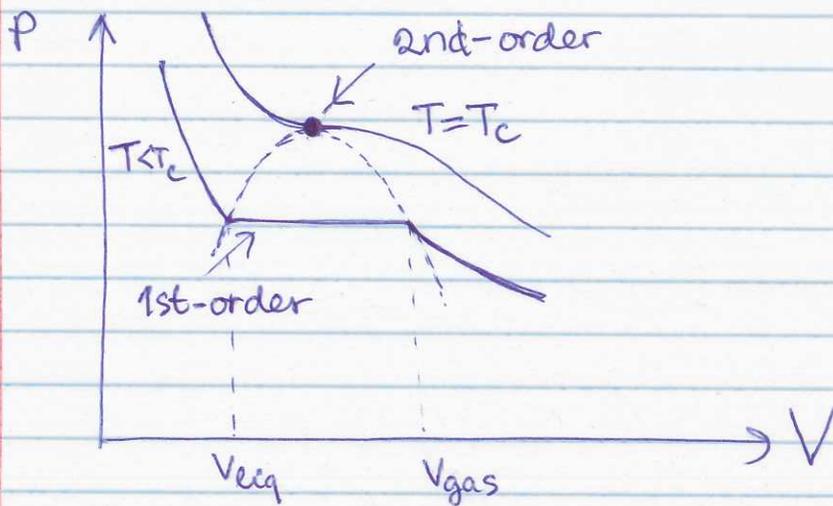


# Critical point behavior



For any liquid/gas phase transition (1st, 2nd order)

for total system

$$\left\{ \begin{array}{l} \alpha_p \equiv \frac{1}{V} \left( \frac{\partial V}{\partial T} \right)_p = \infty \text{ (thermal expansion)} \\ \text{coeff.} \\ \alpha_T \equiv -\frac{1}{V} \left( \frac{\partial V}{\partial p} \right)_T = \infty \text{ (isothermal compressibility)} \\ C_p \equiv \left( \frac{\partial H}{\partial T} \right)_p = \infty \text{ (heat capacity at const } p \text{)} \\ \text{enthalpy} \\ C_v = \left( \frac{\partial U}{\partial T} \right)_v \leq \infty \end{array} \right.$$

Proof: Along the horizontal <sup>(liquid-gas)</sup> coexistence segment, we have

$$\left. \begin{array}{l} \Delta T = \Delta p = 0, \text{ while } \Delta V \neq 0 \end{array} \right\} \Rightarrow \begin{array}{l} \alpha_p = \infty \\ \alpha_T = \infty \end{array}$$

Also, we supply heat to convert liquid to gas while  $\Delta T = 0$

$$C_p = \left( \frac{\delta Q}{\delta T} \right)_p = \infty$$

## For 2nd-order liquid-gas transition

| Crit. exponents $\alpha, \beta, \gamma, \delta$<br>for $T \rightarrow T_c$  | VdW                   | experiment         |
|---|-----------------------|--------------------|
| <ul style="list-style-type: none"> <li><math>C_v \sim  T - T_c ^{-\alpha}</math>, (<math>V = V_c</math>)</li> </ul>   | $\alpha = 0$          | $0 < \alpha < 0.4$ |
| <ul style="list-style-type: none"> <li><math>n_{\text{liq}} - n_{\text{gas}} \sim (T_c - T)^\beta</math>, <math>T &lt; T_c</math><br/><math>p = p_{\text{coex.}}(T)</math></li> </ul> | $\beta = \frac{1}{2}$ | $\beta = 0.35$     |
|   |                       |                    |
| "order parameter" exponent  |                       |                    |
| <ul style="list-style-type: none"> <li><math>\chi_T \sim  T - T_c ^{-\gamma}</math>, <math>V = V_c</math></li> </ul>  | $\gamma = 1$          | $\gamma = 1.26$    |
| "susceptibility" exponent   |                       |                    |
| <ul style="list-style-type: none"> <li><math> p - p_c  \sim  n - n_c ^\delta</math>, <math>T = T_c</math><br/>critical isotherm</li> </ul>  | $\delta = 3$          | $\delta = 4.4$     |

## Proof of VdW exponents $\beta, \gamma, \delta$

VdW equation of state in reduced variables:

$$\bar{p} = \frac{8\bar{T}}{3\bar{v} - 1} - \frac{3}{\bar{v}^2}, \quad \bar{p}_c = \bar{T}_c = \bar{v}_c = 1$$

Shifted variables

$$\delta \bar{p} = \bar{p} - 1$$

$$\delta \bar{v} = \bar{v} - 1$$

$$\delta T = T - 1$$

$$\delta \bar{p} = -1 + \frac{8(\delta T + 1)}{3\delta \bar{v} + 2} - \frac{3}{(1 + \delta \bar{v})^2}$$

Drop "x" from now on; expand rhs for small  $\delta T$ ,  $\delta v$ :

$$\left[ \delta p \approx -1 + \frac{8(1 + \delta T)}{2} \left( 1 - \frac{3}{2} \delta v + \frac{9}{4} \delta v^2 - \frac{27}{8} \delta v^3 + \dots \right) \right.$$

$$\left. - 3 \left( 1 - (2\delta v + \delta v^2) + (2\delta v + \delta v^2)^2 - (2\delta v + \delta v^2)^3 \right) \right] =$$

↓  
drop

$$= \delta T \left( 4 - 6\delta v + 9\delta v^2 - \frac{27}{2} \delta v^3 \right)$$

$$- 6\delta v + 9\delta v^2 - \frac{27}{2} \delta v^3 + 6\delta v + 3\delta v^2 - 3.4\delta v^2 -$$

$$- 3.4\delta v^3 + 3.8\delta v^3 =$$

$$= \delta T \cdot \left( 4 - 6\delta v + 9\delta v^2 - \frac{27}{2} \delta v^3 \right) - \frac{3}{2} \delta v^3$$

Critical isotherm :  $\delta T = 0$

$$T = T_c$$

quick proof:

$T_c, v_c$  was defined

$$\text{by } \frac{\partial p}{\partial v} = \frac{\partial^2 p}{\partial v^2} = 0$$

$$\Rightarrow p - p_c \sim \frac{1}{3!} \frac{\partial^3 p}{\partial v^3} (v - v_c)^3$$

$$\delta \bar{p} \approx -\frac{3}{2} \delta \bar{v}^3$$

$$\left[ \frac{p - p_c}{p_c} \approx -\frac{3}{2} \left( \frac{v - v_c}{v_c} \right)^3 = -\frac{3}{2} \left( \frac{\frac{1}{p} - \frac{1}{p_c}}{\frac{1}{p_c}} \right)^3 = \right.$$

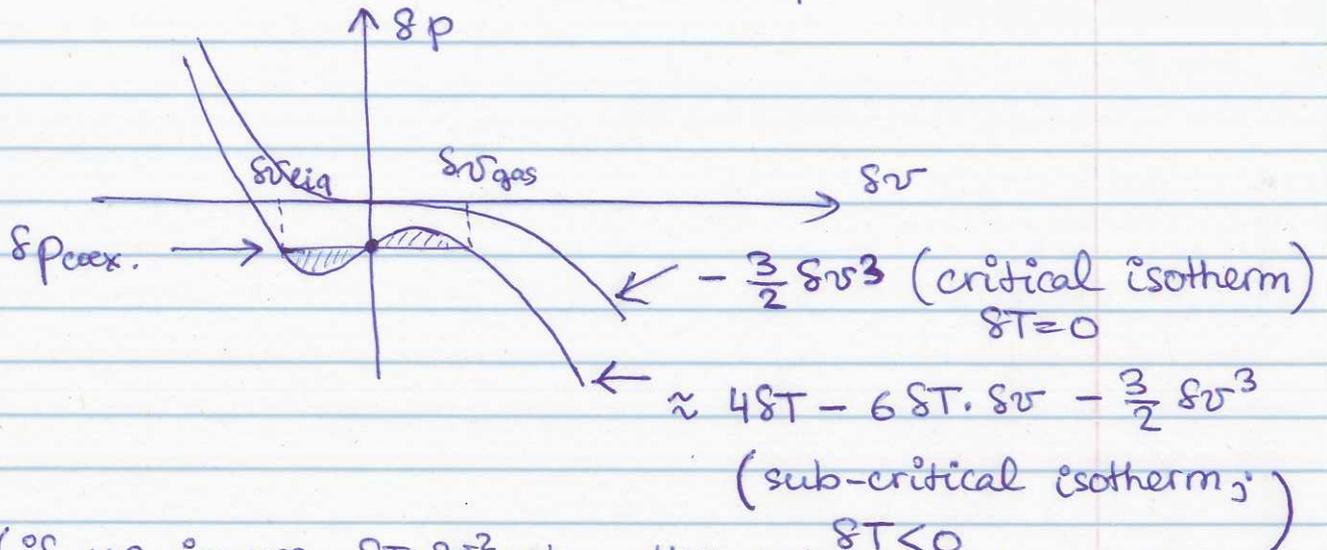
$$\left. = -\frac{3}{2} \left( \frac{p_c - p}{p} \right)^3 \approx -\frac{3}{2} \left( \frac{p_c - p}{p_c} \right)^3 \right]$$

$\Rightarrow$  exponent

$$\boxed{\delta = 3}$$

## Order parameter:

Maxwell construction near critical point



(if we ignore  $\delta T \cdot \delta v^2$ , etc., the subcrit isotherm is "odd" wrt point  $(0, 4\delta T)$  and the Maxwell construction is trivial)

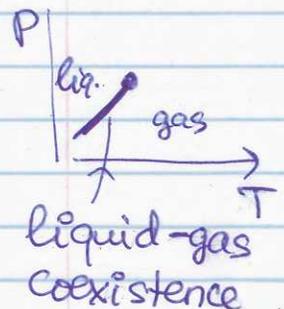
$$\Rightarrow \delta p_{cox} = 4\delta T \stackrel{!}{=} 4\delta T - 6\delta T \delta v_{liq} - \frac{3}{2} \delta v_{liq}^3$$

$$\Rightarrow \boxed{\delta v_{liquid/gas} = \pm 2\sqrt{-\delta T}}$$

$$\Rightarrow \boxed{\frac{n_{liq} - n_{gas}}{n_c} = \frac{1}{\bar{v}_{liq}} - \frac{1}{\bar{v}_{gas}} = \frac{1}{1 + \delta v_{liq}} - \frac{1}{1 + \delta v_{gas}} \approx}$$

$$\approx \delta v_{gas} - \delta v_{liq} = 4\sqrt{-\delta T} = 4 \left( \frac{T_c - T}{T_c} \right)^{1/2}$$

$$\Rightarrow \boxed{\beta = \frac{1}{2}} \text{ - order parameter exponent}$$



$$-\frac{1}{P_c(1+\delta\bar{v})} \cdot \left[ -3 \frac{8(1+\delta\bar{T})}{(3\delta\bar{v}+2)^2} + \frac{6}{(1+\delta\bar{v})^3} \right]^{-1}$$

## Compressibility

$$\alpha_T = -\frac{1}{v} \left[ \left( \frac{\partial P}{\partial v} \right)_T \right]^{-1} = -\frac{1}{(1+\delta\bar{v})P_c} \left[ \left( \frac{\partial \delta\bar{P}}{\partial \delta\bar{v}} \right)_{\delta\bar{T}} \right]^{-1}$$

$$\left( \frac{\partial P}{\partial v} \right)_T = \frac{P_c}{v_c} \left( \frac{\partial \bar{P}}{\partial \bar{v}} \right)_T = \frac{P_c}{v_c} \left( \frac{\partial \delta\bar{P}}{\partial \delta\bar{v}} \right)_{\delta\bar{T}}$$

$$v = v_c(1+\delta\bar{v})$$

Above critical point  $\delta\bar{T} > 0$   
~~at  $v=v_c$~~  at  $v=v_c$  (point  $\boxtimes$ )

$$\left( \frac{\partial \delta\bar{P}}{\partial \delta\bar{v}} \right)_{\delta\bar{T}} = \delta\bar{T} \cdot (-6)$$

$$\alpha_T = -\frac{1}{P_c} \cdot \frac{1}{-6\delta\bar{T}} = \frac{1}{6P_c} \frac{1}{\delta\bar{T}} = \frac{1}{6P_c} \left( \frac{T-T_c}{T_c} \right)^{-1}$$

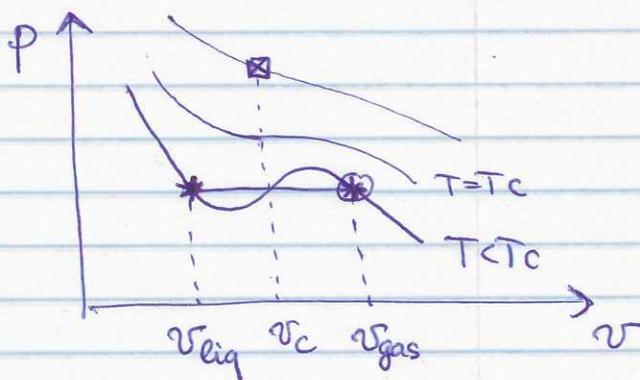
$T > T_c$

$$\Rightarrow \gamma = 1$$

Below the critical point

At points  $*$ ,  $\boxtimes$

$$\alpha_{T, \text{liq or gas}} \approx -\frac{1}{P_c(1+\delta\bar{v}_{\text{liq/gas}})}$$



$$\alpha_T \approx -\frac{1}{6P_c} \left( \frac{1}{1+3\delta\bar{v}+3\delta\bar{v}^2+\delta\bar{v}^3} - \frac{1+\delta\bar{T}}{1+2 \cdot \frac{3}{2}\delta\bar{v} + \left(\frac{3}{2}\delta\bar{v}\right)^2} \right)$$

$$\approx -\frac{1}{6P_c} \times \frac{1}{1+3\delta\bar{v} + \frac{9}{4}\delta\bar{v}^2 - (1+3\delta\bar{v} + 3\delta\bar{v}^2 + O(\delta\bar{v}^3)) - \delta\bar{T} - O(\delta\bar{T} \cdot \delta\bar{v})}$$

$$= -\frac{1}{6P_c} \times \frac{1}{-\delta\bar{T} - \frac{3}{4}\delta\bar{v}_{\text{liq/gas}}^2} = -\frac{1}{6P_c} \times \frac{1}{-\delta\bar{T} - 3(-\delta\bar{T})} = \frac{1}{12P_c \cdot (-\delta\bar{T})}$$

$$\chi_{T, \text{liq or gas}} \approx \frac{1}{12pc} \left( \frac{T_c - T}{T} \right)^{-1}, \quad T < T_c$$

⇒ critical exponent for  $\chi_T$  is the same above and below  $T_c$  ( $\gamma=1$ ), but critical amplitude (prefactor  $\frac{1}{6pc}$  and  $\frac{1}{12pc}$ ) is not!

⇒ behavior of  $\chi_{T, \text{gas}}$  is not symmetric wrt  $T_c$ , i.e.  $\chi_{T, \text{gas}}$  is not only a fnctn of  $|T - T_c|$ .

Mnemonic

Terminology for result amplitude $\boxtimes$  = 2 amplitude $\otimes$  :

$\boxtimes$  is twice as close to crit point as  $\otimes$

$\boxtimes$  is removed from crit point only in temperature  
 $\otimes$  —||— both in temperature and volume

Remark: General relations among critical exponents (valid for any system):

$$\left. \begin{array}{l} \gamma = \beta(\delta - 1) \quad \text{Widom scaling law} \\ \alpha + 2\beta + \gamma = 2 \quad \text{Rushbrooke -||-} \end{array} \right\} \text{satisfied by VdW exponents}$$

↑  
Will discuss later from scaling theory

VdW exponents = mean field exponents.