

# Valence fluctuation and strange metal phase in $\text{YbAlB}_4$

**Yosuke Matsumoto**

*Institute for Solid State Physics, University of Tokyo, Japan*

**Collaborators:**

S. Nakatsuji, K. Kuga, Y. Karaki, Y. Shimura, T. Sakakibara, T. Tomita,  
Y. Uwatoko, A. H. Nevidomskyy, P. Coleman

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- Kondo lattice behavior in the mixed valence systems:  $\beta$ -,  $\alpha$ - $\text{YbAlB}_4$
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- Strange metal phase detached from magnetism in  $\beta$ - $\text{YbAlB}_4$

# Kondo lattice vs. Mixed valence systems

## “Kondo lattice” with integer valence

- Doniach like phase diagram  
**Kondo vs. RKKY**
- Quantum criticality based on critical spin fluctuations associated with **magnetic QCP**.

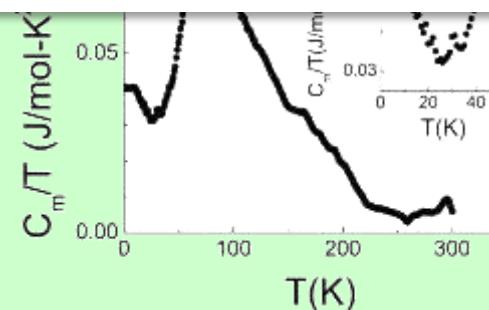
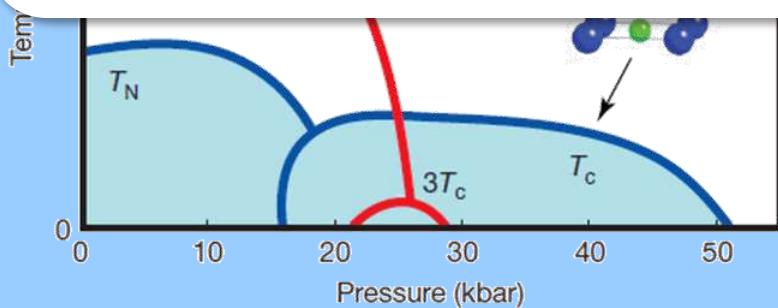
## Mixed valence systems

- Large  $T_K$
- Pauli paramagnetism below  $T_K$
- $\gamma$  ( $= C/T$ ) value is not large.

**Basically, no QC has been expected.**

In contrast,

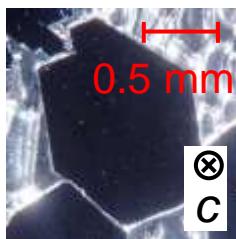
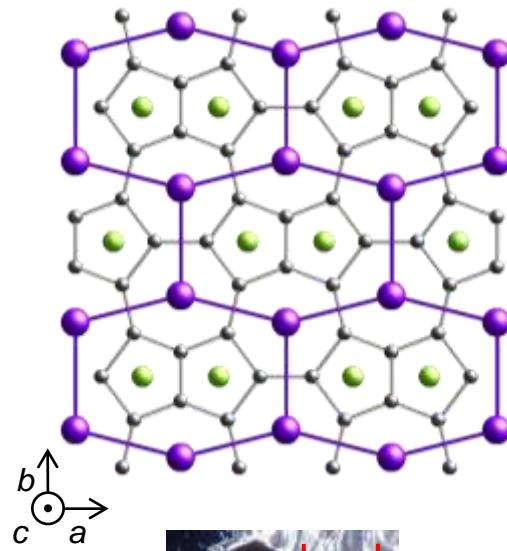
- $\alpha$ -,  $\beta$ -YbAlB<sub>4</sub> provide unique examples of quantum criticality in the mixed valence systems.
- What is the origin ?
- Novel metallic state ?



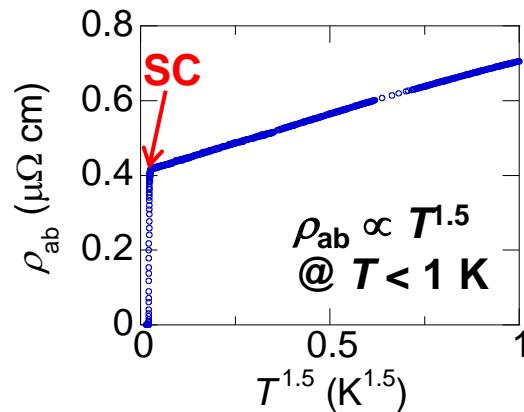
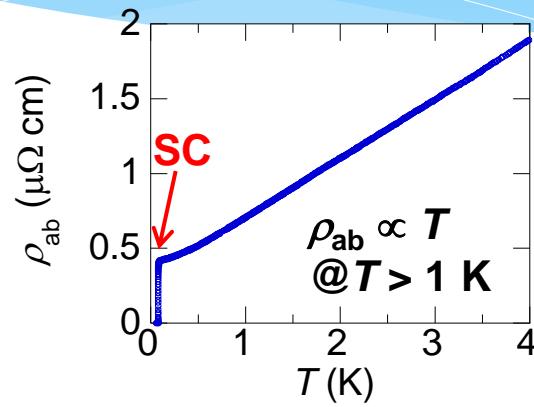
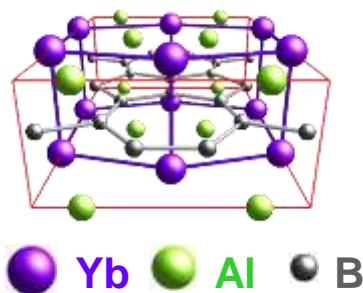
A. L. Cornelius *et al.*,  
PRL **88**, 117201 (2002).

# HF Compounds: $\beta$ -YbAlB<sub>4</sub>

- Yb form distorted hexagonal lattice.
- Layered but Yb-Yb distance is shortest in c axis.
- **The first heavy fermion SC ( $T_c = 80$  mK) in Yb based system.**
- Pronounced NFL behavior in the normal state @  $B = 0$ .

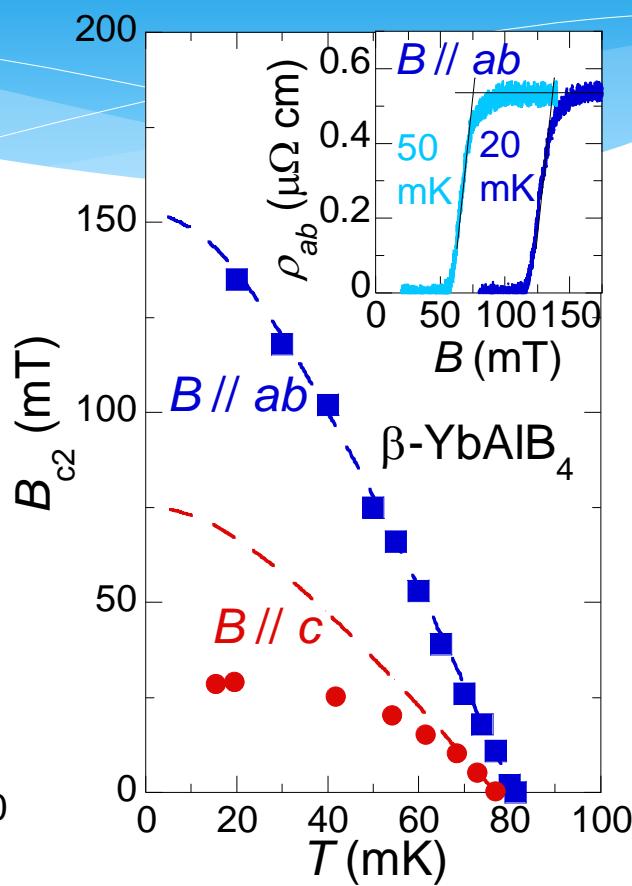
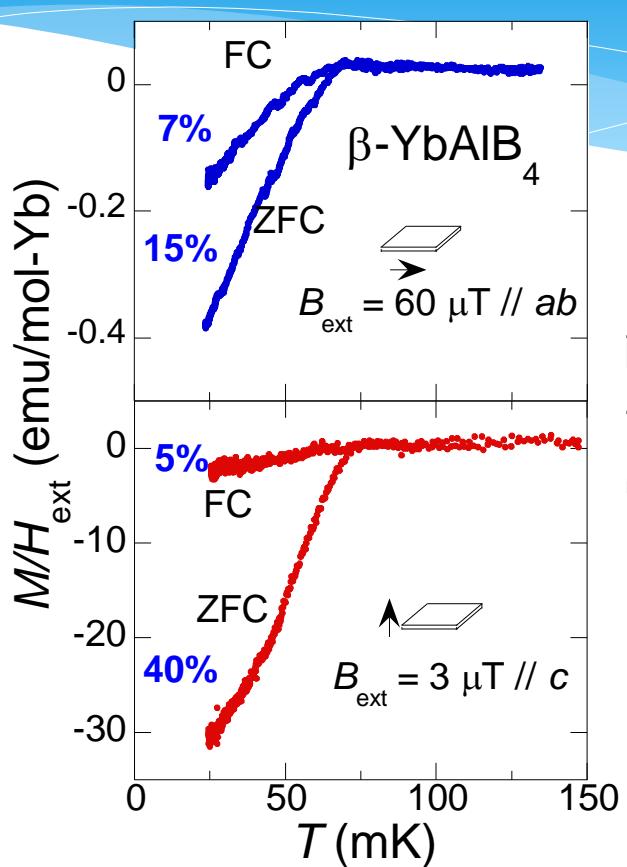
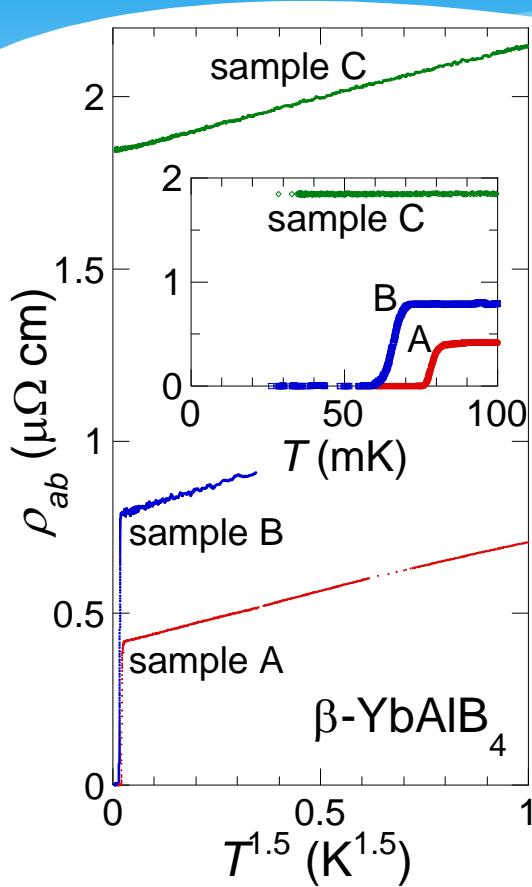


Thickness:  $\sim 10 \mu\text{m}$



# SC in $\beta$ -YbAlB<sub>4</sub>

K. Kuga, Y. Karaki, Y. Matsumoto, Y. Machida, and S. Nakatsuji,  
PRL 101, 137004 (2008).

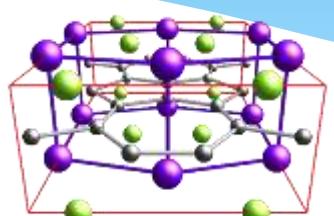
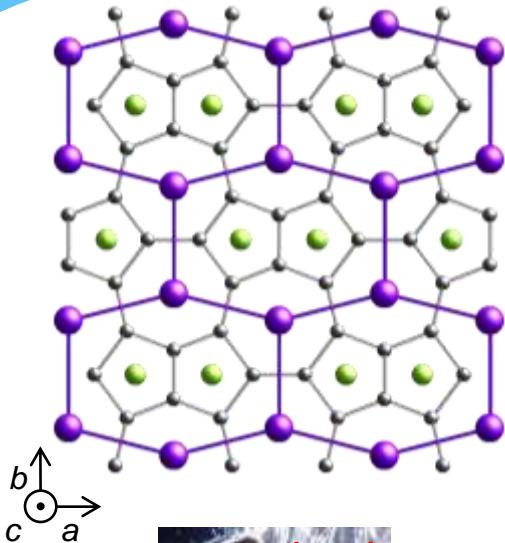


- Type-II SC in the clean limit ( $l = 1.2 \mu\text{m} > \xi = 0.06 \mu\text{m}$ ) with  $T_c = 80 \text{ mK}$
- RRR dependence of  $T_c \Rightarrow$  **unconventional, non-s-wave SC ?**
- Considerable amounts of volume fractions  $\rightarrow$  Bulk SC
- Anisotropic  $B_{c2}$  ( $B \parallel ab$ : orbital critical fields,  $B \parallel c$ : paramagnetic effect)
- $\partial B_{c2} / \partial T|_{T_c} = -2.6 \text{ T/K}$  for  $B \parallel ab \rightarrow m^*/m_0 \sim 180$

# HF Compounds: $\beta$ -, $\alpha$ -YbAlB<sub>4</sub>

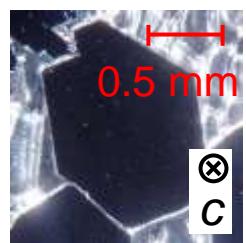
- Yb form distorted hexagonal lattice.
- Layered but Yb-Yb distance is shortest in c axis.

$\beta$ -YbAlB<sub>4</sub>



Yb    Al    B

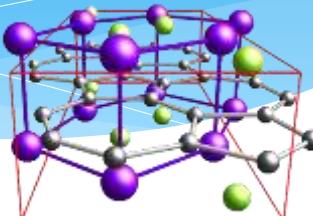
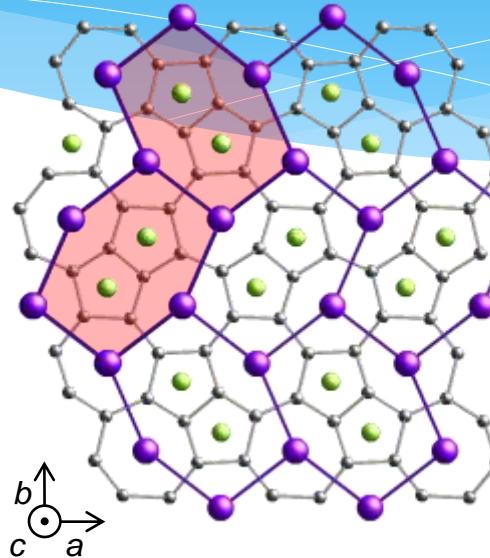
Space Group : *Cmmm*,  
 $a = 7.3080 \text{ \AA}$ ,  
 $b = 9.3150 \text{ \AA}$ ,  
 $c = 3.4980 \text{ \AA}$



Thickness:  $\sim 10 \mu\text{m}$

- The first HFSC ( $T_c = 80 \text{ mK}$ ) in Yb based system.
- Pronounced NFL behavior in the normal state @  $B = 0$ .

$\alpha$ -YbAlB<sub>4</sub>



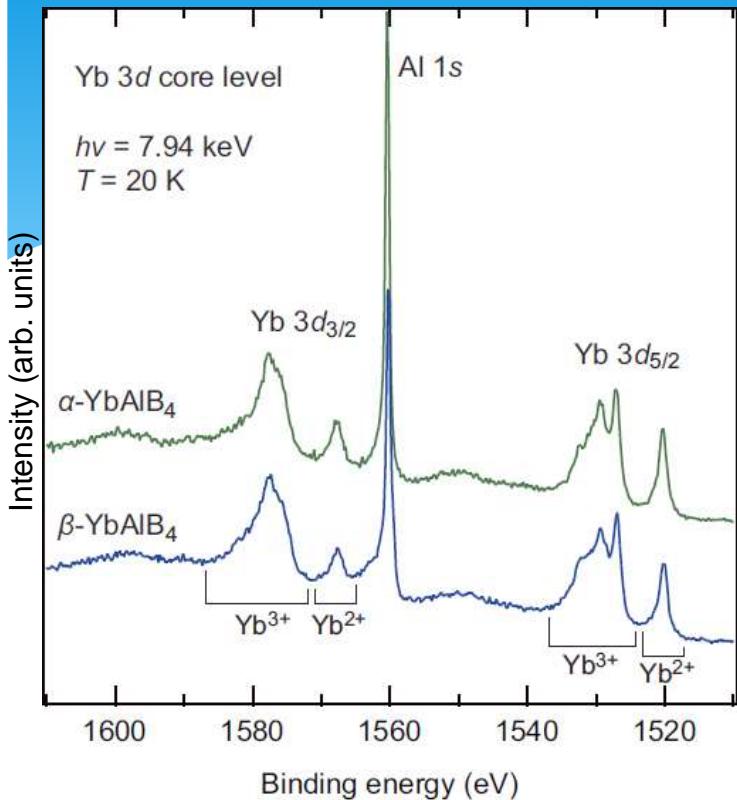
Space Group : *Pbam*,  
 $a = 5.9220(2) \text{ \AA}$ ,  
 $b = 11.468(5) \text{ \AA}$ ,  
 $c = 3.5060(5) \text{ \AA}$



$\sim 1 \times 1 \times 3 \text{ mm}^3$

- FL ground state @  $B = 0$
- Close to QC ?
- Metamagnetic behavior @  $B \sim 3 \text{ T}$

# Strong Valence Fluctuation in $\alpha$ , $\beta$ -YbAlB<sub>4</sub>



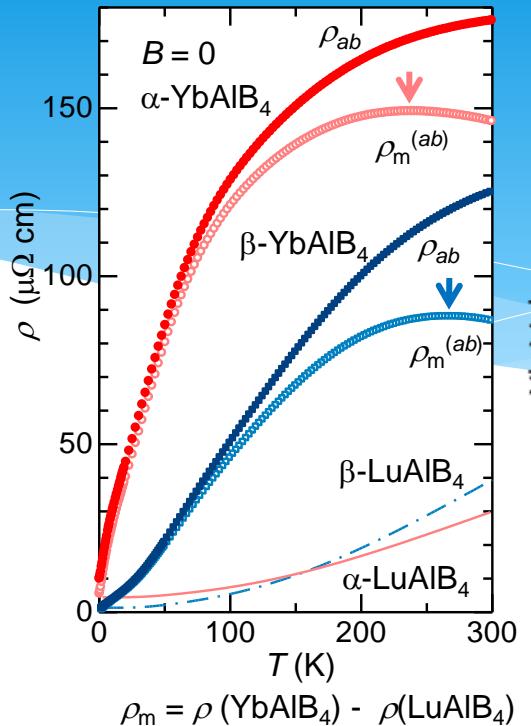
M. Okawa *et al.*, PRL **104**, 247201 (2010).

$\alpha\text{-YbAlB}_4 : 2.73 \pm 0.02$

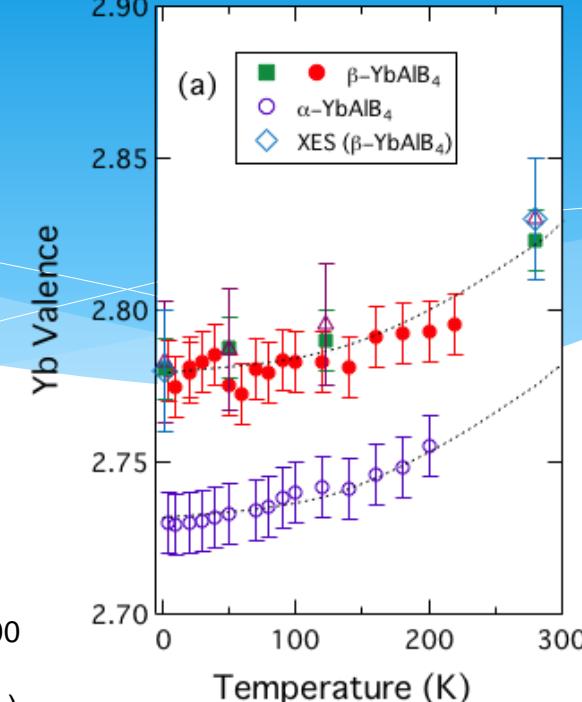
$\beta\text{-YbAlB}_4 : 2.75 \pm 0.02$

from HXPES

$\beta$  phase is the first example of **QC** in mixed valence system.



Y. Matsumoto *et al.*, PRB **84**, 125126 (2011).



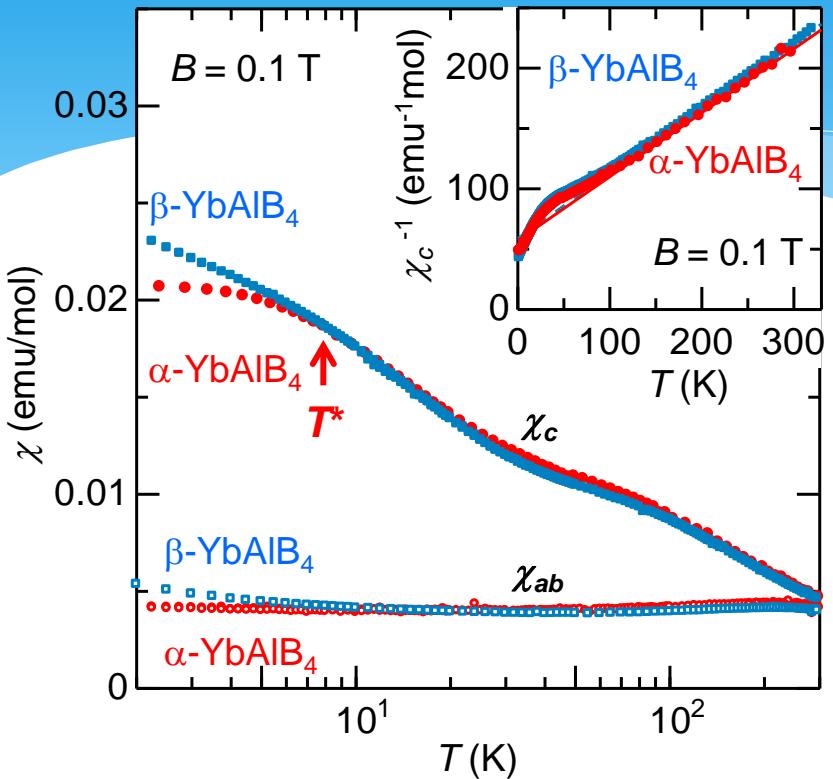
Y. H. Matsuda *et al.*, JKPS **62**, 1778 (2013).

Large  $T$  scale peculiar to mixed valence systems

- Coherence peak of a resistivity @  $T \sim 200 \sim 250 \text{ K}$
- $T_0 \sim 200 \text{ K}$  obtained from  $C$ .  
 $C_M/T = S_0/T_0 \ln(T_0/T)$
- Yb valence measured by X-ray adsorption spectra revealed  $T$  scale of  $\sim 290 \text{ K}$ .

# Kondo lattice behavior with small $T$ scale of $\sim 8$ K

Y. Matsumoto et al., Science 331, 316 (2011), Y. Matsumoto et al., PRB 84, 125126 (2011).

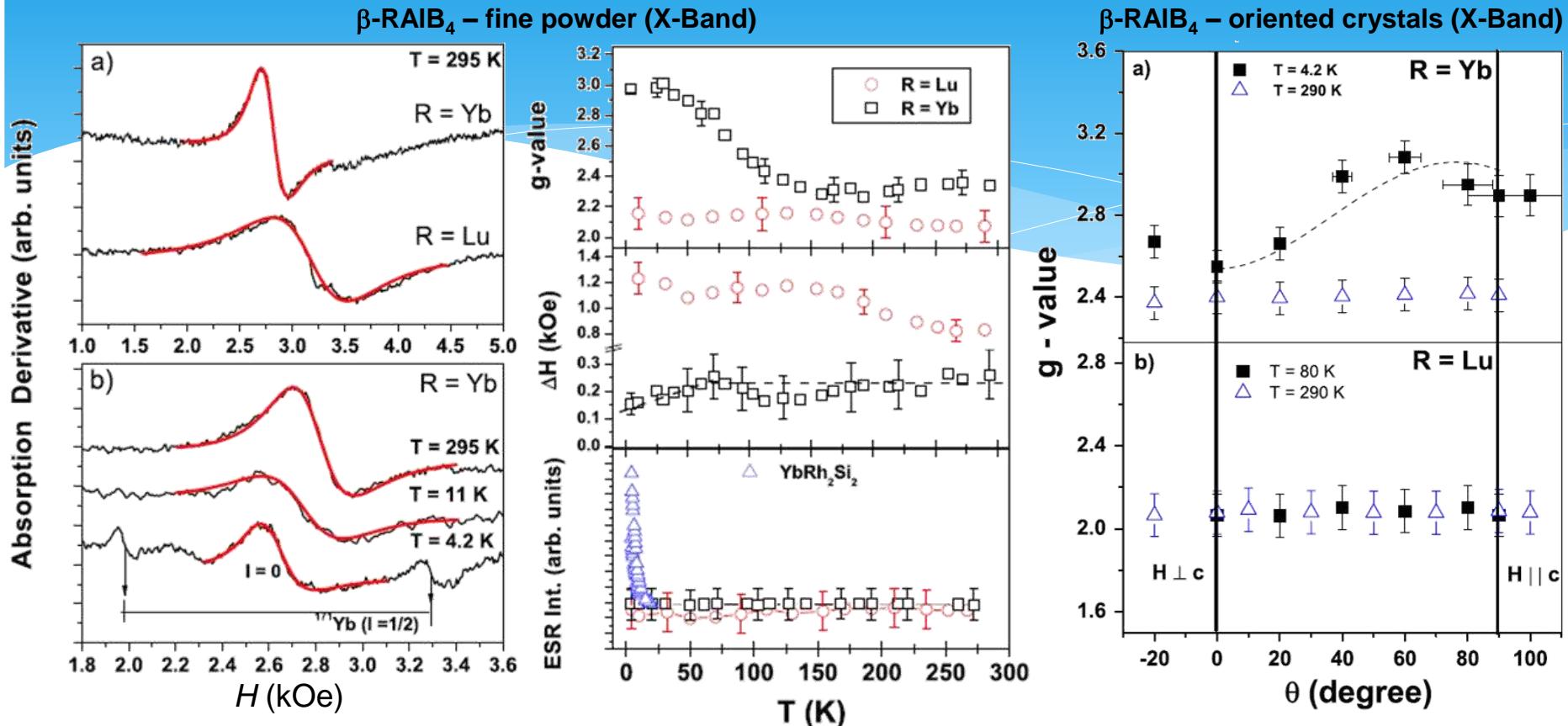


- Almost identical  $\chi$  for  $\alpha$  and  $\beta$  @  $T > 8$  K
- Ising like susceptibility with moments aligned along the c-axis
- **Curie-Weiss behavior in  $\chi_c$  far below the valence fluctuation  $T$  scale**
- **Shoulder in  $\chi_c$  @  $T^* \sim 8$  K**
  - ✓ Energy scale of Kondo lattice which emerges in low  $T$ .
  - ✓ Below  $T^* \sim 8$  K,  
 **$\alpha$  phase: FL  $\Leftrightarrow$   $\beta$  phase: NFL**
- **$C/T \sim 130$  mJ/K<sup>2</sup> mol @ 0.4 K**

How does the small  $T$  scale 8 K coexists with large valence  $T$  scale of 200 ~ 300 K ?

# Local moment behavior in ESR spectra of $\beta$ -YbAlB<sub>4</sub>

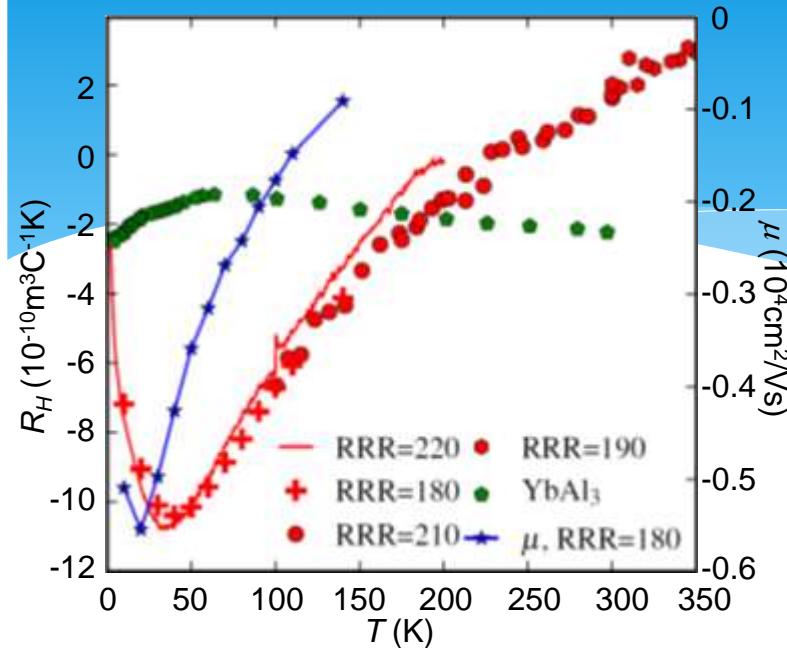
L. M. Holanda *et al.*, PRL 107, 026402 (2011).



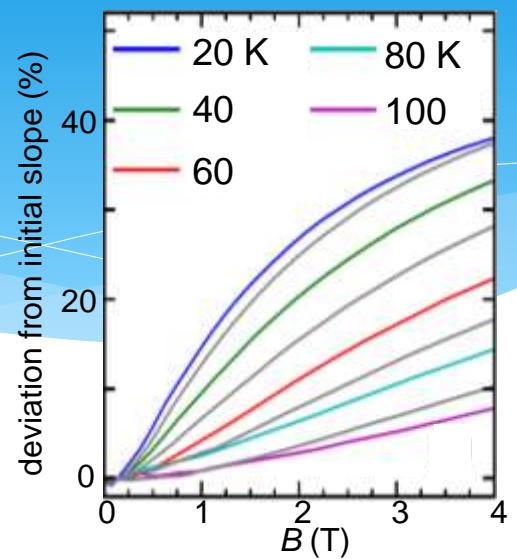
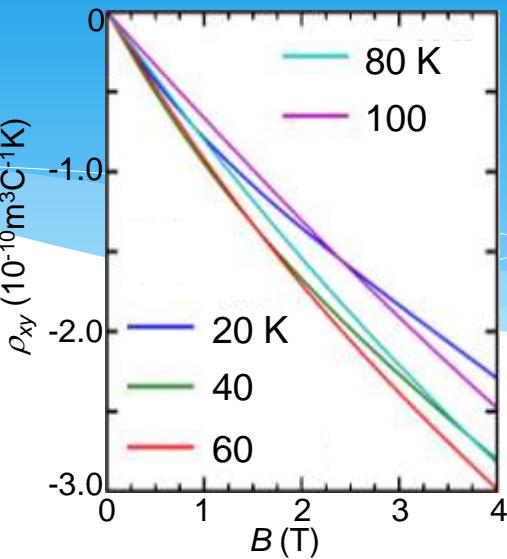
- $^{171}\text{Yb}$  hyperfine lines.
- $g$  value increases from 2.4 to 3.0 below  $T \sim 100 \text{ K}$ .
- anisotropy in  $g$  value
- $T$  independent intensity which is typical in ESR signal for conduction electron.

Unique dual behavior

# Two component Hall effect

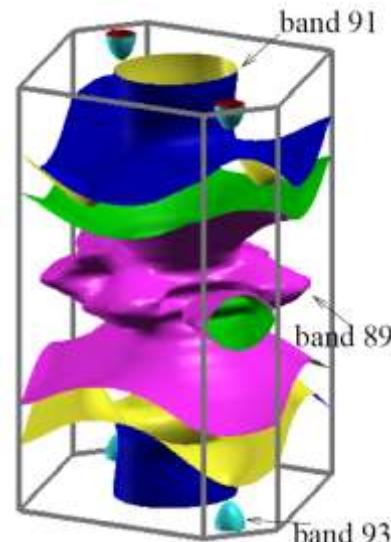


E. C. T. O'Farrell, Y. Matsumoto, S. Nakatsuji, PRL **109**, 176405 (2012).



- Large  $T$  dependence of  $R_H$
- $R_H$  has a peak at  $T \sim 40$  K ⇒ **Coherence develops below 40 K.**
- $\rho_{xy}$  becomes non-linear in  $B$  below 100 K. This can be explained by **two band model** with
  1. **hole-like component** (band 89?) and
  2. **electron-like component with low carrier density** (band 91?).

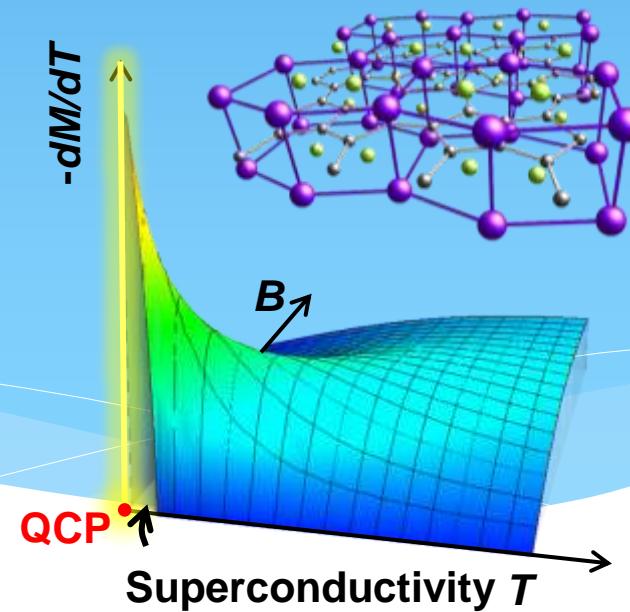
**The electron-like component is responsible for the Kondo lattice behavior and QC ?**



A. Nevidomsky and P. Coleman, PRL **102**, 077202 (2009). Also by E. C. T. O'Farrell et al., H. Harima et al., and W. Pickett et al.

# Quantum criticality without tuning in the mixed valence system $\beta\text{-YbAlB}_4$

Y. Matsumoto, S. Nakatsuji, K. Kuga, Y. Karaki, N. Horie, Y. Shimura, T. Sakakibara,  
A. H. Nevidomskyy, P. Coleman, Science **331**, 316 (2011).



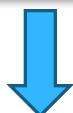
# NFL behavior @ $B = 0$ in $\beta\text{-YbAlB}_4$

S. Nakatsuji et al., Nature Phys. 4, 603-607 (2008).

$B = 0$ :

Non Fermi Liquid

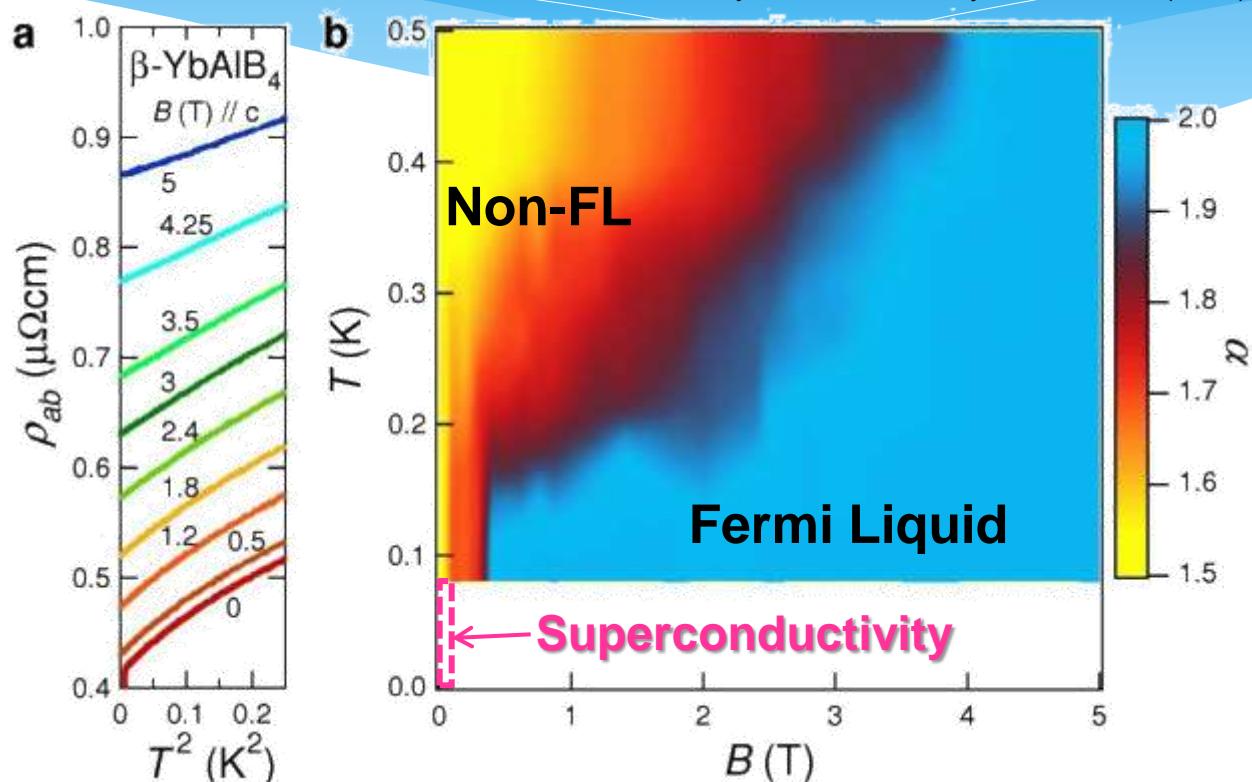
$$\left\{ \begin{array}{l} \rho_{ab} \propto T^{1.5} \\ M/H \propto T^{-0.5} \\ C/T \propto \ln(T^*/T) \end{array} \right.$$



In magnetic field:

Fermi Liquid

$$\left\{ \begin{array}{l} \rho_{ab} \propto T^2 \\ M/H = \text{const.} \\ C/T = \text{const.} \end{array} \right.$$



$$\Delta\rho = \rho(T) - \rho(0) = T^\alpha$$

FL behavior is recovered in magnetic field.

**QCP at  $B = 0$  under  $P = 0$  ?**

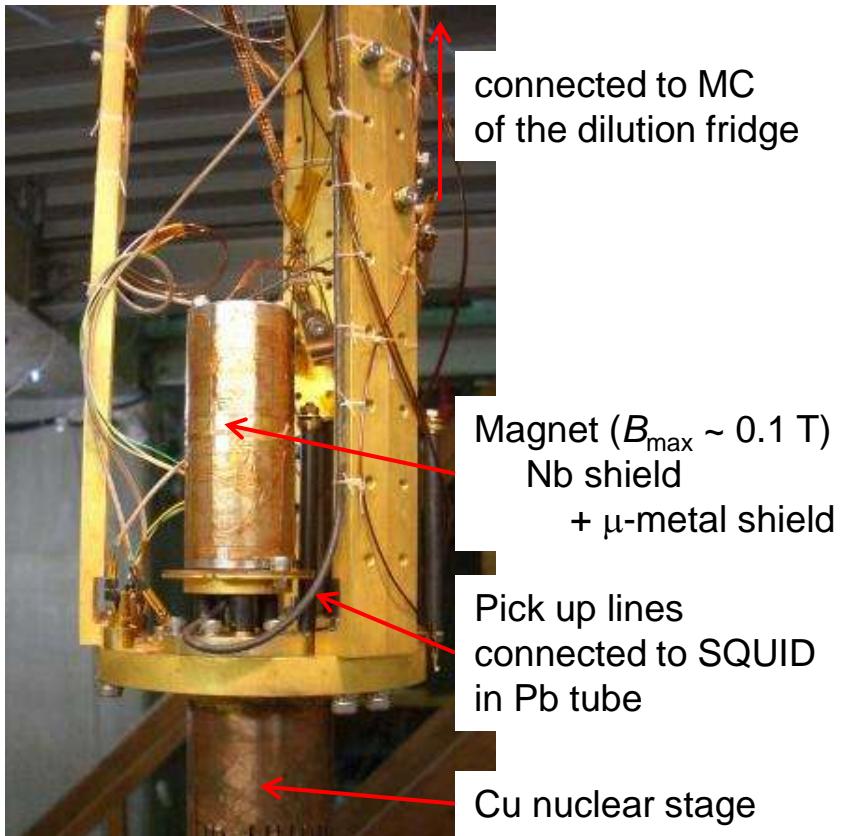
Normally, we have to tune  $B$  or  $P$  or doping to approach QCP.

ex)  $\text{YbRh}_2\text{Si}_2$ ,  $\text{CeCoIn}_5$ ,  $\text{CeCu}_{5.9}\text{Au}_{0.1}$ , ...

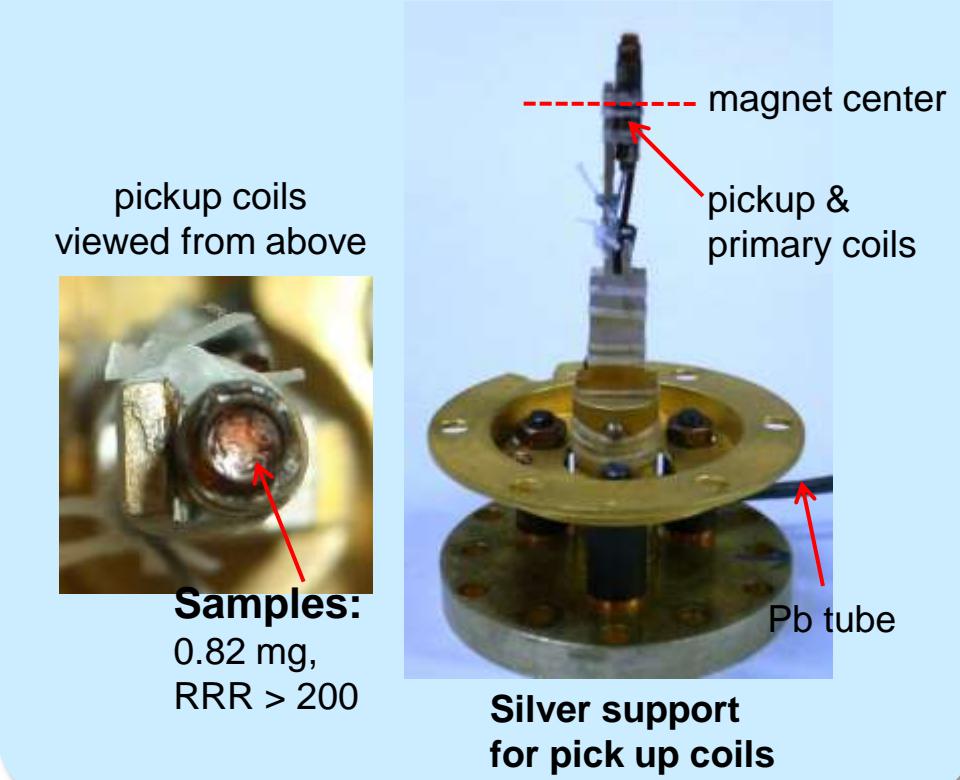
# SQUID magnetometer installed in a dilution fridge

ULT group @ ISSP, Univ. of Tokyo

- High precision (**resolution:  $\sim 10^{-8}$  emu**)
- $T$  range: **20 mK  $\leq T \leq 4$  K,  $B < 0.1$  T**  
 $T$  down to  $\sim 0.1$  mK is also available by using nuclear demag
- AC susceptibility and DC magnetization can be measured.

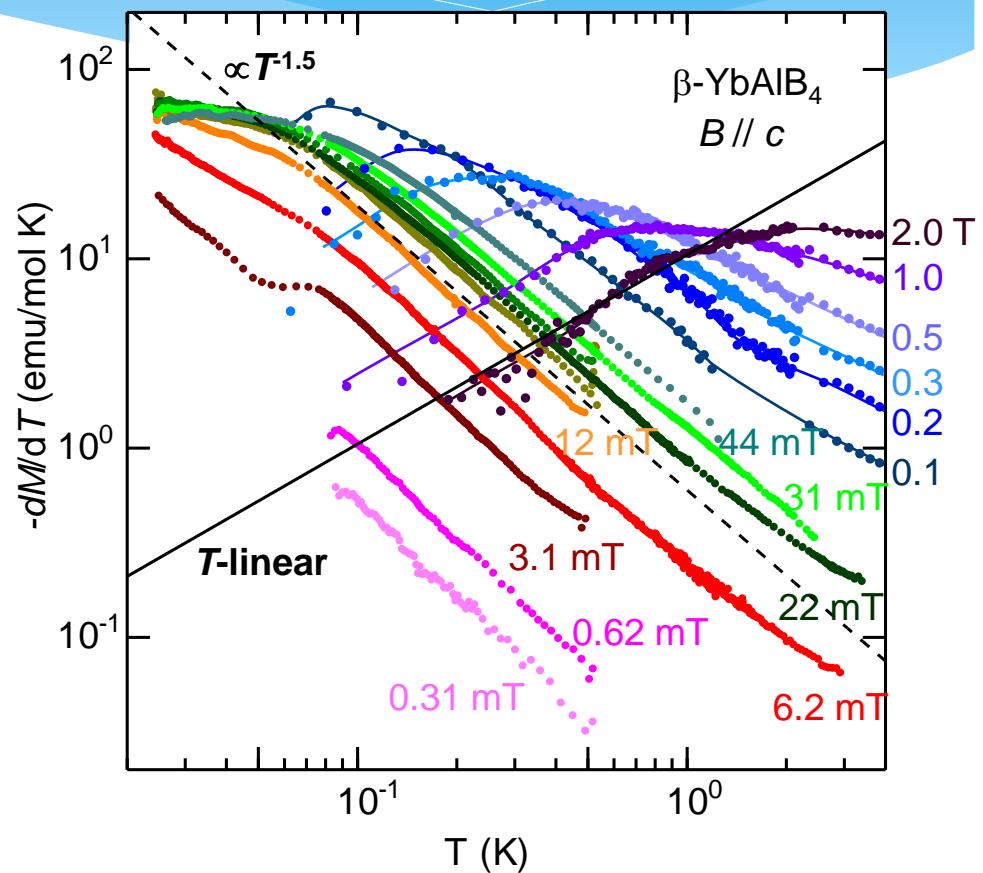
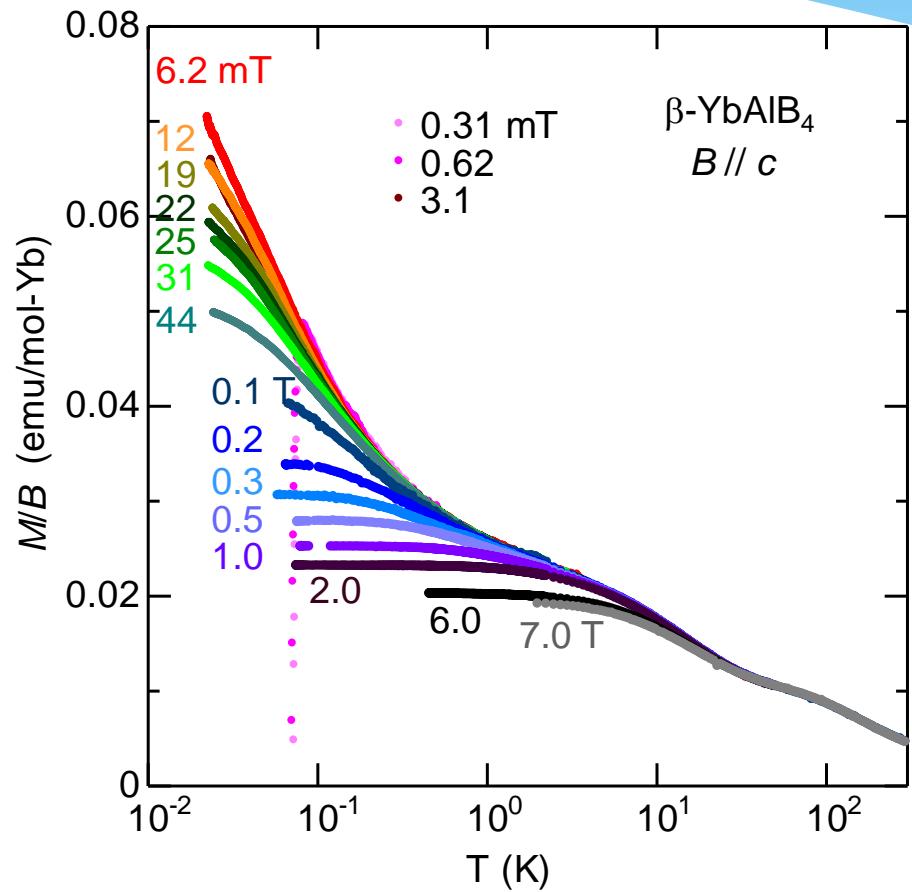


Setup inside of the shielded magnet



# Diverging Susceptibility at $B \rightarrow 0$ ( $B \parallel c$ )

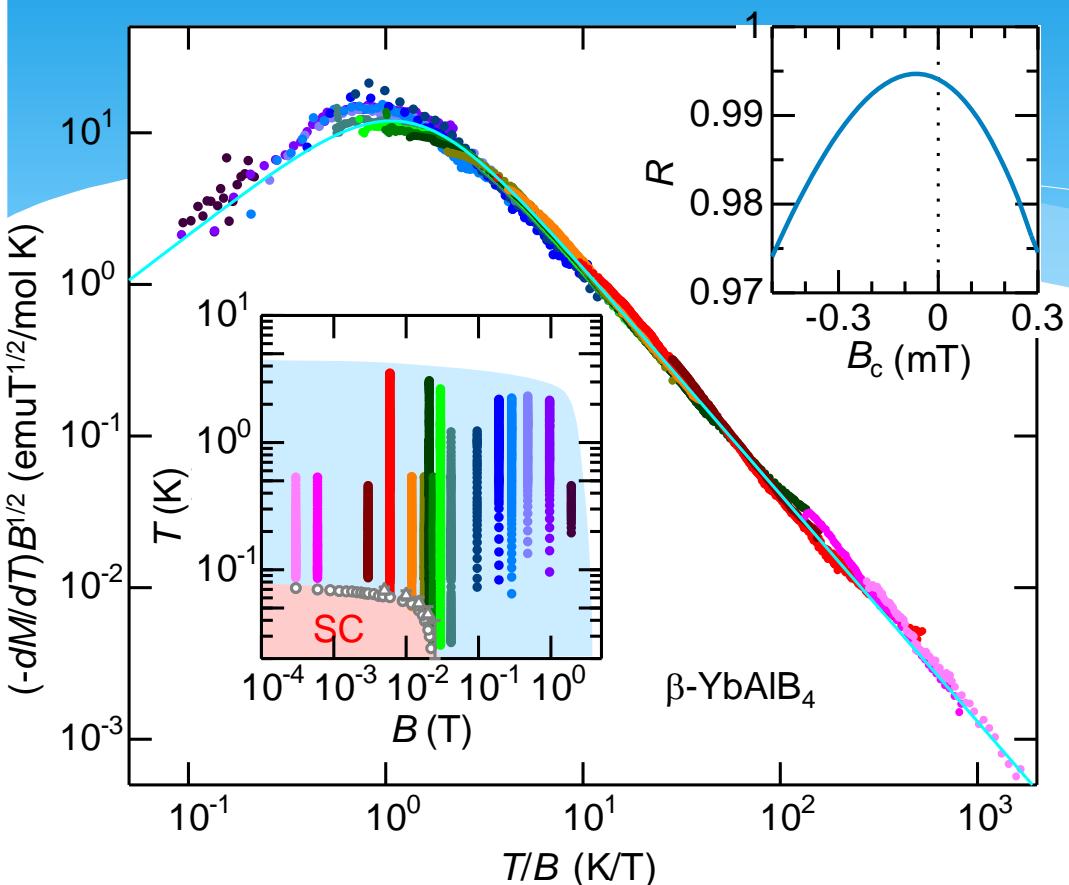
Y. Matsumoto et al., Science 331, 316 (2011).



- At  $T/B \gg 1$  (but below  $T^* \sim 8$  K) (NFL):  $-dM/dT \propto T^{-1.5}$  ( $M \propto \text{const.} + T^{-0.5}$ )
- At  $T/B \ll 1$  (FL):  $-dM/dT \propto T$

# T/B Scaling of Magnetization

Y. Matsumoto et al., Science 331, 316 (2011).



- dM/dT @  $B \leq 2$  T,  $T \leq 3$  K satisfy following scaling equations.

$$-\frac{\partial M}{\partial T} = B^{\alpha-2} \phi(T/B)$$

with  $\alpha = 3/2$

Integrating both parts,

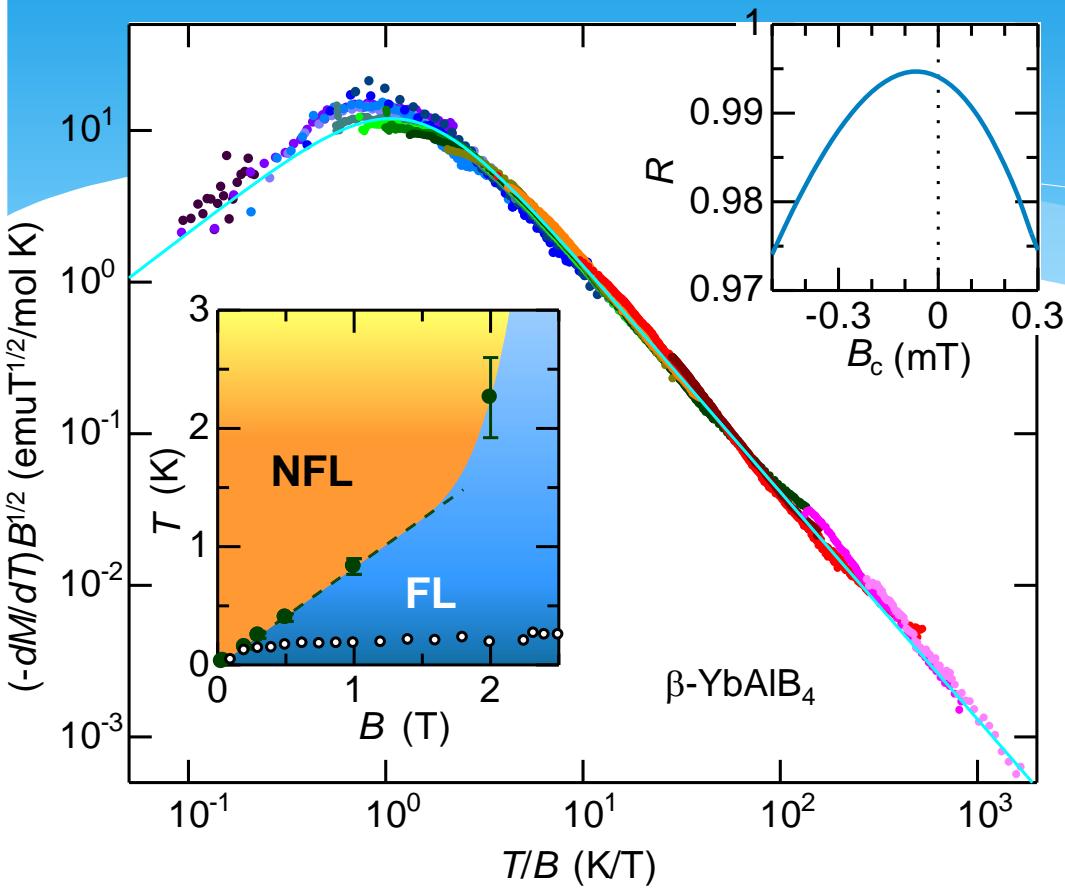
$$F = B^\alpha f(T/B)$$

$$\phi(x) = (\alpha-1)f'(x) - xf''(x)$$

- Scaling observed over 4 orders of  $T/B$
- No intrinsic energy scale ( $T$  and  $B$  are interchangeable) @  $B \leq 2$  T,  $T \leq 3$  K.
- QCP is not of a SDW type

# T/B Scaling of Magnetization

Y. Matsumoto et al., Science 331, 316 (2011).



➤ The data can be fitted to

$$\phi(x) = \Lambda x (A + x^2)^{-n}$$

with  $x = T / (B - B_c)$

This satisfy the limiting behavior of  $f(x)$ ,  
i.e.,  $f(x) \propto \begin{cases} x^\alpha & (x \gg 1, \text{NFL}) \\ \text{const} + x^2 & (x \ll 1, \text{FL}) \end{cases}$

The best fit was obtained with

$$\begin{cases} n = 1.25 \pm 0.01 \\ (\alpha = 3/2 \pm 0.02) \\ B_c = 0.1 \pm 0.1 \text{ mT} \end{cases}$$

➤  $|B_c| \leq 0.2 \text{ mT (2 G)}$

- ✓  $B_c$  is practically zero
- ✓  $\mu_0 H_{c2} / B_c > 100$
- ✓ Zero-field Quantum Criticality

The first example of  
quantum criticality without tuning !!

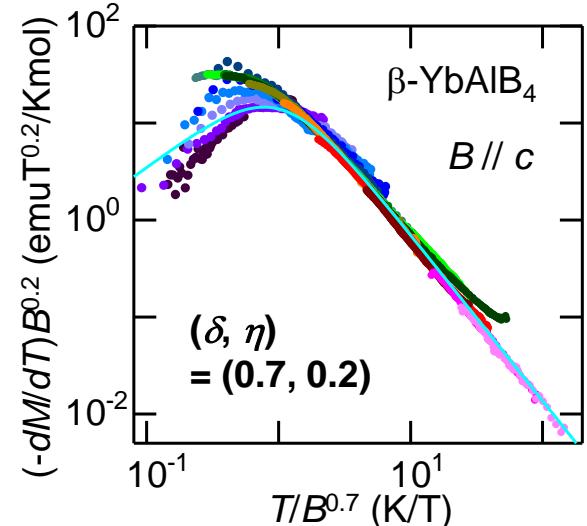
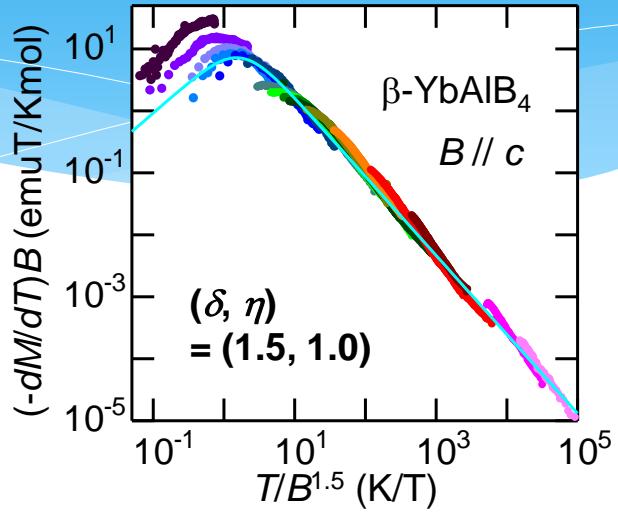
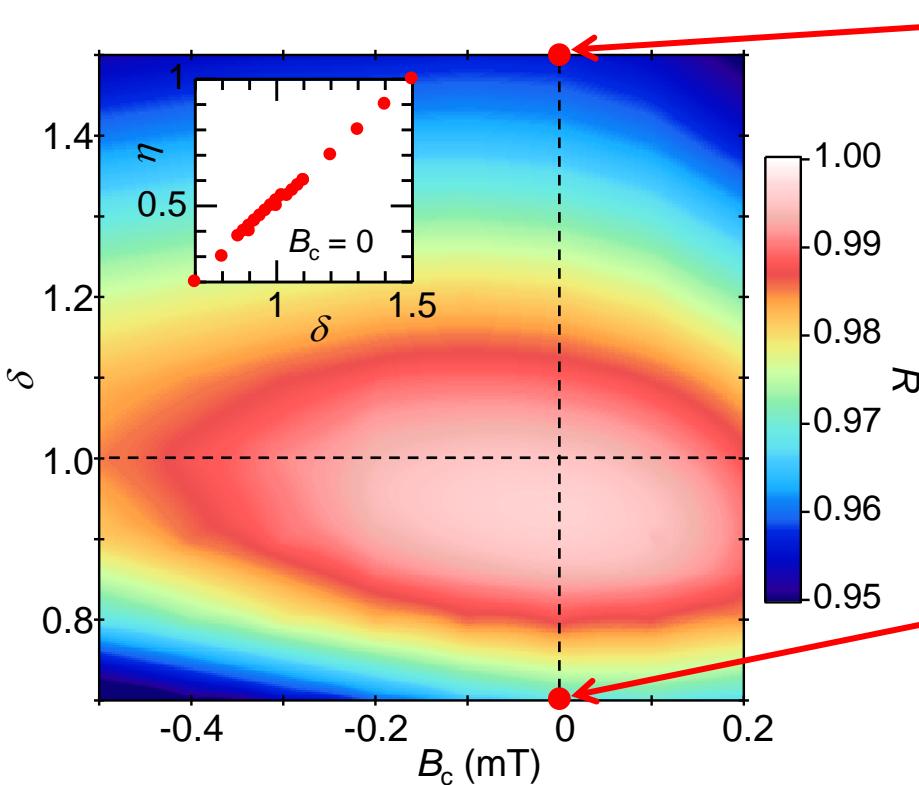
# $TI(B-B_c)^\delta$ scaling with $\delta \neq 1$ , $B_c \neq 0$ ?

More general form was checked.

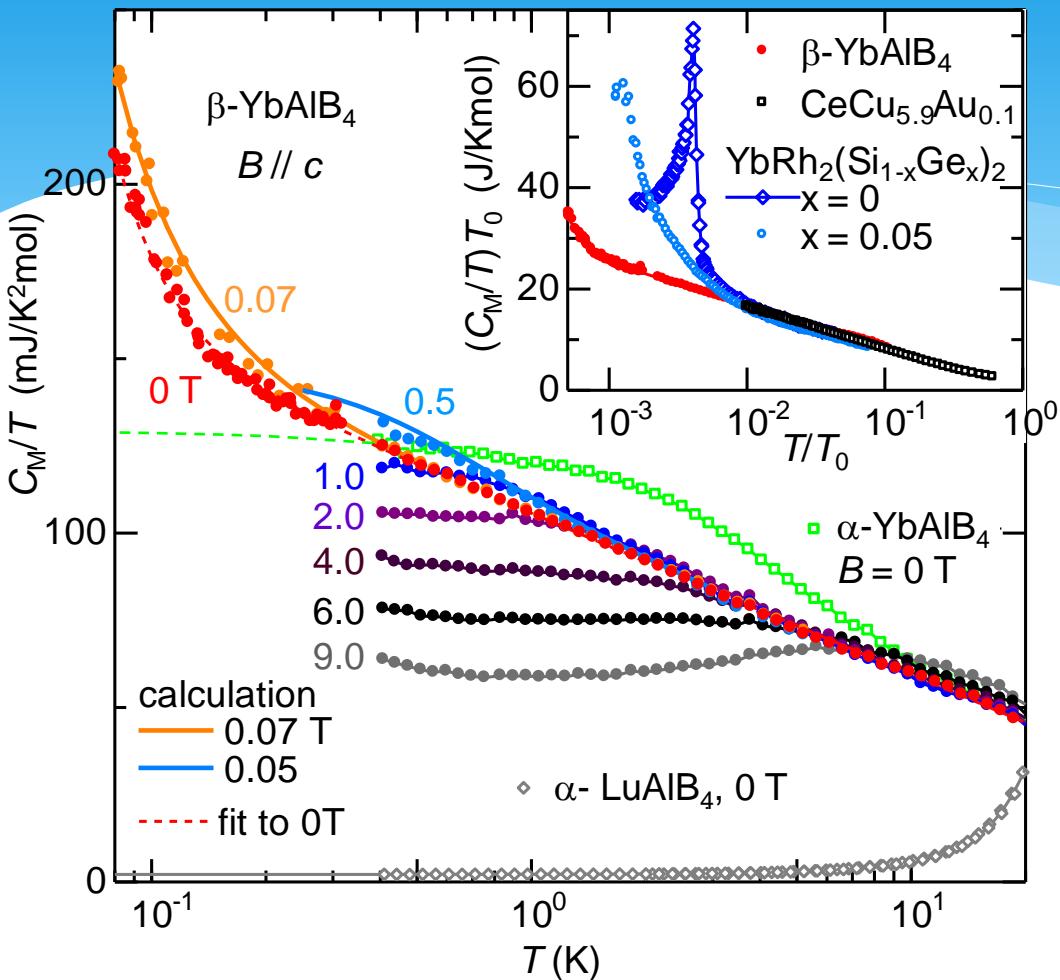
$$-\frac{dM}{dT} = (B - B_c)^{-\eta} \phi \left( \frac{T}{(B - B_c)^\delta} \right)$$

The best fit is obtained with  
 $(B_c, \delta) = (-0.02 \pm 0.20 \text{ mT}, 0.94 \pm 0.12)$

⇒  **$T/B$  scaling is most likely!**



# Specific Heat



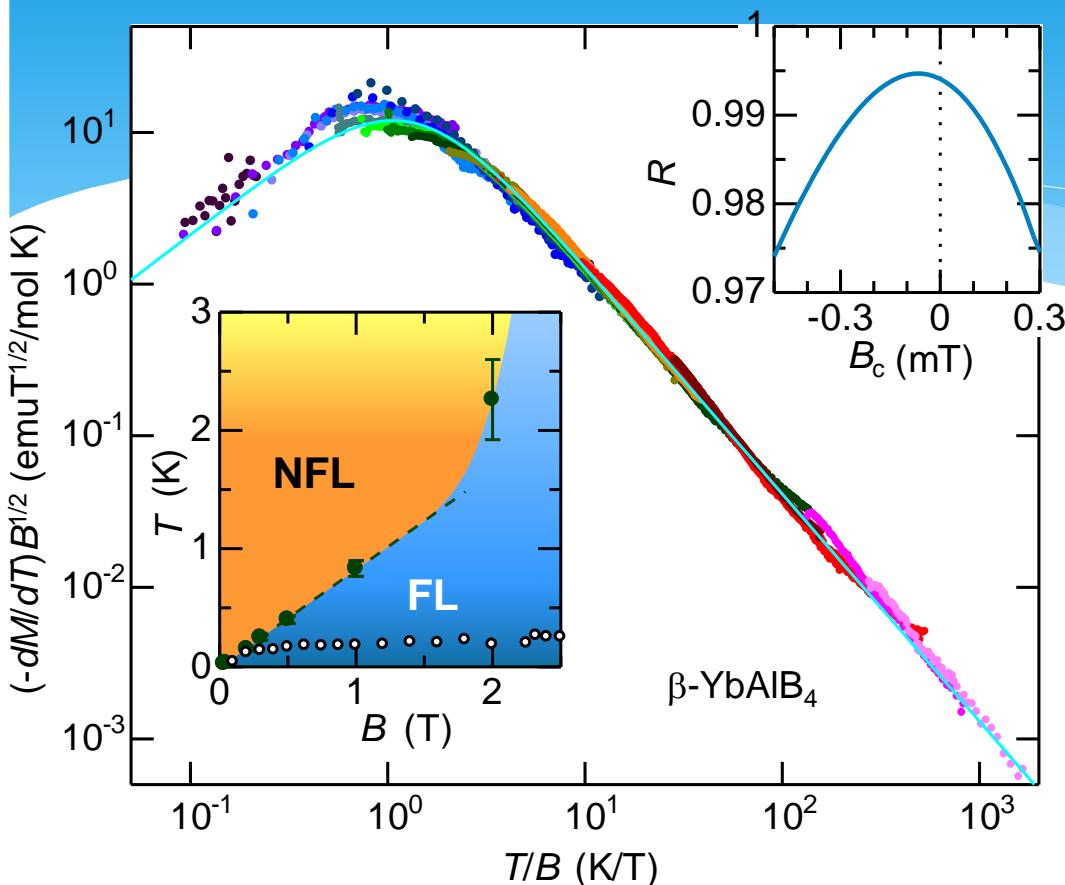
Y. Matsumoto *et al.*, Science 331, 316 (2011).

- $C_M/T \sim \ln(T_0/T)$  @  $T \geq 0.2\text{ K}$   
with  $T_0 = 200\text{ K}$   
 $\text{CeCu}_{5.9}\text{Au}_{0.1}: 6.2\text{ K}$   
 $\text{YbRh}_2\text{Si}_2: 24\text{ K}$
- Further increase below 0.2 K  
 QC without tuning
- The scaling indicates  
 $C_{\text{QC}}/T \propto T^{-0.5}$ .  
 Difficult to extract due to  
the nuclear contribution.
- Small increase @  $B \leq 0.5\text{ T}$  is  
consistent with the scaling.  
 This can be checked by  
Maxwell's relations.

$$\left( \frac{\partial \mathbf{S}}{\partial \mathbf{B}} \right)_T = \left( \frac{\partial \mathbf{M}}{\partial T} \right)_B$$

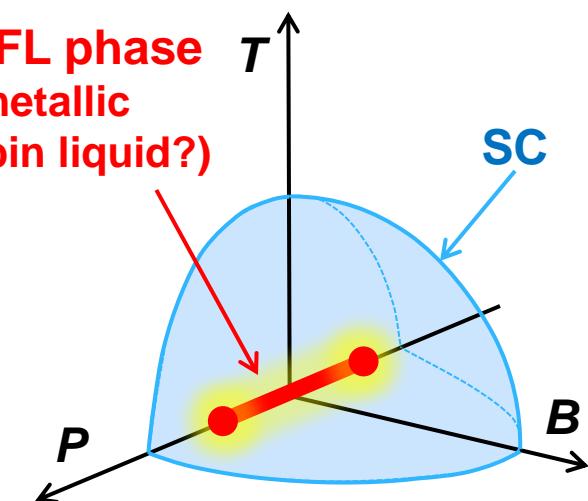
$$\Rightarrow \frac{\mathbf{C}}{T} = \int_0^B \frac{\partial^2 \mathbf{M}}{\partial T^2} d\mathbf{B} + \left( \frac{\mathbf{C}}{T} \right)_{B=0}$$

# Why $B_c = 0$ ?



Y. Matsumoto *et al.*, Science 331, 316 (2011).

1. Fine tuned to QCP?
2. **Non-fermi liquid phase or Quantum Critical Phase,**  
More natural than the above accidental scenario.

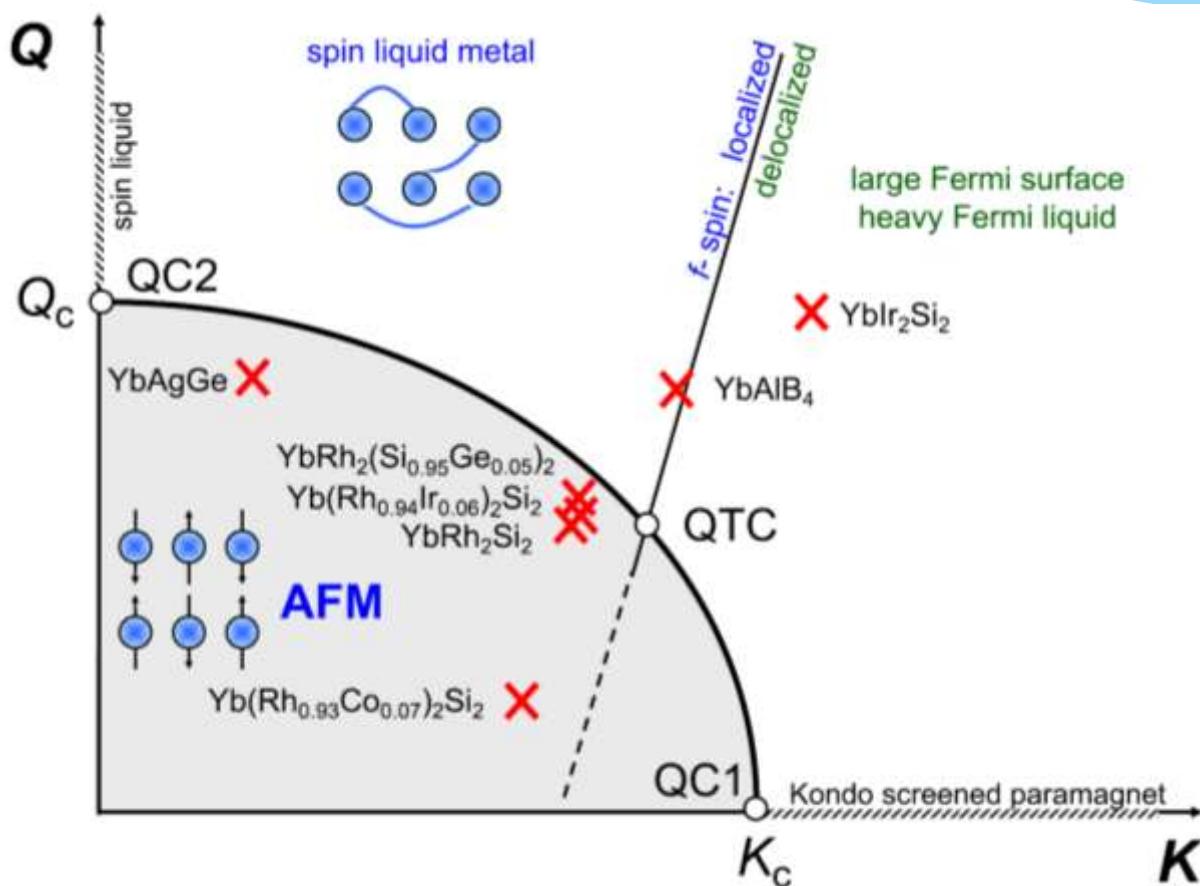


- What happens under pressure ?
- NFL behavior robust against pressure ?

# What is the origin ?

## Topological phase transition at $P_c$ ?

- Partial Mott localization of the  $f$ -electrons ?
- Spin liquid metal @ ambient pressure ?



J. Custers *et al.*, PRL 104, 186402 (2010).  
References therein.

# What is the origin ?

## Kondo break down QCP at mixed valence

J. H. Pixley, S. Kirchner, K. Ingersent, Q. Si, PRL **109**, 086403 (2010).

## Valence quantum criticality?

S. Watanabe, K. Miyake, PRL **105**, 186403 (2010).

## Anisotropic hybridization effect ?

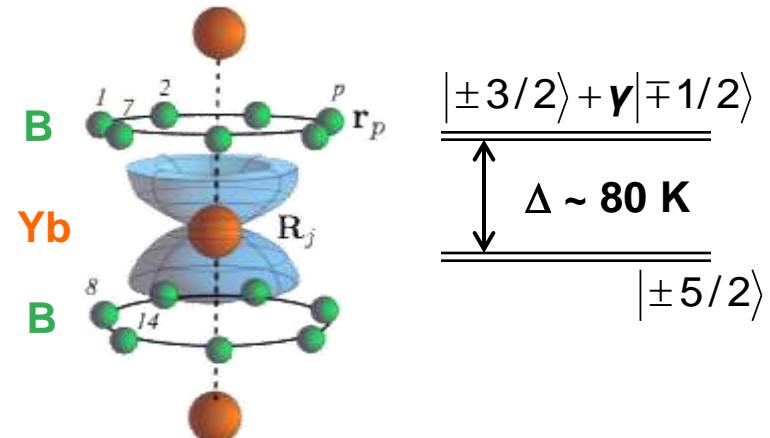
A. Nevidomsky and P. Coleman, PRL **102**, 077202 (2009), A. Ramires *et al.*, PRL **109**, 176404 (2012).

### ➤ A hybridization node along the c-axis

CEF ground doublet  $|J_z = \pm 5/2\rangle$   
based on the local Yb site symmetry

- $c-f$  hybridization  $V(\mathbf{k}) \sim (k_x \pm ik_y)^2$
- 2D heavy band with a dispersion  $E(k) \sim |V(k)|^2 \sim (k_{\perp})^4$  explains QC in  $\beta$ -YbAlB<sub>4</sub>, if  $E_f = 0$  ?

- Topological phase transition
- Highly anisotropic resistivity in  $\alpha$ -YbAlB<sub>4</sub> with  $\rho_{ab}/\rho_c \sim 11$  @ LT  
Consistent with hybridization node along the c-axis



Y. Matsumoto, K. Kuga, T. Tomita, Y. Karaki and S. Nakatsuji, PRB **84**, 125126 (2011).

# Summary

- In spite of the **strong valence fluctuation** with characteristic  $T$  scale of **200-300 K**, the both  $\alpha$ - and  $\beta$ -YbAlB<sub>4</sub> exhibit **Kondo lattice behavior** with low- $T$  scale of  **$T^* \sim 8$  K**.
- Contrasting behavior below  $T^* \sim 8$  K
  - $\alpha$  phase: FL**         **$\beta$  phase: QC without tuning + SC**
- **$T/B$  scaling of magnetization**  
observed over 4 decades of  $T/B$  in  
 $\beta$ -YbAlB<sub>4</sub> → **QC without tuning**
- Resistivity measurements under  $P$   
revealed **strange metal phase**  
**extends up to  $P_c \sim 0.4$  GPa**
- **Strange metal phase detached  
from magnetic order**
- What is the origin ?
  - ✓ topological phase transition?
  - ✓ valence? ...