}}_i|^2 + \frac{m(n_i)\omega_i^2}{2} |\mathbf{q}_i - \bar{\mathbf{q}}_i|^2 \right)$$

- Minimize  $\mathcal{F}(\{\pi\}, \beta, \Gamma)$  w.r.t.  $\{\pi\} \equiv (\{\bar{\mathbf{q}}\}, \{\bar{\mathbf{p}}\}, \{\omega\})$
- Compute  $\langle \cdot \rangle_0$  using Gaussian quadrature

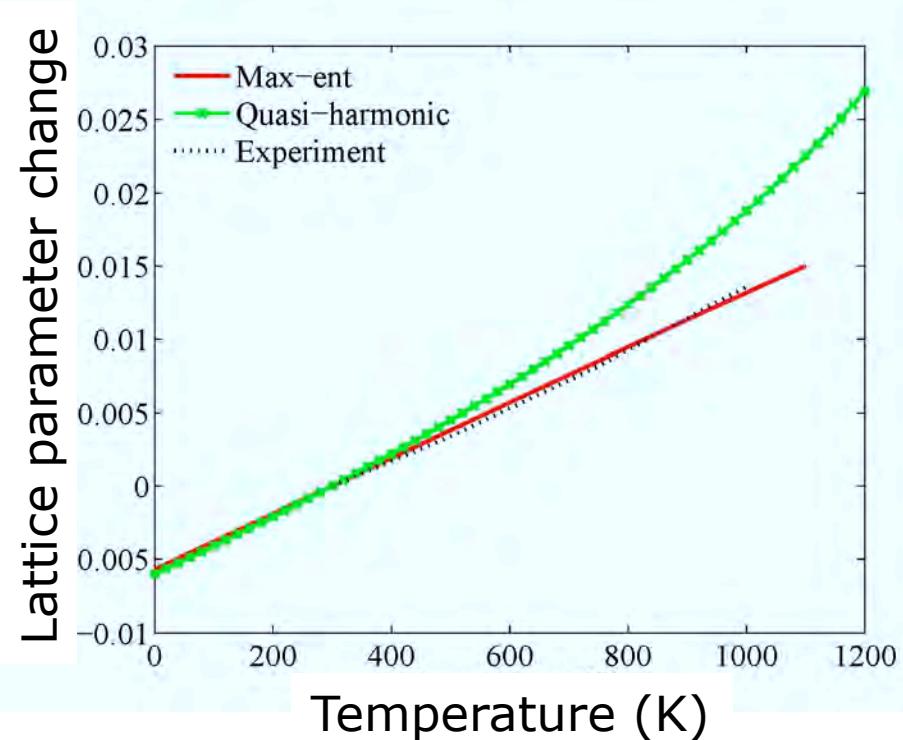


# Equilibrium SM – Meanfield validation

Specific heat



Thermal expansion



Johnson, R.A., *Phys. Rev. B*, **37**:3924, 1988

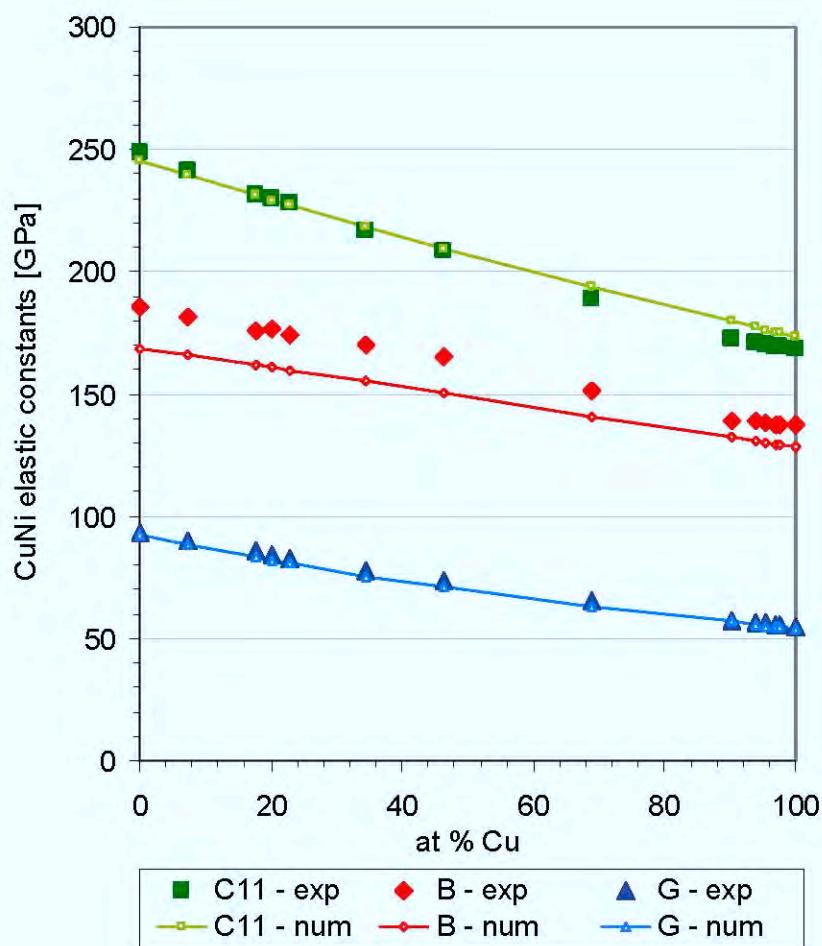
Sutton, A.P. and Chen, J., *Phil. Mag. Letters*, **61**: 139, 1990

Kulkarni, Y., Knap, J. & MO, J. *Mech. Phys. Solids*, **56**:1417, 2008

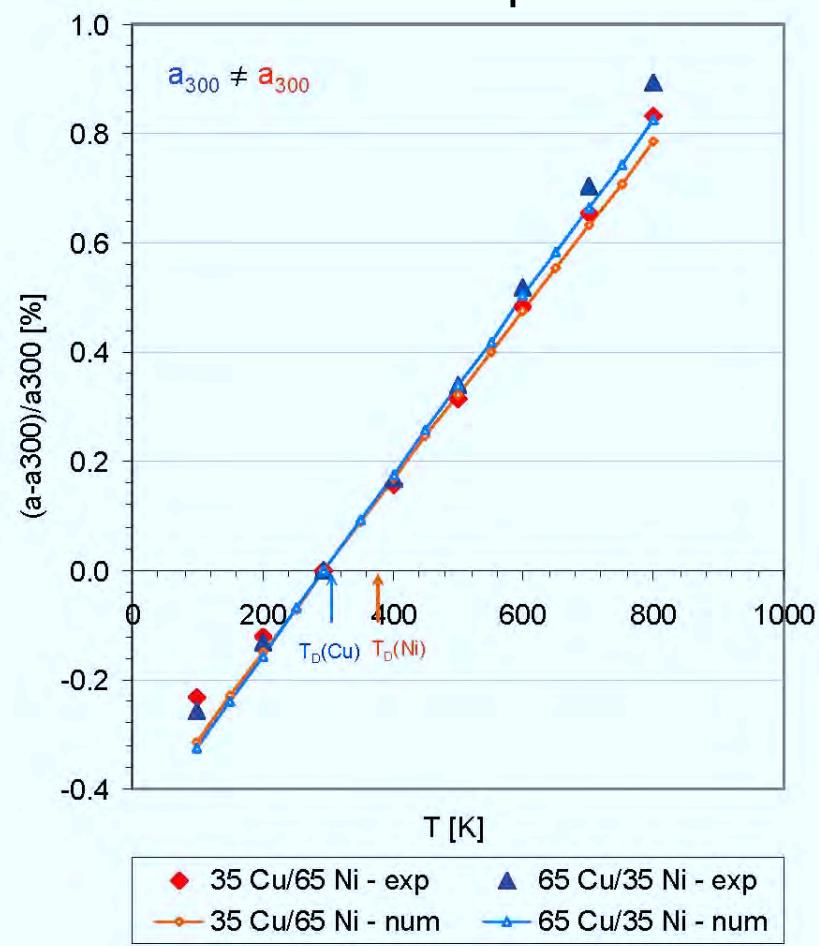


# Equilibrium SM – Meanfield validation

Elastic constants



Thermal expansion



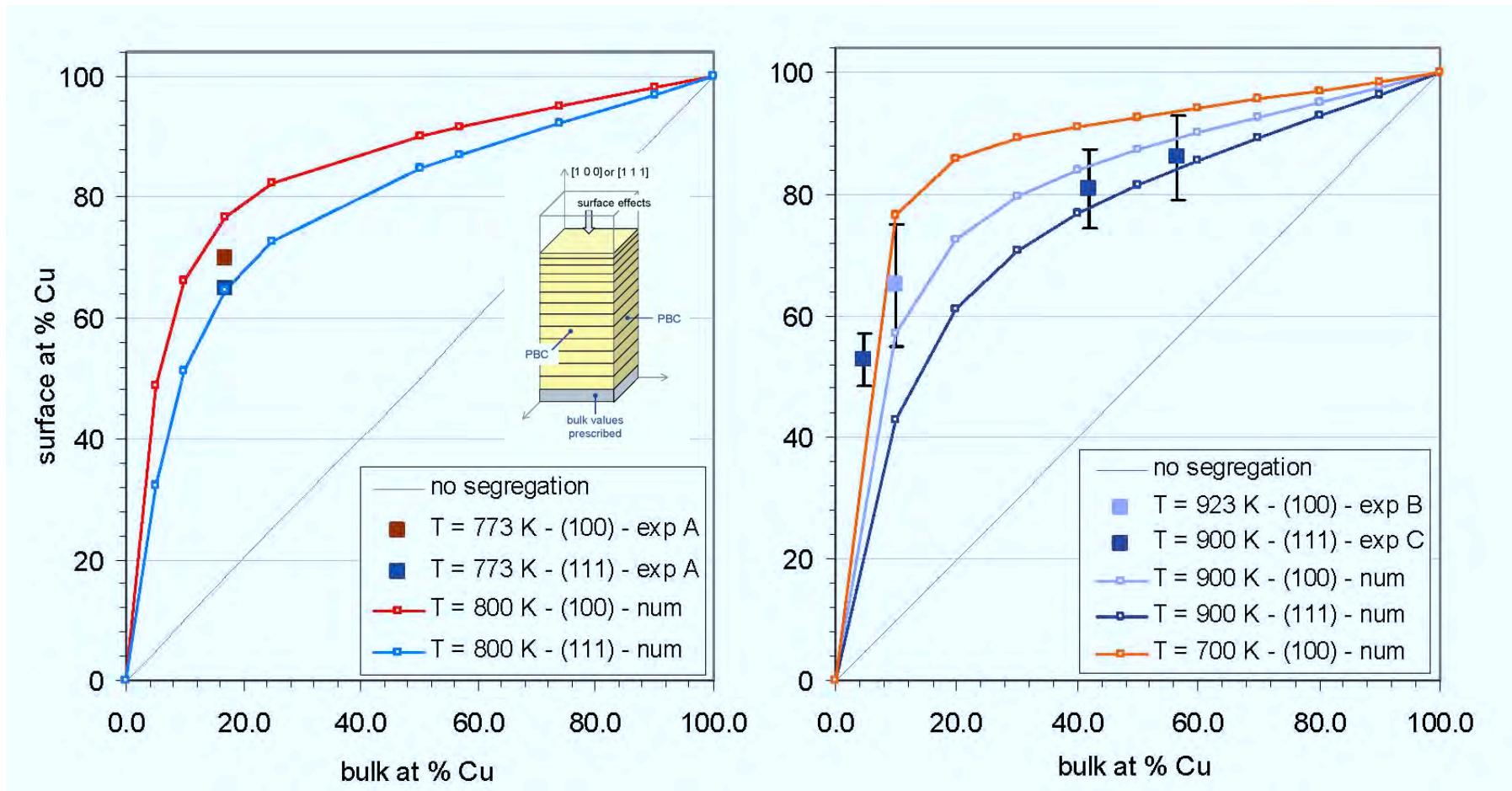
Simmons, G., The MIT Press (1971)

Johnson, R., *Phys. Rev. B*, **39**(17):12554, 1989

Venturini, G., *Doctoral Dissertation*, Caltech, 2011

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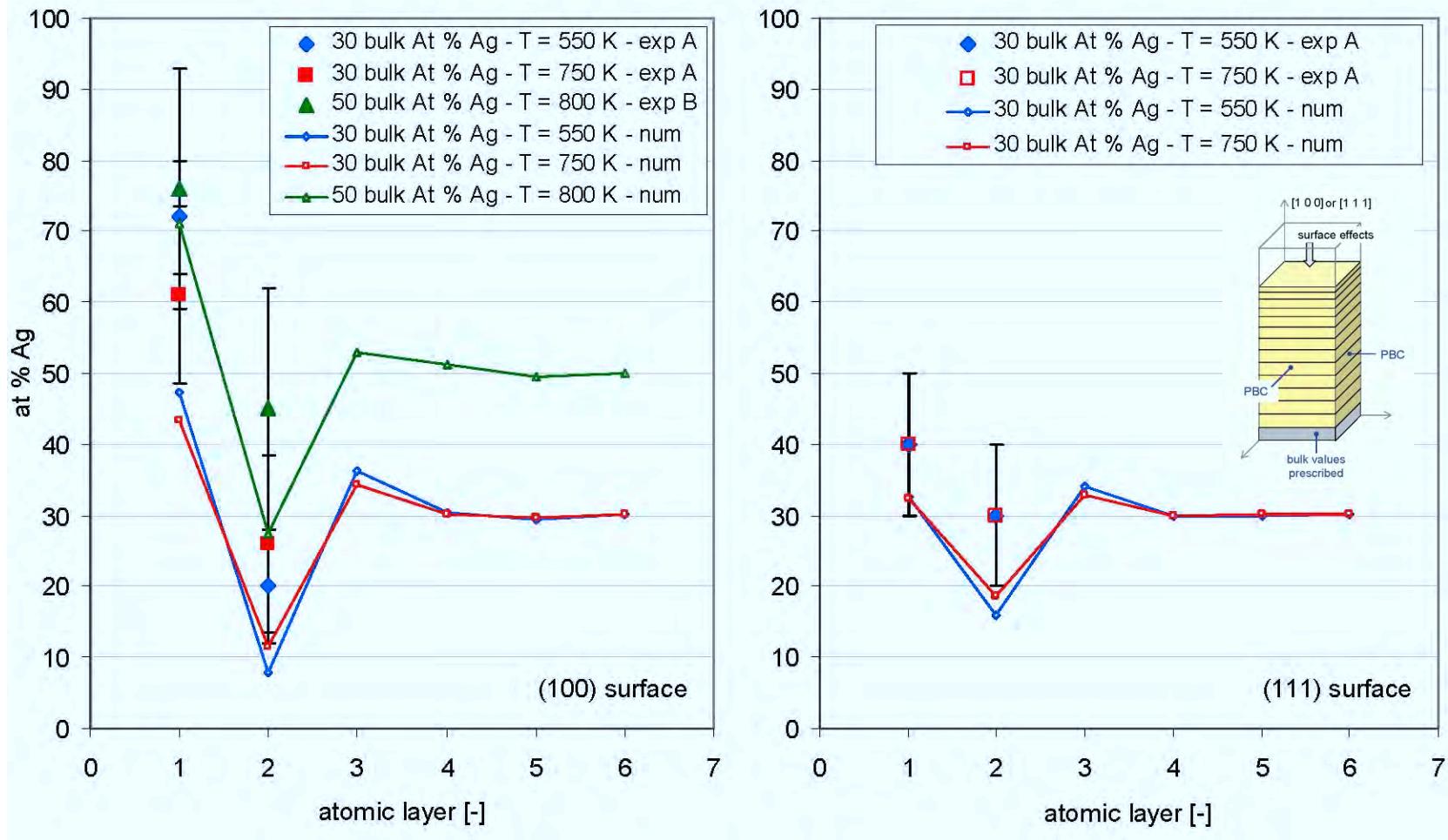
# Equilibrium SM – Meanfield validation



Helms C. and Yu, K., *J. Vac. Sci. Tech.*, **12**:276, 1975  
Sakurai, T. et al., *Phys. Rev. Lett.*, **55**(5):514, 1985  
Venturini, G., *Doctoral Dissertation*, Caltech, 2011

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# Equilibrium SM – Meanfield validation



King, T. and Donnelly, R., Surf. Sci., **151**:374, 1985  
Derry, G. and Wan R., Surf. Sci., **556**:862, 2004  
Venturini, G., *Doctoral Dissertation*, Caltech, 2011



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# Extension to non-equilibrium SM

- Assume  $H = \sum_{i=1}^N h_i$ , otherwise arbitrary.
- Principle of maximum entropy:  $\mathcal{S}[p] \rightarrow \max!$
- subject to:  $\langle h_i \rangle = e_i, \quad \langle n_{ik} \rangle = x_{ik}$
- Lagrangian: reciprocal temperature field  
$$\mathcal{L}[p, \{\beta\}, \{\gamma\}] = \mathcal{S}[p] - k_B \{\beta\}^T \{\langle h \rangle\} - k_B \{\gamma\}^T \{\langle n \rangle\}$$
- Gran-canonical distribution: chemical potential field

$$p = \frac{1}{\Xi} e^{-\{\beta\}^T \{h\} - \{\gamma\}^T \{n\}}$$

$$\Xi = \sum_{\{n\} \in \mathcal{O}_{NM}} \int e^{-\{\beta\}^T \{h\} - \{\gamma\}^T \{n\}} dq dp$$



# Non-equilibrium SM – Macroscopic equilibrium

- Macroscopic variables  $\equiv X$  (e.g., volume, strain...)
- Parametrized Hamiltonians:  $H(\{q\}, \{p\}, \{n\}, X)$
- Gran-canonical thermodynamic potentials:
  - Free entropy:  $\Phi(\{\beta\}, \{\gamma\}, X) = k_B \log \Xi$
- Statement of macroscopic equilibrium:  
 $\Phi(\{\beta\}, \{\gamma\}, \cdot) \rightarrow \max!$  (free entropy principle)



# Non-equilibrium SM – Meanfield theory

- Space of trial *local Hamiltonians*:  $\mathcal{H}_0$
- Free-entropy inequality: For all  $\{h_0\} \in \mathcal{H}_0$ ,
  - i)  $\Phi \geq \Phi_0 - k_B \{\beta\}^T \{\langle h - h_0 \rangle_0\} \equiv -\mathcal{F}[\{h_0\}, \{\beta\}, \{\gamma\}]$
  - ii)  $\Phi = -\mathcal{F}[\{h_0\}, \{\beta\}, \{\gamma\}] \Leftrightarrow \{h_0\} = \{h\}$
- Best approximation:  $\min \mathcal{F}[\{h_0\}, \{\beta\}, \{\gamma\}]$  over  $\mathcal{H}_0$
- Example:  $\mathcal{H}_0 \equiv$  local harmonic oscillators,

$$h_{0i} = \frac{1}{2m(n_i)} |\mathbf{p}_i - \bar{\mathbf{p}}_i|^2 + \frac{m(n_i)\omega_i^2}{2} |\mathbf{q}_i - \bar{\mathbf{q}}_i|^2$$

- Minimize  $\mathcal{F}[\{h_0\}, \{\beta\}, \{\gamma\}]$  w.r.t.  $\{\pi\} \equiv (\{\bar{\mathbf{q}}\}, \{\bar{\mathbf{p}}\}, \{\omega\})$
- Compute  $\langle \cdot \rangle_0$  using Gaussian quadrature



# Non-equilibrium SM – Kinetics (Onsager)

- Need equations of evolution for  $\{e\}$  and  $\{\gamma\}$ .
- Free entropy:  $\Phi = k_B \log \Xi(\{\beta\}, \{\gamma\})$
- Thermodynamic driving-force identities:

$$\beta_i = \frac{1}{k_B} \frac{\partial \Phi^*}{\partial e_i}(\{e\}, \{x\}), \quad \gamma_{ik} = \frac{1}{k_B} \frac{\partial \Phi^*}{\partial x_{ik}}(\{e\}, \{x\})$$

- General kinetic relations (*a la* Onsager):

$$\dot{e}_i = \underbrace{\frac{\partial \psi}{\partial \beta_i}(\{\beta\}, \{\gamma\})}_{\text{Discrete Fourier law}},$$

$$\dot{x}_{ik} = \underbrace{\frac{\partial \psi}{\partial \gamma_{ik}}(\{\beta\}, \{\gamma\})}_{\text{Discrete Fick's law}}$$

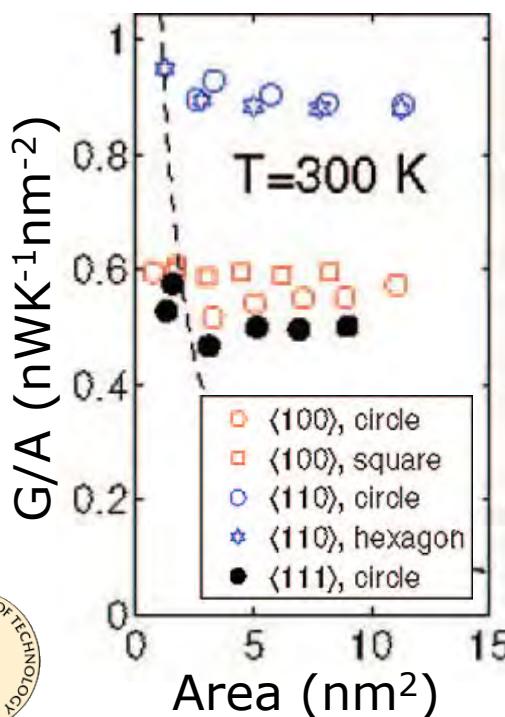
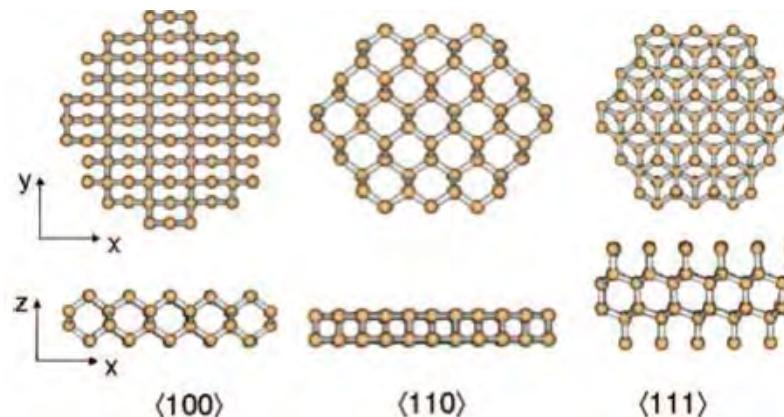
Discrete Fourier law

Discrete Fick's law

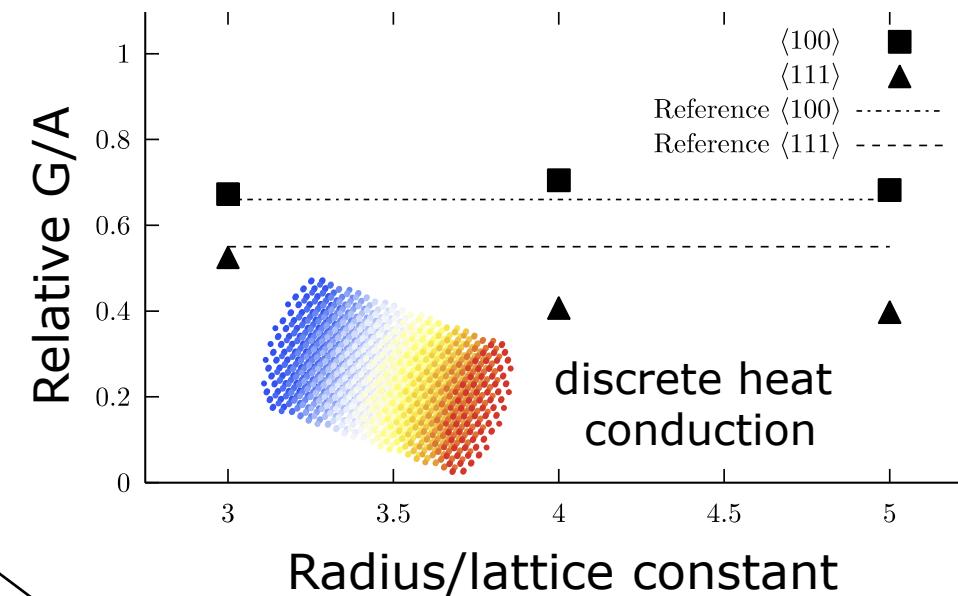
- $\psi$  formulated using discrete differential operators



# Non-equilibrium SM – Validation



Markussen, T., et al.,  
“Heat Conductance Is Strongly  
Anisotropic for Pristine Silicon  
Nanowires”, *Nano Letters*,  
8(11):3771, 2008.



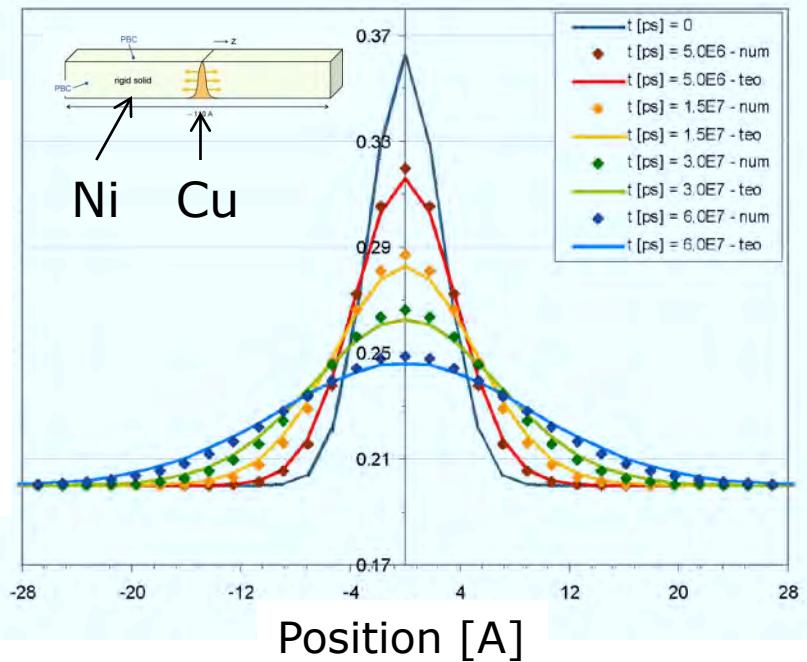
Markussen, T., et al., 2008.

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# Non-equilibrium SM – Validation



Cu atomic fraction

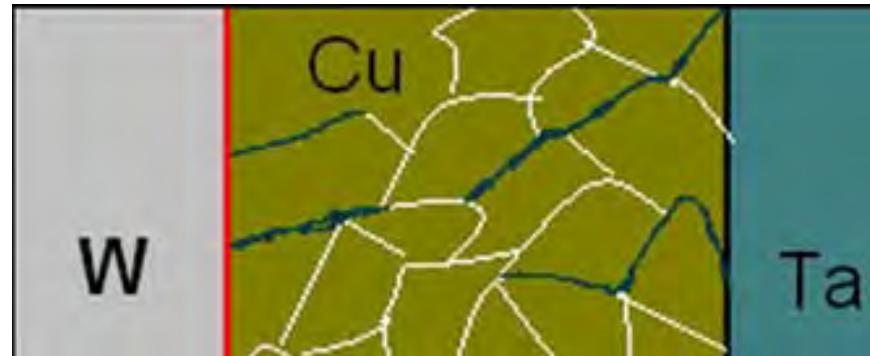
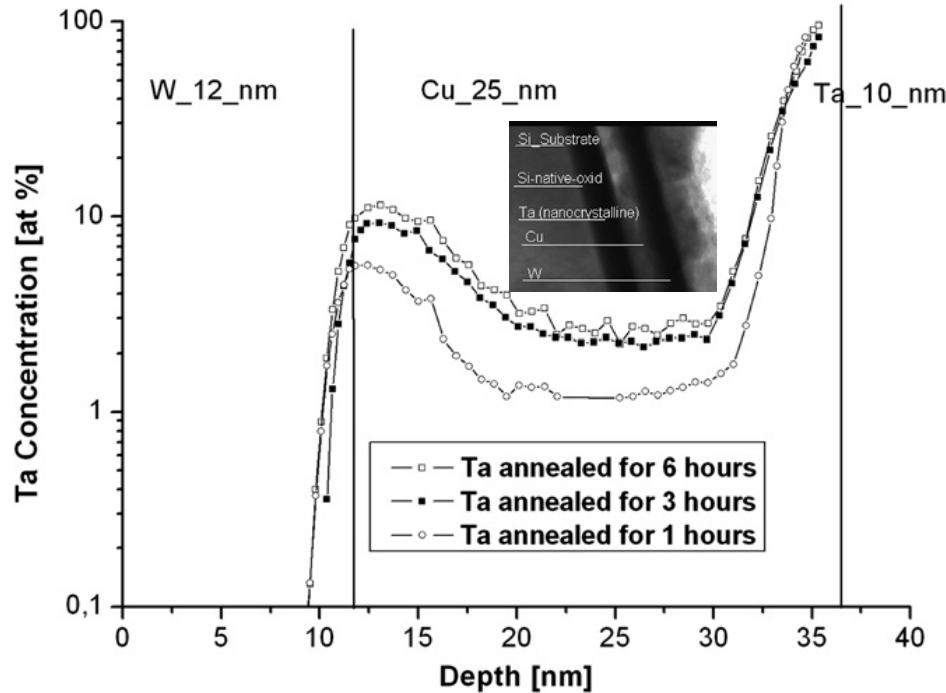


Venturini, G., *Doctoral Dissertation*,  
Caltech, 2011

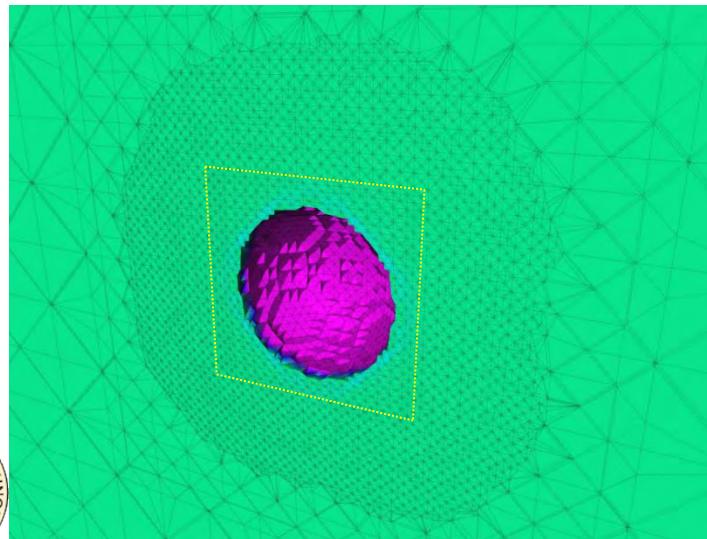
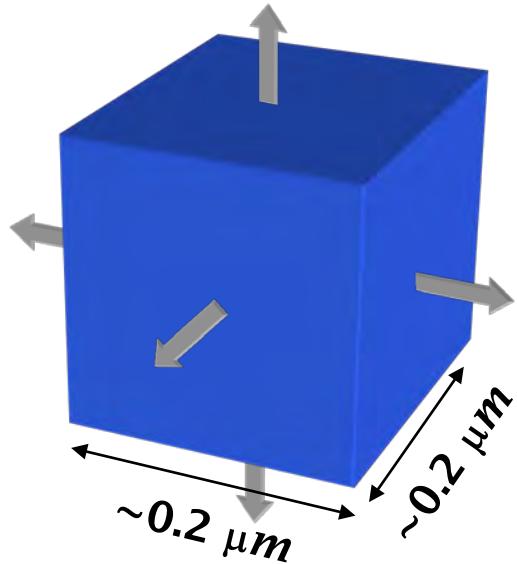
Lakatos, A. et al., "Investigations of diffusion kinetics in Si/Ta/Cu/W and Si/Co/Ta systems by secondary neutral mass spectrometry",  
*Vacuum*, 85:493, 2010



The Log(Concentration)-depth profile of Ta



# Application to nanovoid plastic cavitation



- Parameters:
  - $T_0 = 300K$
  - Full Size =  $72a_0$
  - Atomistic Zone =  $14a_0$
  - Diameter =  $12a_0$
  - Strain Rate =  $< 10^{10} \text{ 1/s}$
- Loading:
  - Triaxial, uniaxial
- Material:
  - Copper, Mishin, Y.,  
Phys. Rev. B. 2001,

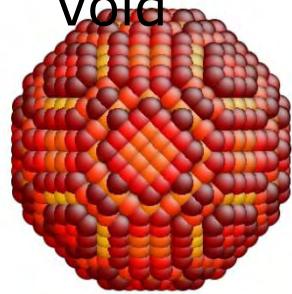
Initial quasicontinuum  
mesh with full atomistic  
resolution near void



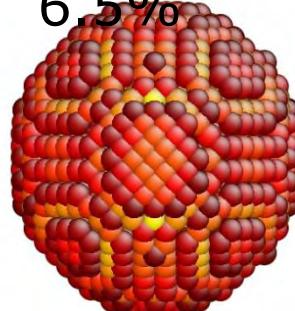
# Nanovoid plastic cavitation in Cu

Uniaxial loading

Initial  
Void

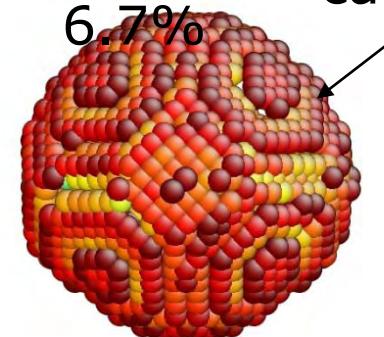


Void at  
6.5%

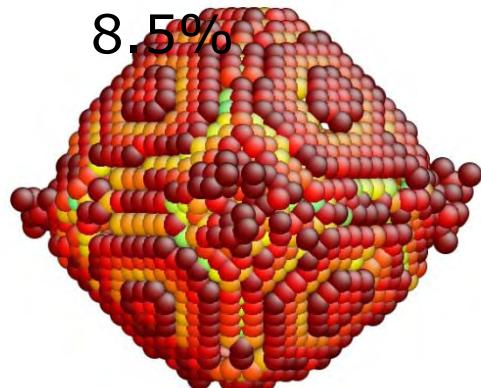


Growth on transversal  
direction when  
dislocation are emitted

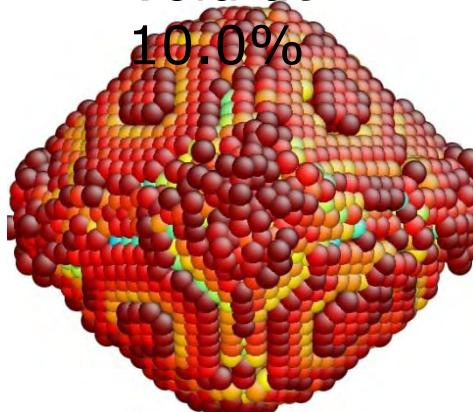
Void at  
6.7%



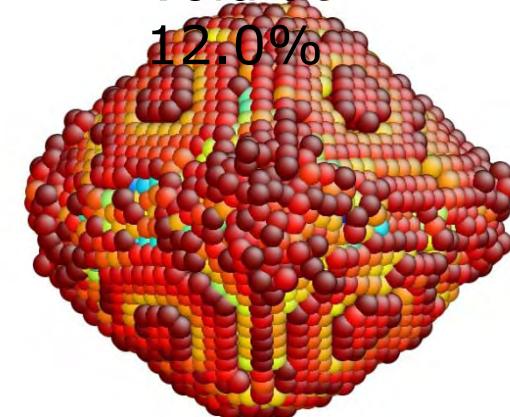
Void at  
8.5%



Void at  
10.0%



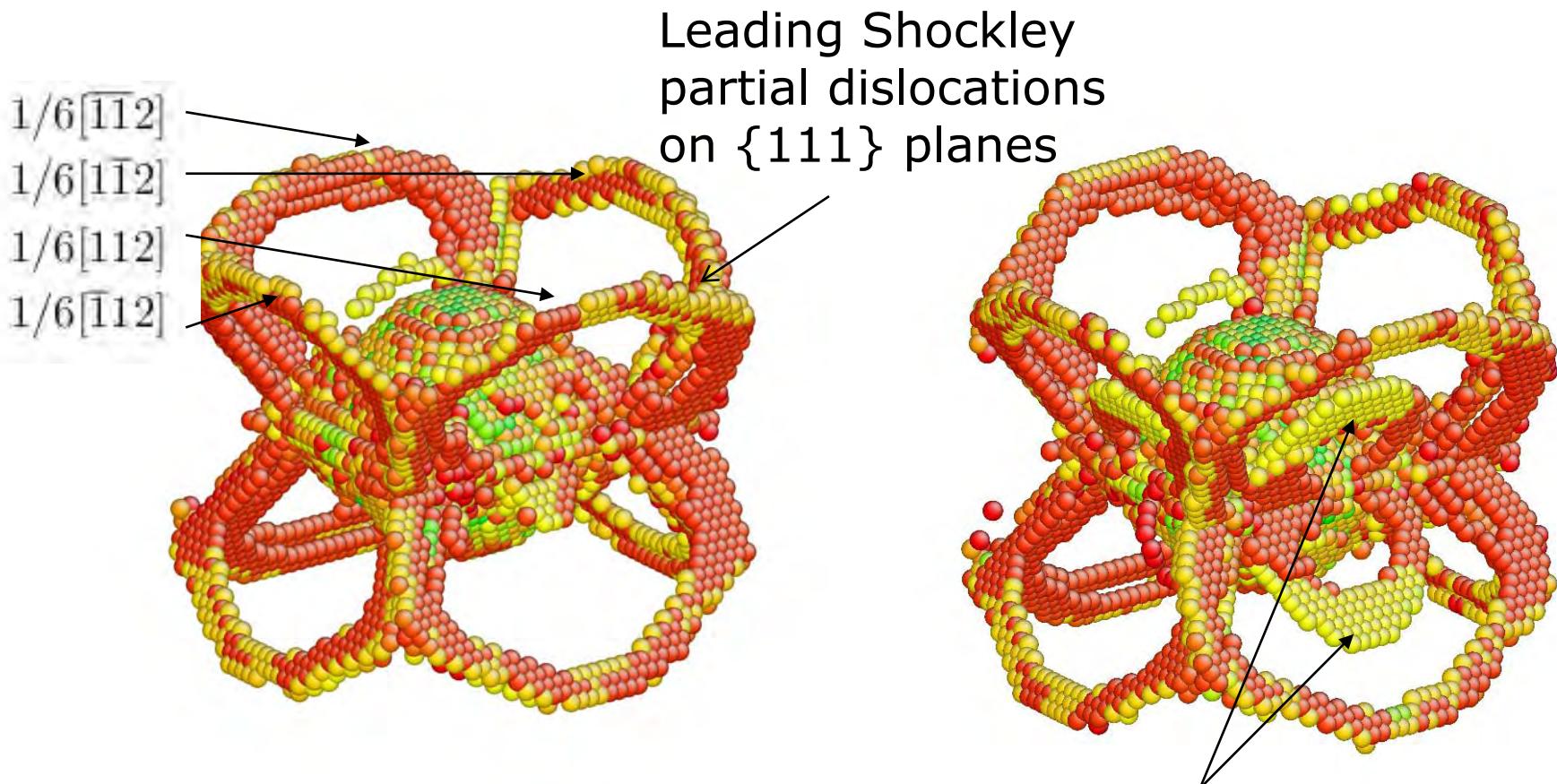
Void at  
12.0%



Final shape of void,  
octahedral on 8 planes on  $\{111\}$  Michael Ortiz  
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# Nanovoid plastic cavitation in Cu



Uniaxial loading,  $\epsilon$   
= 6.6%

Leading Shockley  
partial dislocations  
on  $\{111\}$  planes

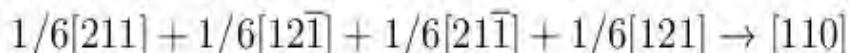
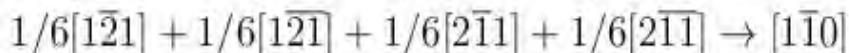
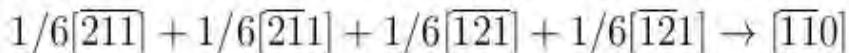
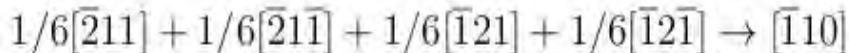
Trailing Shockley partial  
dislocations on  $\{111\}$  planes



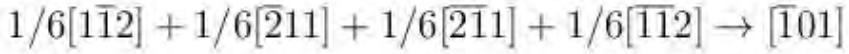
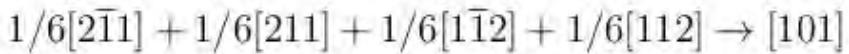
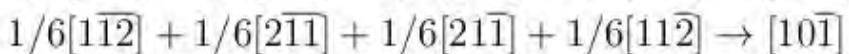
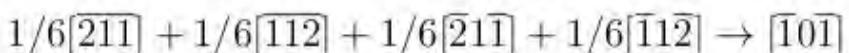
# Nanovoid plastic cavitation in Cu

Shear to Prismatic loop reactions:

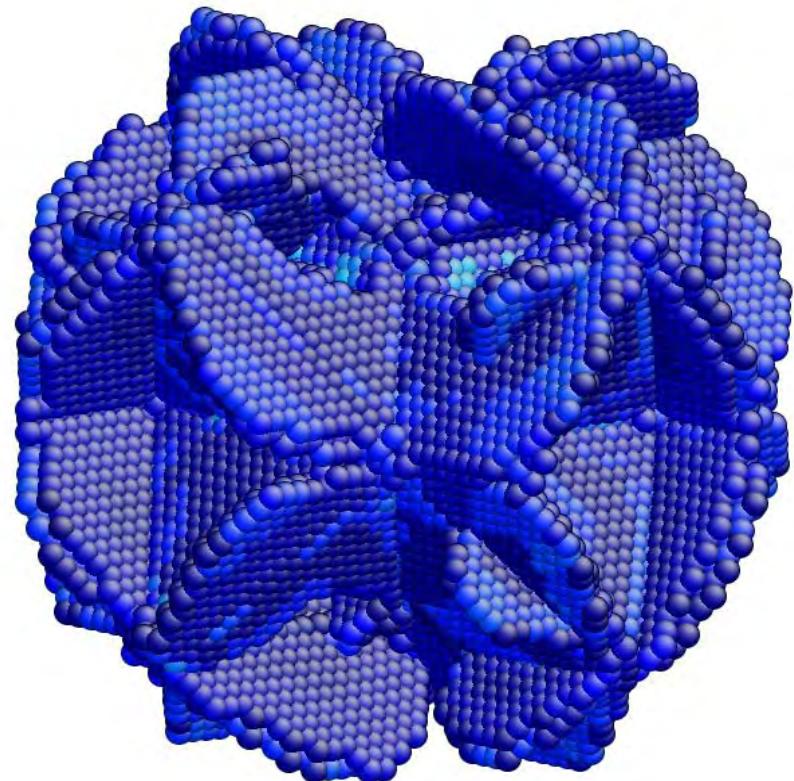
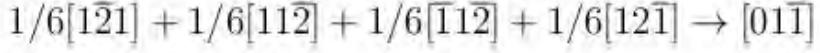
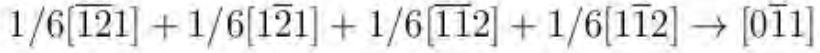
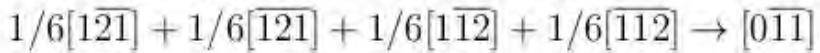
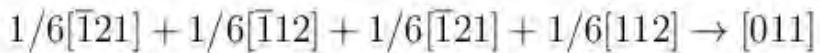
On  $<110>$  directions



On  $<110>$  directions



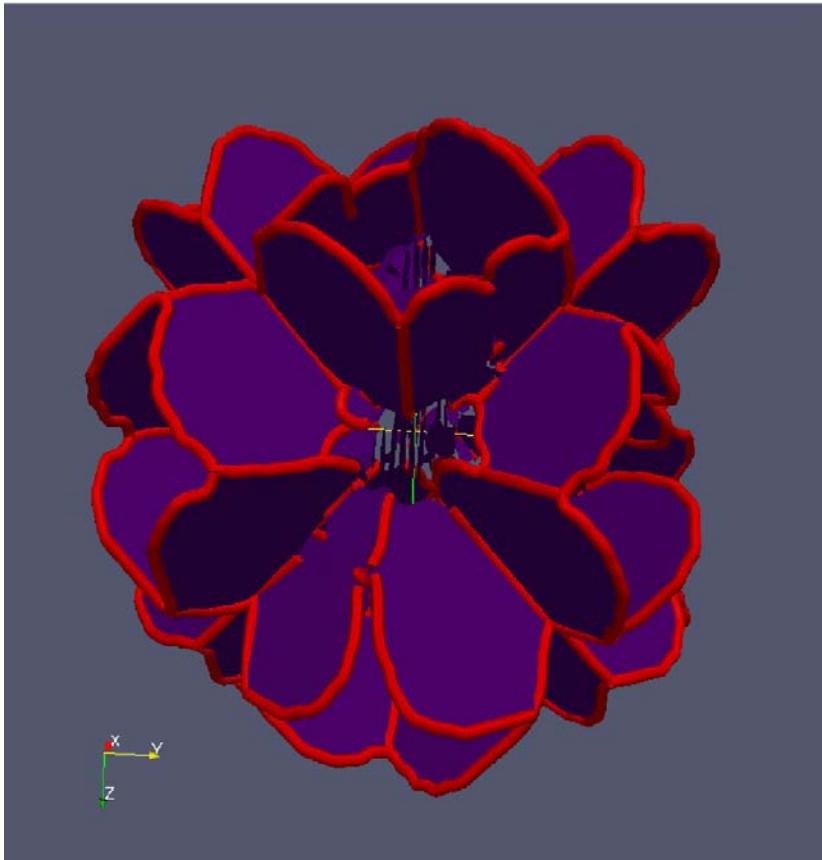
On  $<110>$  directions



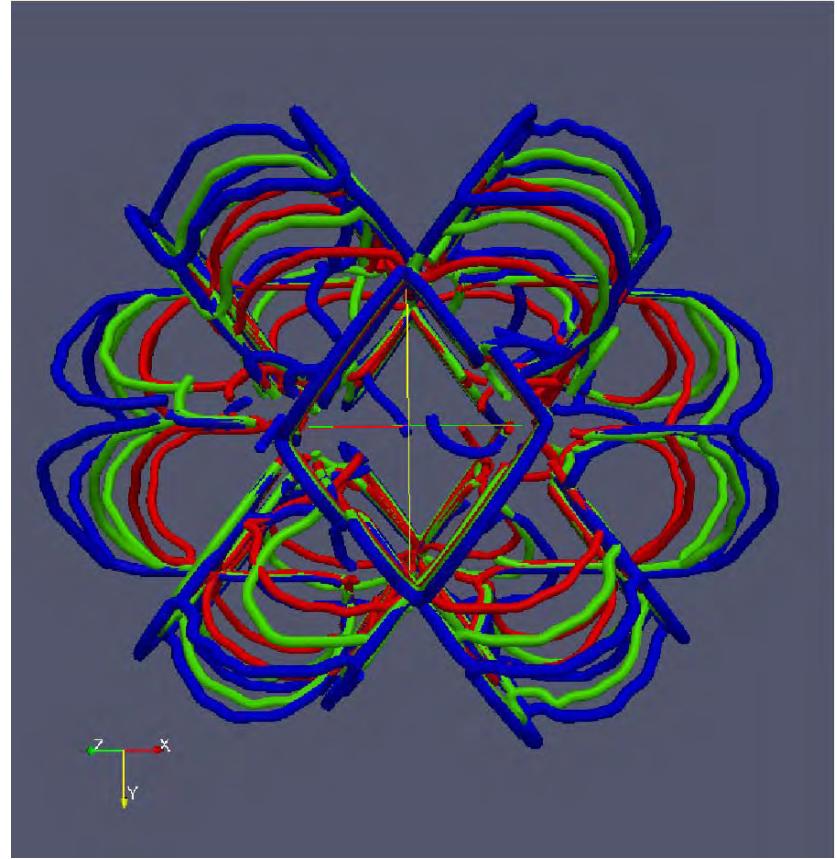
Triaxial loading



# Nanovoid plastic cavitation in Cu



Prismatic loop structure,  
triaxial loading



Prismatic loop evolution  
( $\epsilon = 5, 6, 7\%$ )

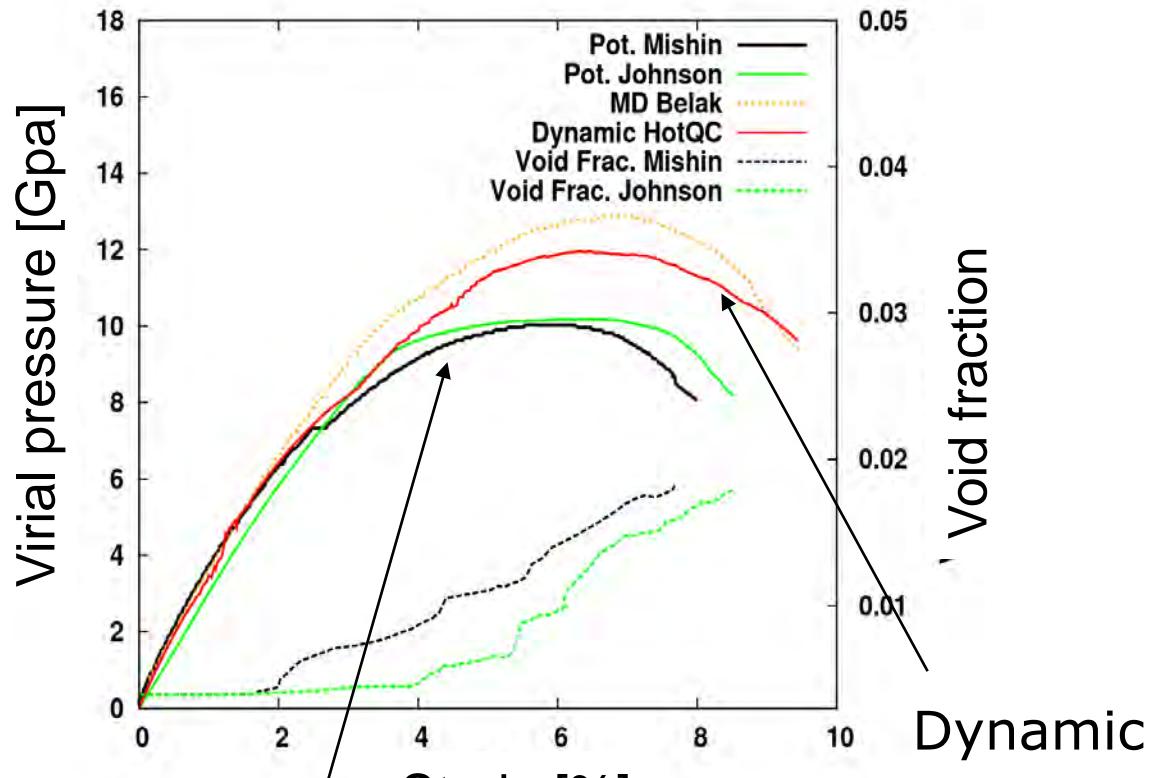


(Images obtained with **DXA** and **Paraview**)

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# Nanovoid plastic cavitation in Cu

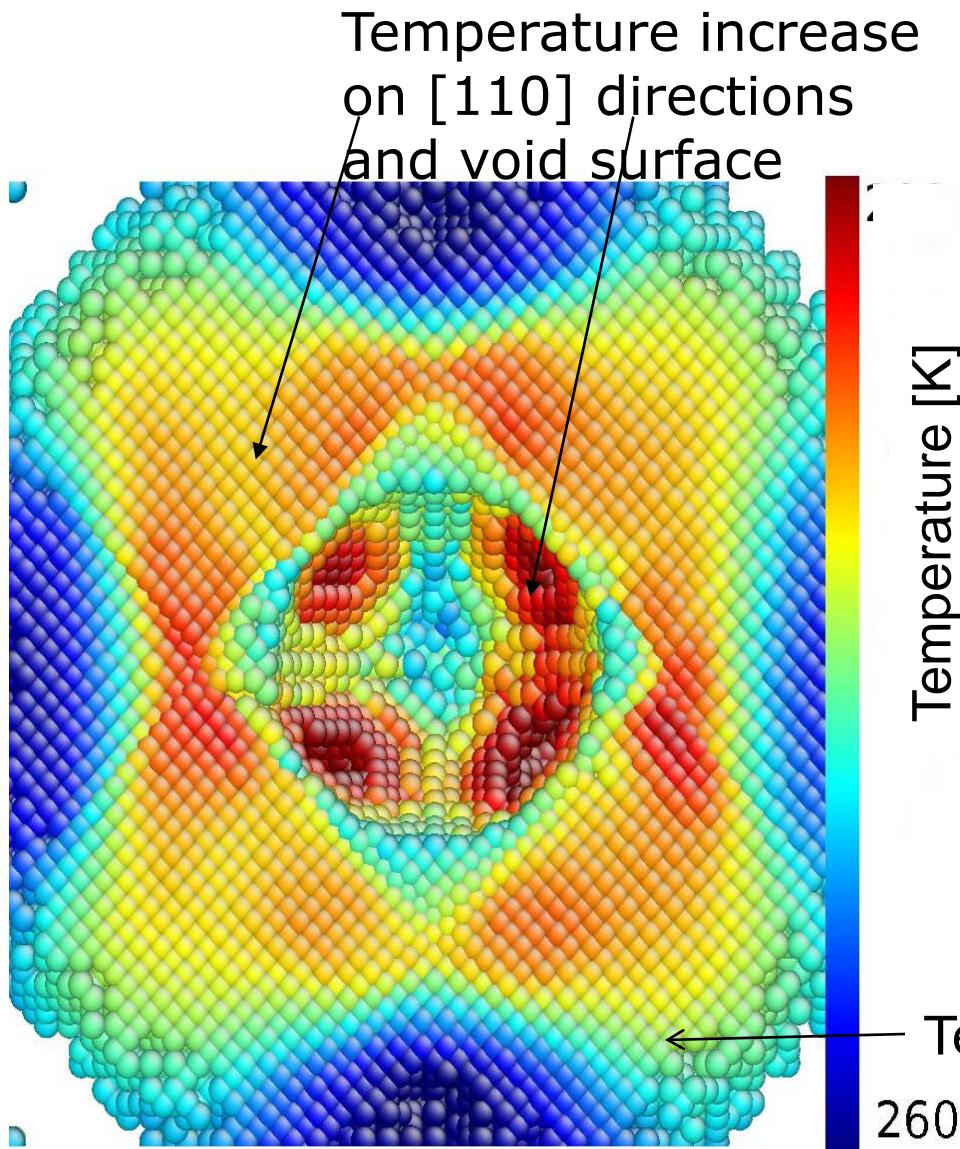
Triaxial loading ( $d\varepsilon/dt = 10^{10} \text{ 1/s}$ )



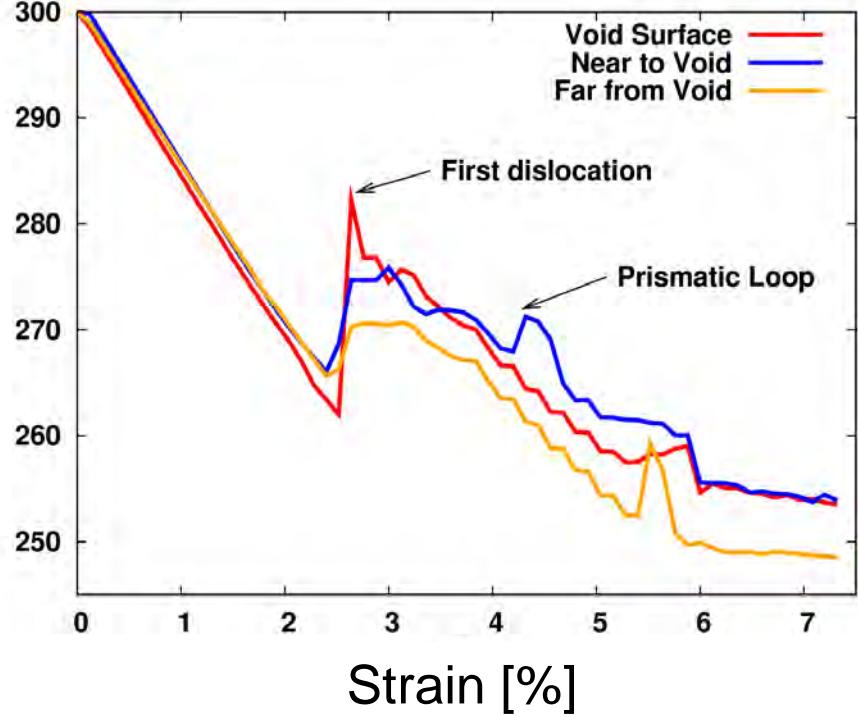
Quasistatic



# Nanovoid plastic cavitation in Cu



Temperature evolution



# To Professor Rodney J. Clifton



Scientist, teacher, mentor



Michael Ortiz  
RJC 06/12