

Unitary equivalences

$$\hat{\phi}_{X_\varepsilon} = \frac{\partial}{\partial t} \sigma_0 + e^t \frac{\partial}{\partial x} \sigma_1 + e^{-t} \frac{\partial}{\partial y} \sigma_2$$

$$\begin{aligned} \psi((x,y),t) &= \sum_{\lambda} \psi_{\lambda} e^{2\pi i \langle (a,b), \lambda \rangle} \\ &= \sum_{\lambda} \psi_{\lambda} e^{2\pi i \langle \Theta_{-t}(x,y), \lambda \rangle} = \sum_{\lambda} \psi_{\lambda} E_{\lambda} \end{aligned}$$

where

$$(x,y) = \Theta_{-t}(a,b) \quad t \in [0, \log \varepsilon)$$

↑ standard torus

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fours $T_t = \mathbb{R}^2 / \Lambda_t$

$$\Theta_{-t}(x,y) = (e^{-t}x, e^ty) \quad \langle (a,b), \lambda \rangle = a\lambda_1 + b\lambda_2$$

$$\hat{\phi}_{X_\varepsilon} E_{\lambda} = \left(\frac{\partial}{\partial t} \sigma_0 + 2\pi i \lambda_1 \sigma_1 + 2\pi i \lambda_2 \sigma_2 \right) E_{\lambda}$$

Thus, passing to Fourier modes in the fiber T^2 -direction gives unitary equivalence

$$\hat{\phi}_{X_\varepsilon} \psi_{\lambda} = \left(\frac{\partial}{\partial t} \sigma_0 + 2\pi i \lambda_1 \sigma_1 + 2\pi i \lambda_2 \sigma_2 \right) \psi_{\lambda}$$

with NC torus action

$$\hat{\pi}(R_{\eta}^{\sigma}) \psi_{\lambda} = \sigma(\eta, \lambda) \psi_{\lambda + \eta} \quad \text{in Connes-Landi deformation}$$

Then other unitary equivalence:

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$$U \psi_\lambda = \sigma_\lambda \psi_\lambda \quad \text{where}$$

$$\sigma_\lambda = \text{prod of Pauli matrices } \sigma_i \quad i=1,2$$

where σ_i is in σ_λ iff $\text{corresp. } \lambda_i < 0$

$$\Rightarrow U \hat{\phi}_{X_\varepsilon} U^* = \text{sign}(N(\lambda)) \left(\frac{\partial}{\partial t} \sigma_0 + 2\pi i |\lambda_1| \sigma_1 + 2\pi i |\lambda_2| \sigma_2 \right)$$

and with

$$U \hat{\pi}(R_\eta) U^* : \sigma_\lambda \psi_\lambda \mapsto \sigma_{\lambda+\eta} \psi_{\lambda+\eta}$$

Then one more unitary transformation:

$$F_V = \text{fundam domain of action of } V$$

on lattice Λ

$$\text{write then } \lambda \in \Lambda \text{ as } \lambda = A_\varepsilon^k(\mu) \quad \begin{array}{l} \text{Some } \mu \in F_V \\ \text{Some } k \in \mathbb{Z} \end{array}$$

then for $\lambda \neq 0$ time shift

$$\tilde{U}(\sigma_\lambda \psi_\lambda)(t) = \sigma_\lambda \psi_\lambda \left(t - \log \frac{|\mu|}{|N(\mu)|^{1/2}} \right)$$

$$\Rightarrow \hat{\psi}_\lambda := \tilde{U}(\sigma_\lambda \psi_\lambda) = \sigma_\lambda \psi_{|N(\lambda)|^{1/2}} (\text{sign}(\lambda_1) \varepsilon^k, \text{sign}(\lambda_2) \varepsilon^{-k})$$

This gives a unitarily equivalent spectral triple (3)

$$\tilde{\mathcal{D}} = \tilde{\mathcal{D}}^{(0)} + \sum_{\mu \in \Lambda \setminus \{0\}} \tilde{\mathcal{D}}^\mu \quad \text{where}$$

$$\tilde{\mathcal{D}}^\mu \tilde{\psi}_{A_\varepsilon^k(\mu)} \approx \text{sign}(N(\mu)) |N(\mu)|^{1/2}$$

$$\left(|N(\mu)|^{-1/2} \frac{\partial}{\partial t} \sigma_0 + 2\pi i \varepsilon^k \sigma_1 + 2\pi i \varepsilon^{-k} \sigma_2 \right) \tilde{\psi}_{A_\varepsilon^k(\mu)}$$

$$\tilde{\pi}(R_\gamma^\sigma) \tilde{\psi}_\lambda = \tilde{\psi}_{\lambda+\gamma}$$

$$\tilde{\mathcal{D}}^\mu = D_\mu B_\mu \quad \text{where} \quad D_\mu \tilde{\psi}_{A_\varepsilon^k(\mu)} = \text{sign}(N(\mu)) |N(\mu)|^{1/2} \tilde{\psi}_{A_\varepsilon^k(\mu)}$$

$$B_\mu \tilde{\psi}_{A_\varepsilon^k(\mu)} = \left(|N(\mu)|^{-1/2} \frac{\partial}{\partial t} \sigma_0 + 2\pi i \varepsilon^k \sigma_1 + 2\pi i \varepsilon^{-k} \sigma_2 \right) \tilde{\psi}_{A_\varepsilon^k(\mu)}$$

Relation to NC torus

$$\psi_\lambda \mapsto 2\pi i (\lambda_1 \sigma_1 + \lambda_2 \sigma_2) \psi_\lambda \quad \text{is}$$

$$D_{\theta, \theta'} = \begin{pmatrix} 0 & \delta_{\theta'} - i\delta_\theta \\ \delta_\theta + i\delta_{\theta'} & 0 \end{pmatrix}$$

$$\delta_\theta \psi_{n,m} \mapsto \psi_{(n+m\theta)} \psi_{n,m}$$

$$\delta_{\theta'} \psi_{n,m} \mapsto \psi_{(n+m\theta')} \psi_{n,m}$$

Under same unitary equivalences

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$$D_{\theta, \theta', 0}^{\mu} = \sum_{\mu \in \Lambda_{\frac{1}{V}}} D_{\theta, \theta'}^{\mu}$$

$$D_{\theta, \theta'}^{\mu} \psi_{A_{\varepsilon}^k(\mu)} = (\lambda_1 \sigma_1 + \lambda_2 \sigma_2) \psi_{A_{\varepsilon}^k(\mu)}$$

becomes after unitary equivalences

$$\tilde{D}_{\theta, \theta'}^{\mu} \tilde{\psi}_{A_{\varepsilon}^k(\mu)} = \text{sign}(N(\mu)) |N(\mu)|^{1/2} (\varepsilon^k \sigma_1 + \varepsilon^{-k} \sigma_2) \tilde{\psi}_{A_{\varepsilon}^k(\mu)}$$

$$\tilde{D}_{\theta, \theta'}^{\mu} = D_{\theta}^{\mu} B_{\theta}$$

$$D_{\theta}^{\mu} \tilde{\psi}_{A_{\varepsilon}^k(\mu)} = \text{sign}(N(\mu)) |N(\mu)|^{1/2} \tilde{\psi}_{A_{\varepsilon}^k(\mu)}$$

$$B_{\theta} \tilde{\psi}_{A_{\varepsilon}^k(\mu)} = (\varepsilon^k \sigma_1 + \varepsilon^{-k} \sigma_2) \tilde{\psi}_{A_{\varepsilon}^k(\mu)}$$

$$\zeta_{D_{\theta, \theta'}}(s) = 2 Z_{\varepsilon}\left(\frac{s}{2}\right) \sum_{\mu} |N(\mu)|^{-s/2}$$

$$\zeta_{\tilde{D}_{\theta, \theta'}}(s) = 2 \left(\Lambda, V, \frac{s}{2} \right) Z_{\varepsilon}\left(\frac{s}{2}\right)$$

with $Z_{\varepsilon}(s) = \sum_k (\varepsilon^{2k} + \varepsilon^{-2k})^{-s/2}$

Lorentzian spectral triples

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(another look at NC tori w/ real multipl. and the Shimizu L-function)

$$\Lambda \subset \mathbb{R}^2 \text{ as before } \quad K = \mathbb{Q}(\sqrt{d}) \quad \alpha_i: K \hookrightarrow \mathbb{R}$$

$$V = \varepsilon^{\mathbb{Z}} \quad \lambda \mapsto A_{\varepsilon}^k(\lambda) = (\varepsilon^k \lambda_1, \varepsilon^{-k} \lambda_2)$$

$$x \in K \Rightarrow c(x) = x' \text{ Galois conjugate}$$

$$\lambda_2 = c(\lambda_1)$$

$$N(\lambda) = \lambda_1 \lambda_2 = (m + m\theta)(m + m\theta')$$

quadratic form but not p.s. def.

Note: think wave operator

$$\square = p_1^2 - p_2^2 \text{ instead of Laplace operator } p_1^2 + p_2^2 = \Delta$$

(in momentum variables)

$$\mathcal{D}_{\lambda} = \begin{pmatrix} 0 & \mathcal{D}_{\lambda}^+ \\ \mathcal{D}_{\lambda}^- & 0 \end{pmatrix} = \begin{pmatrix} 0 & \lambda_1 \\ \lambda_2 & 0 \end{pmatrix} \quad \text{first order (Dirac) factorization}$$

$$\mathcal{D}_{\lambda}^2 = \square_{\lambda} \quad \square_{\lambda} = \begin{pmatrix} N(\lambda) & 0 \\ 0 & N(\lambda) \end{pmatrix}$$

Assemble together these Fourier modes

$$\mathcal{H} = \ell^2(\Lambda) \oplus \ell^2(\Lambda) \quad \mathcal{D} e_{\pm, \lambda} = \mathcal{D}_{\lambda} e_{\pm, \lambda}$$

$$\mathcal{D} \gamma = -\gamma \mathcal{D} \text{ w/ grading } \gamma = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \text{ on } \mathcal{H}$$

and $[D, R_\eta^\sigma] e_{\pm, \lambda} = \sigma(\lambda, \eta) \eta_{\pm} e_{\pm, \eta + \lambda}$

where $\eta_1 = \eta_+$ & $\eta_2 = \eta_-$

So far properties similar to usual Dirac of spectral triples on NC for but other properties different

1) D not self-adjoint!

but invariant under a different involution

$D = c(D^\dagger)$ $c = \text{Galois involution on } K$

[So see this only if working on K not on \mathbb{R} or \mathbb{C}]

2) More serious problem: D has infinite multiplicities in the spectrum

(hence cannot have compact resolvent as for spectral triple) : typical problem in Lorentzian geom. infinite non-compact group of isometries (here V)

Idea: use Krein spaces instead of Hilbert spaces & Wick rotation to Hilbert spaces

here: variation on this idea: want to keep arithm. structure so not field/Krein spaces but analogous vect. spaces over K w/ bilin forms

Arithmetic Krein spaces:

$$c: K \rightarrow K \text{ Galois invol. } c: x \mapsto x'$$

$V = K$ -vector space

$T: V \rightarrow V$ is c -linear if

$$T(av + bw) = c(a)T(v) + c(b)T(w)$$

Lorentzian pairing on a K -vector space V

$$(\cdot, \cdot): V \times V \rightarrow K$$

c -lin. in first, linear in second var.

Krein K -space: $(V, (\cdot, \cdot))$ st.

$\exists \kappa: V \rightarrow V$ c -lin. involution st.

$(\kappa \cdot, \cdot)$ satisfies

$$(1) (\kappa \cdot, \cdot) = c(\cdot, \kappa \cdot)$$

$$(2) \forall v \neq 0 \text{ in } V \quad (\kappa v, v) \in K \text{ is totally positive}$$

soy then $(\kappa \cdot, \cdot)$ "pos-def." inn prod (in K -sense)

Krein K -adjoint, T^\dagger

$$(v, Tw) = (T^\dagger v, w)$$

The \mathbb{C} -lin. invol. is the Wick rotation \mathcal{K}

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Given $\mathcal{V}, \langle \cdot, \cdot \rangle = (\mathcal{K} \cdot, \cdot)$

\Rightarrow real Hilbert spaces

$$\mathcal{V}_{\alpha_i} = \mathcal{V} \otimes_{\alpha_i(\mathbb{K})} \mathbb{R}$$

$$\langle v, w \rangle = \frac{1}{2} \alpha_1 ((\mathcal{K}v, w) + (v, \mathcal{K}w))$$

$$= \frac{1}{2} \alpha_2 ((\mathcal{K}v, w) + (v, \mathcal{K}w))$$

gives inner prod

$$\langle v, w \rangle = \frac{1}{2} \alpha_1 (\mathcal{K}v, w) + \frac{1}{2} \alpha_2 (\mathcal{K}v, w)$$

positive definite in usual sense since $(\mathcal{K}v, v)$ totally pos.

T \mathbb{K} -lin op. on \mathcal{V}

$$M_{\alpha_i}(T) \geq -\infty$$

$$M_{\alpha_i}(T) = \inf_{(v, v)=1} \alpha_i(Tv, Tv)$$

Krein \mathbb{K} -triple $(\mathcal{A}, \mathcal{V}, \mathcal{D})$

\mathcal{A} invol. \mathbb{K} -algebra

\mathcal{V} Krein \mathbb{K} -space (\cdot, \cdot)

\mathcal{D} densely defined

\mathbb{K} -linear

$\mathcal{D}^\dagger = \mathcal{D}$ Krein self-adj.

the "bounded commutator" condition becomes

$$G = [\mathcal{D}, \pi(a)] \text{ has } M_{\alpha_i}(G) > -\infty$$

w/ repres. of \mathcal{A}

$$\pi(a^*) = \pi(a)^\dagger \text{ Krein adjoint}$$

$$M_{\alpha_i}(\pi(a)) > -\infty$$

Then a Lorentzian \mathbb{K} -spectral triple
 $(A, \mathcal{V}, \mathcal{D})$ is as above with

(1) \exists densely def. \mathbb{K} -linear $op.$

$$U: \mathcal{V} \rightarrow \mathcal{V} \quad (Uv, Uv) = (v, v)$$

$$U^\dagger = U^{-1} \text{ s.t.}$$

$$U^\dagger \mathcal{D} U = \mathcal{D}$$

(2) $C_{\alpha, U} = [\mathcal{D}, \pi_U(\alpha)]$ has $M_{\alpha_i}(C_{\alpha, U}) > -\infty$

$$\pi_U(\alpha) = U^\dagger \pi(\alpha) U$$

(3) U unbounded self-adjoint $U = U^*$
on associated real Hilbert space $\langle \cdot, \cdot \rangle$

(4) $(A, \mathcal{V}, \mathcal{D})$ satisfies "p-summability"

$$\sum_n |\langle U e_n, |\mathcal{D}|^2 U e_n \rangle|^{-s/2} < \infty \quad \forall s \geq p$$

on real Hilbert space $\overset{\text{b.n.s.}}{\text{a.n.}} \mathcal{D} \langle \cdot, \cdot \rangle$

(Krein isometries need not be bounded operators!)

Now back to our example $e_\lambda \quad \lambda \in \Lambda$
basis of \mathcal{V}_Λ over \mathbb{K}

$$(v, w) := \sum_\lambda c(a_\lambda) b_\lambda \quad \text{for } v = \sum a_\lambda e_\lambda; w = \sum b_\lambda e_\lambda$$

$$K(v) = \sum_\lambda c(a_\lambda) e_\lambda \quad \text{Krein involution}$$

$$\alpha_{\lambda} \langle v, v \rangle = \sum_{\lambda} \alpha_{\lambda} (a_{\lambda}^2)$$

where

$$\langle v, w \rangle = (\kappa v, w) = c(v, \kappa w)$$

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$K[\lambda]$ acting on $V_{\lambda} : R_{\lambda} e_{\eta} = e_{\lambda+\eta}$

$M_{\alpha_i}(R_{\lambda}) > -\infty$ & induce bounded op's on associated real Hilbert spaces

$$(R_{\lambda} v, R_{\lambda} w) = (v, w) \text{ Krein isometry}$$

analog of twisted group ring:

$$\omega \in K^{\times} \text{ with } N(\omega) = \omega \omega' = 1$$

$$\tilde{\omega}(\lambda, \eta) = \omega^{(m, m) \wedge (k, k)} \text{ cocycle}$$

$$R_{\lambda}^{\tilde{\omega}} e_{\eta} = \tilde{\omega}(\lambda, \eta) e_{\lambda+\eta}$$

but unbounded op's on associated Hilbert space

$$R_{\lambda}^{\tilde{\omega}} e_{\eta, \pm} = A_{\omega}^{(r, k) \wedge (n, m)} e_{\lambda+\eta, \pm}$$

Lorentzian Dirac operators

$$\mathcal{D}_K e_{\lambda, \pm} = \mathcal{D}_{K, \lambda} e_{\lambda, \pm} = \begin{pmatrix} 0 & \mathcal{D}_{\lambda}^+ \\ \mathcal{D}_{\lambda}^- & 0 \end{pmatrix} e_{\lambda, \pm} = \begin{pmatrix} 0 & \rho \\ c(\lambda) & 0 \end{pmatrix} e_{\lambda, \pm}$$

$$\lambda = (\alpha_1(\lambda), \alpha_2(\lambda))$$

$$\mathcal{D}_{\lambda} = \begin{pmatrix} 0 & \lambda_1 \\ \lambda_2 & 0 \end{pmatrix} \quad c(\mathcal{D}_{\lambda}) = \begin{pmatrix} 0 & \lambda_2 \\ \lambda_1 & 0 \end{pmatrix}$$

Then Krein triple with involution:

$$T_{\varepsilon} e_{\lambda, \pm} = \begin{pmatrix} \varepsilon^{\rho(\lambda)} & 0 \\ 0 & \varepsilon^{-\rho(\lambda)} \end{pmatrix} e_{\lambda, \pm} \text{ for } A_{\varepsilon}^{\rho(\lambda)}(\mu) = \lambda$$

$$U = T_{\varepsilon} J$$

$$J e_{\lambda, \pm} = e_{J(\lambda), \pm} \quad J(\lambda) = A_{\varepsilon}^{-k}(\mu) \text{ for } \lambda = A_{\varepsilon}^k(\mu)$$

Then

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$$\sum_{\lambda \neq 0} \langle U_\varepsilon e_{\lambda, \pm}, |D_{\mathbb{K}}^2| U_\varepsilon e_{\lambda, \pm} \rangle^{-s/2}$$

$$= \sum_{\lambda \neq 0} \left(\varepsilon^{2p(\lambda)} + \varepsilon^{-2p(\lambda)} \right)^{-s/2} |N(\lambda)|^{-s/2}$$

$$= \sum_{k \in \mathbb{Z}} \left(\varepsilon^{2k} + \varepsilon^{-2k} \right)^{-s/2} \sum_{\substack{\mu \in \Lambda \setminus \{0\} \\ \frac{\mu}{\sqrt{V}}} } |N(\mu)|^{-s/2} \quad \text{zeta function}$$

$$\rightarrow \bullet \quad Z_\varepsilon\left(\frac{s}{2}\right) L(\Lambda, V, \frac{s}{2}) = \sum_n \frac{\text{sgn}(\langle U_{en}, D^2 U_{en} \rangle)}{|\langle U_{en}, |D^2| U_{en} \rangle|^{-s/2}}$$

Similarly for η function $=: \eta_D(s)$ (Lorentzian)
