

# Motives of intersections of quadrics and the Feynman integral of the massive sunset graph

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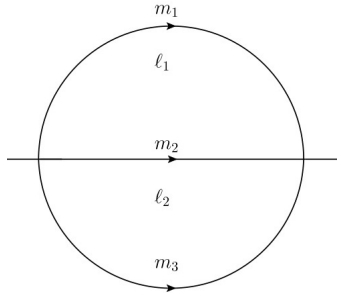
## This lecture is based on:

- Matilde Marcolli, Gonçalo Tabuada, *Feynman quadrics motive of the massive sunset graph*, arXiv:1705.10307

## Related work:

- M. Bernardara, G. Tabuada, *Chow groups of intersections of quadrics via homological projective duality and (Jacobians of) non-commutative motives*, Izv. Math. 80 (2016) no. 3, 463–480.
- M. Marcolli, G. Tabuada, *Jacobians of noncommutative motives*, Moscow Mathematical Journal 14 (2014) no. 3, 577–594.

## Sunset Graph



perturbative scalar QFT with masses  $m_i$  (massive propagators along the edges)

## Previously known (via different method)

- S. Bloch P. Vanhove, *The elliptic dilogarithm for the sunset graph*. J. Number Theory 148 (2015), 328–364
- S. Bloch, M. Kerr and P. Vanhove, *A Feynman integral via higher normal functions*. Compos. Math. 151 (2015) no. 12, 2329–2375.

The motive of the massive sunset graph (computed using graph hypersurfaces) in dimension  $D = 2$  and  $D = 4$  (case of equal masses and more general case) is non-mixed Tate (expressed in terms of elliptic curves).

We work in general dimension  $D$  and with Feynman quadrics instead of graph hypersurfaces

## Feynman graphs

- $D > 0$  spacetime dimension (Euclidean)
- $(\Gamma, m, \kappa)$  Feynman graph equipped with mass parameters  $m = (m_e)$  and external momenta  $\kappa = (\kappa_i)$
- internal edges  $e_j \in E_{\text{int}}(\Gamma)$  carry momentum variables  $k_j = (k_{j,r}) \in \mathbb{A}^D$
- edge propagator

$$q_j(k_j) = \sum_{r=1}^D k_{j,r}^2 + m_j^2$$

## Feynman integral $\mathcal{I}(\Gamma, m, \kappa)$

$$C \int \frac{\prod_{v \in V_{\text{int}}(\Gamma)} \delta(\sum_{e_i \in E_{\text{int}}(\Gamma)} \epsilon_{v,i} k_i + \sum_{e_j \in E_{\text{ext}}(\Gamma)} \epsilon_{v,j} \kappa_j)}{\prod_{e_i \in E_{\text{int}}(\Gamma)} q_i(k_i)} \prod_{e_i \in E_{\text{int}}(\Gamma)} \frac{d^D k_i}{(2\pi)^D}$$

- $C = \prod_v \lambda_v (2\pi)^{-D}$  with  $\lambda_v$  coupling constant at vertex  $v$
- $\epsilon_{v,i}$  incidence matrix with entries 1, -1, or 0, for  $v = s(e)$ ,  $v = t(e)$ ,  $v \notin \partial(e)$
- $\prod_{e_i} d^D k_i$  standard volume form in  $\mathbb{A}^{nD}(\mathbb{R})$
- $n := \#E_{\text{int}}(\Gamma)$  number of internal edges

Unrenormalized (usually divergent) Feynman integral

## Feynman quadrics

- Notation:  $v = (v_{i,r}) \in \mathbb{A}^{nD}$  and  $v' = (v'_{i,r}) \in \mathbb{A}^{nD}$ , let  $\langle v, v' \rangle := \sum_{i=1}^n \sum_{r=1}^D v_{i,r} v'_{i,r}$  and  $v^2 := \langle v, v \rangle = \sum_{i,r} v_{i,r}^2$
- Homogeneous polynomial ( $nD + 1$  variables):

$$q'_i(k_i, x) := \sum_{r=1}^D k_{i,r}^2 + m_i^2 x^2 = k_i^2 + m_i^2 x^2$$

(identify  $k_i = (k_{i,r}) \in \mathbb{A}^D$  with  $v = (v_{j,r})$  of  $\mathbb{A}^{nD}$  with  $k_{i,r}$  for  $i = j$  and 0 otherwise)

- Quadric Hypersurface:  $Q'_i \subset \mathbb{P}^{nD}$

## Linear relations

- delta function in Feynman integral imposes linear relations at vertices between the momentum variables

$$\sum_{\substack{e_i \in E_{\text{int}}(\Gamma) \\ s(e_i)=v}} k_i + \sum_{\substack{e_j \in E_{\text{ext}}(\Gamma) \\ s(e_j)=v}} \kappa_j = \sum_{\substack{e_i \in E_{\text{int}}(\Gamma) \\ t(e_i)=v}} k_i + \sum_{\substack{e_j \in E_{\text{ext}}(\Gamma) \\ t(j)=v}} \kappa_j$$

- $N$  number of *independent* linear relations
- choose  $n - N$  *independent* variables  $\ell = \{\ell_i\}$  among  $\{k_1, \dots, k_n\}$  (loop variables)
- have  $N = \#V_{\text{int}}(\Gamma) - 1$  so difference  $n - N = \#E_{\text{int}}(\Gamma) - \#V_{\text{int}}(\Gamma) + 1$  is first Betti number  $L = b_1(\Gamma)$



## Vanishing external momenta $\kappa = 0$

- linear subspace of momentum conservation

$$H_\Gamma := \bigcap_{v \in V_{\text{int}}(\Gamma)} \left\{ \sum_{\substack{e_i \in E_{\text{int}}(\Gamma) \\ s(e_i) = v}} k_i - \sum_{\substack{e_i \in E_{\text{int}}(\Gamma) \\ t(e_i) = v}} k_i = 0 \right\} \subset \mathbb{P}^{nD}$$

- Feynman quadrics in loop variables

$$Q_i := Q'_i \cap H_\Gamma = \{q_i(\ell, x) = 0\}$$

- The quadrics  $Q_i$  are usually singular (cones)
- Notation: coordinates  $u = (u_0 : \dots : u_{LD})$  on  $\mathbb{P}^{LD}$  with

$$u_0 := x, \quad (u_1, \dots, u_D) := \ell_1 \quad \dots \quad (u_{(L-1)D}, \dots, u_{LD}) := \ell_L$$

## Nets of quadrics

- parameterizing space of all quadric hypersurfaces in  $\mathbb{P}^{LD}$  is the projective space  $\mathbb{P}^{\binom{LD+2}{2}-1}$  of symmetric  $(LD+1) \times (LD+1)$ -matrices up to scalar multiples
- inside this parameterizing space discriminant hypersurface  $\mathcal{D}$ : quadratic forms with non-trivial kernel
- a *net of  $n$  quadric hypersurfaces* in  $\mathbb{P}^{LD}$  consists of an embedding  $\rho: \mathbb{P}^{n-1} \hookrightarrow \mathbb{P}^{\binom{LD+2}{2}-1}$

## Net of Feynman quadrics of a graph $\Gamma$ :

$$\rho: \mathbb{P}^{n-1} \hookrightarrow \mathbb{P}^{\binom{LD+2}{2}-1} \quad (0 : \dots : 0 : \underset{i}{1} : 0 : \dots : 0) \mapsto Q_i$$

- quadric hypersurfaces  $Q_i$  belong to  $\mathbb{P}^{\binom{LD+2}{2}-1}(\mathbb{R})$  (the defining quadratic form  $q_i$  of the quadric  $Q_i$  is real)
- symmetric matrices  $A_i$ , defined by  $q_i(u) = \langle u, A_i u \rangle$ , can be written as  $A_i = T_i^\dagger T_i$ , with  $T_i^\dagger$  adjoint of  $T_i$  with respect to the bilinear form  $\langle v, v' \rangle$
- momentum conservation condition:

$$\sum_{s(e_i)=v} \bar{T}_i = \sum_{t(e_i)=v} \bar{T}_i$$

$$\bar{T}_i = P T_i P \text{ with projection } P : (u_0, \dots, u_{LD}) \mapsto (u_1, \dots, u_{LD})$$

## Deformations of nets of quadrics

- one-parameter deformation of  $\rho : \mathbb{P}^{n-1} \hookrightarrow \mathbb{P}^{\binom{LD+2}{2}-1}$  is a morphism  $\tilde{\rho} : \mathbb{P}^{n-1} \times \mathbb{A}^1 \rightarrow \mathbb{P}^{\binom{LD+2}{2}-1}$  with  $\rho = \tilde{\rho}|_{\mathbb{P}^{n-1} \times \{0\}}$
- given  $\epsilon \in \mathbb{A}^1(\mathbb{Q})$ ,  $\epsilon \neq 0$ , write  $\rho_\epsilon$  for net  $\tilde{\rho}|_{\mathbb{P}^{n-1} \times \{\epsilon\}}$  ( $\epsilon$ -deformation of  $\rho$ )

For any Feynman graph  $\Gamma$  there is always a one-parameter deformation  $\tilde{\rho}$  of the net of Feynman quadrics

$$\rho_\epsilon : \mathbb{P}^{n-1} \hookrightarrow \mathbb{P}^{\binom{LD+2}{2}-1} \quad (0 : \dots : 0 : \underset{i}{1} : 0 : \dots : 0) \mapsto Q_{i,\epsilon}$$

such that for sufficiently small  $\epsilon$

- the quadrics  $Q_{i,\epsilon}$  belong to  $\mathbb{P}^{\binom{LD+2}{2}-1} \setminus \mathcal{D}$  (smooth)
- the quadrics  $Q_{i,\epsilon}$  belong to  $\mathbb{P}^{\binom{LD+2}{2}-1}(\mathbb{R})$  (real)
- the symmetric matrices  $A_{i,\epsilon}$  can be written as  $A_{i,\epsilon} = T_{i,\epsilon}^\dagger T_{i,\epsilon}$  (positive)
- momentum conservation condition:

$$\sum_{s(e_i)=v} \bar{T}_{i,\epsilon} = \sum_{t(e_i)=v} \bar{T}_{i,\epsilon} \text{ for } \bar{T}_{i,\epsilon} = PT_{i,\epsilon}P$$

These deformations  $\rho_\epsilon$  maintain physical properties (real, positive, momentum conservation) while they gain smoothness (replacing cones with smooth quadrics)

- Sketch of the argument for momentum conservation:
  - choose a spanning tree for the Feynman graph  $\Gamma$
  - constructing an  $\epsilon$ -deformation  $q_{i,\epsilon}$  of quadratic forms  $q_i$  associated to the  $L$  edges in the complement of the spanning tree
  - show that there is a unique way to extend the deformation to the remaining quadratic forms  $q_i$  on the edges of the spanning tree so that momentum conservation holds

## The Motive and the Period

Feynman integral in terms of periods of motive associated to the net of Feynman quadrics

- **motive**: in the category of mixed motives  $DM_{\text{gm}}(F)_{\mathbb{Q}}$  with  $F \subseteq \mathbb{C}$  algebraically closed

$$M_{(\Gamma, m)}^{\mathbb{Q}} = M(\mathbb{P}^{LD} \setminus Q_{(\Gamma, m)})_{\mathbb{Q}}$$

$Q_{(\Gamma, m)} := \bigcup_{i=1}^n Q_{i, \epsilon}$  union of the quadric hypersurfaces

- **algebraic differential forms**:  $\alpha \in \mathbb{N}$

$$\omega := \sum_{i=0}^{LD} (-1)^i u_i du_1 \wedge \cdots \wedge \widehat{du}_i \wedge \cdots \wedge du_{LD}$$

$$\eta_{\alpha} := \frac{\omega}{\prod_{i=1}^n q_i^{\alpha}} \quad \eta_{\alpha, \epsilon} := \frac{\omega}{\prod_{i=1}^n q_{i, \epsilon}^{\alpha}}$$

- restriction of  $\omega$  to the affine chart  $\mathbb{A}^{LD}$  coords  $(1, u_1, \dots, u_{LD})$  is affine volume form  $du_1 \wedge \cdots \wedge du_{LD}$

## Properties:

- $\alpha = 1$  (divergent) Feynman integral

$$\frac{C}{(2\pi)^D} \cdot \int_{\mathbb{A}^{LD}(\mathbb{R})} \eta_1$$

- $\alpha > \frac{LD}{2n}$  convergent (regularization)

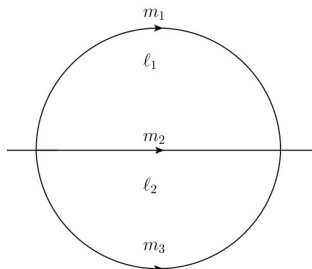
$$\int_{\mathbb{A}^{LD}(\mathbb{R})} \eta_{\alpha,\epsilon} = \int_{\mathbb{P}^{2D}(\mathbb{R})} \eta_{\alpha,\epsilon}$$

period of  $\mathbb{P}^{LD} \setminus Q(\Gamma, m)$

exponent  $\alpha$  changes superficial degree of convergence of Feynman integral from  $\delta(\Gamma) = LD - 2n$  to  $\delta_\alpha(\Gamma) = DL - 2n\alpha$

$\epsilon$ -deformation  $Q_{i,\epsilon}$  ensures differential form  $\eta_{\alpha,\epsilon}$  has no poles on the hyperplane at infinity  $\mathbb{P}^{LN}(\mathbb{R}) \setminus \mathbb{A}^{LN}(\mathbb{R}) = \mathbb{P}^{LN-1}(\mathbb{R})$

## Sunset Graph



in the specific case of the massive sunset graph there is an explicit deformation  $Q_{i,\epsilon}$  of the net of quadrics such that

- $Q_{i,\epsilon}$  smooth, real, positive, satisfying momentum conservation
- the double intersections  $Q_{i,\epsilon} \cap Q_{j,\epsilon}$  and the triple intersection  $Q_{1,\epsilon} \cap Q_{2,\epsilon} \cap Q_{3,\epsilon}$  are all transversal



## main idea

- usual physical perturbation

$$q_{e,\epsilon}(k, x) := k_e^2 + (m_e^2 + i\epsilon)x^2$$

- main purpose to move location of poles in the complex plane
- not good for smoothness and for transversality
- ... but use as model idea

## perturbed quadrics of the Sunset Graph

Zariski open  $W(m) \subset \mathbb{A}^1$  (depending on mass parameter  $m = (m_1, m_2, m_3)$ , with  $m_i \neq 0$ ), for every  $\epsilon \in W(m)$  the deformed quadrics  $Q_{1,\epsilon}, Q_{2,\epsilon}, Q_{3,\epsilon} \subset \mathbb{P}^{2D}$  are smooth, real, positive, satisfying momentum conservation and transverse intersections  $Q_{i,\epsilon} \cap Q_{j,\epsilon}$  and  $Q_{1,\epsilon} \cap Q_{2,\epsilon} \cap Q_{3,\epsilon}$

## Motives

- $M(X)_{\mathbb{Q}}$  mixed motive and  $M^c(X)_{\mathbb{Q}}$  mixed motive with compact support (isomorphic for smooth projective)
- dual motive dual  $M(X)^{\vee} \simeq M^c(X)_{\mathbb{Q}}(-d)[-2d]$
- category of mixed-Tate motives: triangulated subcategory in  $DM_{\text{gm}}(F)_{\mathbb{Q}}$  generated by Tate motives  $\mathbb{L}^k$
- if distinguished triangle in  $DM_{\text{gm}}(F)_{\mathbb{Q}}$  with two out of three terms mixed-Tate  $\Rightarrow$  third one also mixed-Tate
- also if distinguished triangle with one term mixed-Tate and one not  $\Rightarrow$  third one must be non-mixed-Tate

- for any  $X$  smooth  $M(X)_{\mathbb{Q}}$  mixed-Tate iff  $M^c(X)_{\mathbb{Q}}$  mixed-Tate (mixed Tate subcategory stable under duals)
- **Mayer-Vietoris triangle**: Zariski open cover  $X = U \cup V$

$$M^c(X)_{\mathbb{Q}} \longrightarrow M^c(U)_{\mathbb{Q}} \oplus M^c(V)_{\mathbb{Q}} \longrightarrow M^c(U \cap V)_{\mathbb{Q}} \longrightarrow M^c(X)_{\mathbb{Q}}[1]$$

- **Gysin triangle**: Zariski closed subscheme  $Z \subset X$  with open complement  $U$

$$M^c(Z)_{\mathbb{Q}} \longrightarrow M^c(X)_{\mathbb{Q}} \longrightarrow M^c(U)_{\mathbb{Q}} \longrightarrow M^c(Z)_{\mathbb{Q}}[1]$$

**Step 1:** the motives  $M^c(\mathbb{P}^{2D} \setminus Q_{i,\epsilon})_{\mathbb{Q}}$  are mixed-Tate

- $2D$  even, quadric hypersurface  $Q_{i,\epsilon} \subset \mathbb{P}^{2D}$  odd-dimensional
- motivic decomposition of Chow motive (smooth quadrics)

$$h(Q_{i,\epsilon})_{\mathbb{Q}} \simeq 1 \oplus \mathbb{L} \oplus \mathbb{L}^{\otimes 2} \oplus \dots \oplus \mathbb{L}^{\otimes (2D-1)}$$

so  $M^c(Q_{i,\epsilon})_{\mathbb{Q}} \simeq M(Q_{i,\epsilon})_{\mathbb{Q}}$  is mixed-Tate

- Gysin triangle with  $X = \mathbb{P}^{2D}$  and  $Z = Q_{i,\epsilon}$  gives  $M^c(\mathbb{P}^{2D} \setminus Q_{i,\epsilon})_{\mathbb{Q}}$  mixed-Tate

**Step 2:**  $M^c(\mathbb{P}^{2D} \setminus (Q_{i,\epsilon} \cup Q_{j,\epsilon}))_{\mathbb{Q}}$  mixed-Tate iff  
 $M^c(\mathbb{P}^{2D} \setminus (Q_{i,\epsilon} \cap Q_{j,\epsilon}))_{\mathbb{Q}}$  mixed-Tate

- Mayer-Vietoris triangle with  $X := \mathbb{P}^{2D} \setminus (Q_{i,\epsilon} \cap Q_{j,\epsilon})$ ,  
 $U := \mathbb{P}^{2D} \setminus Q_{i,\epsilon}$  and  $V := \mathbb{P}^{2D} \setminus Q_{j,\epsilon}$
- this has  $U \cap V = \mathbb{P}^{2D} \setminus (Q_{i,\epsilon} \cup Q_{j,\epsilon})$
- $M^c(U)_{\mathbb{Q}}$  and  $M^c(V)_{\mathbb{Q}}$  mixed-Tate by Step 1
- if two out of three terms mixed-Tate then third term in the Mayer-Vietoris triangle also mixed-Tate

**Step 3:** assume all  $M^c(\mathbb{P}^{2D} \setminus (Q_{i,\epsilon} \cap Q_{j,\epsilon}))_{\mathbb{Q}}$  are mixed-Tate, then  $M^c(\mathbb{P}^{2D} \setminus Q_{(\Gamma,m)})_{\mathbb{Q}}$  is mixed-Tate iff  $M^c(\mathbb{P}^{2D} \setminus (Q_{1,\epsilon} \cap Q_{2,\epsilon} \cap Q_{3,\epsilon}))_{\mathbb{Q}}$  is mixed-Tate

- take  $U = \mathbb{P}^{2D} \setminus (Q_{1,\epsilon} \cup Q_{2,\epsilon})$  and  $V = \mathbb{P}^{2D} \setminus Q_{3,\epsilon}$  then  $U \cap V = \mathbb{P}^{2D} \setminus Q_{(\Gamma,m)}$
- $M^c(V)_{\mathbb{Q}}$  mixed-Tate by Step 1
- $M^c(\mathbb{P}^{2D} \setminus (Q_{1,\epsilon} \cap Q_{2,\epsilon}))_{\mathbb{Q}}$  mixed-Tate by assumption
- get  $M^c(U)_{\mathbb{Q}}$  mixed Tate by Step 2 and assumption
- Mayer-Vietoris triangle:  $M^c(\mathbb{P}^{2D} \setminus Q_{(\Gamma,m)})_{\mathbb{Q}}$  mixed-Tate iff  $M^c(U \cup V)_{\mathbb{Q}}$  mixed-Tate
- take  $U_{13} := \mathbb{P}^{2D} \setminus (Q_{1,\epsilon} \cap Q_{3,\epsilon})$  and  $U_{23} := \mathbb{P}^{2D} \setminus (Q_{2,\epsilon} \cap Q_{3,\epsilon})$

$$\begin{aligned} U_{13} \cap U_{23} &= \mathbb{P}^{2D} \setminus ((Q_{1,\epsilon} \cap Q_{3,\epsilon}) \cup (Q_{2,\epsilon} \cap Q_{3,\epsilon})) \\ &= \mathbb{P}^{2D} \setminus ((Q_{1,\epsilon} \cup Q_{2,\epsilon}) \cap Q_{3,\epsilon}) = U \cup V \end{aligned}$$

- Mayer-Vietoris triangle:  $M^c(U \cup V)_{\mathbb{Q}}$  mixed-Tate iff  $M^c(U_{13} \cup U_{23})_{\mathbb{Q}}$  mixed-Tate

$$U_{13} \cup U_{23} = \mathbb{P}^{2D} \setminus (Q_{1,\epsilon} \cap Q_{2,\epsilon} \cap Q_{3,\epsilon})$$

**Step 4:** the motives  $M^c(\mathbb{P}^{2D} \setminus (Q_{i,\epsilon} \cap Q_{j,\epsilon}))_{\mathbb{Q}}$  are mixed-Tate

- by transversality the intersections  $Q_{i,\epsilon} \cap Q_{j,\epsilon}$  are smooth complete intersections of two odd-dimensional quadrics
- motivic decomposition of Chow motive  $\mathfrak{h}(Q_{i,\epsilon} \cap Q_{j,\epsilon})_{\mathbb{Q}}$  (Bernardara-Tabuada)

$$1 \oplus \mathbb{L} \oplus \mathbb{L}^{\otimes 2} \oplus \dots \oplus \mathbb{L}^{\otimes (D-2)} \oplus (\mathbb{L}^{\otimes (D-1)})^{\oplus (2D+2)} \oplus \mathbb{L}^{\otimes D} \oplus \dots \oplus \mathbb{L}^{\otimes (2D-2)}$$

so  $M^c(Q_{i,\epsilon} \cap Q_{j,\epsilon})_{\mathbb{Q}} \simeq M(Q_{i,\epsilon} \cap Q_{j,\epsilon})_{\mathbb{Q}}$  mixed-Tate

- Gysin triangle with  $X := \mathbb{P}^{2D}$  and  $Z := Q_{i,\epsilon} \cap Q_{j,\epsilon} \Rightarrow M^c(\mathbb{P}^{2D} \setminus (Q_{i,\epsilon} \cap Q_{j,\epsilon}))_{\mathbb{Q}}$  mixed-Tate

**Step 5:** the motive  $M^c(\mathbb{P}^{2D} \setminus (Q_{1,\epsilon} \cap Q_{2,\epsilon} \cap Q_{3,\epsilon}))_{\mathbb{Q}}$  is **not** mixed-Tate

- same Gysin triangle argument:  $M^c(\mathbb{P}^{2D} \setminus (Q_{1,\epsilon} \cap Q_{2,\epsilon} \cap Q_{3,\epsilon}))_{\mathbb{Q}}$  mixed-Tate iff  $M^c(Q_{1,\epsilon} \cap Q_{2,\epsilon} \cap Q_{3,\epsilon})_{\mathbb{Q}}$  mixed-Tate
- by transversality  $Q_{1,\epsilon} \cap Q_{2,\epsilon} \cap Q_{3,\epsilon}$  smooth complete intersection of three odd-dimensional quadrics
- motivic decomposition of Chow motive  $\mathfrak{h}(Q_{1,\epsilon} \cap Q_{2,\epsilon} \cap Q_{3,\epsilon})_{\mathbb{Q}}$  (Bernardara-Tabuada)

$$1 \oplus \mathbb{L} \oplus \mathbb{L}^{\otimes 2} \oplus \dots \oplus \mathbb{L}^{\otimes (2D-3)} \oplus (\mathfrak{h}^1(J_a^{D-2}(Q_{1,\epsilon} \cap Q_{2,\epsilon} \cap Q_{3,\epsilon}))_{\mathbb{Q}} \otimes \mathbb{L}^{\otimes (D-1)})$$

- $J_a^{D-2}(Q_{1,\epsilon} \cap Q_{2,\epsilon} \cap Q_{3,\epsilon})$  is the  $(D-2)$ -th intermediate algebraic Jacobian of  $Q_{1,\epsilon} \cap Q_{2,\epsilon} \cap Q_{3,\epsilon}$
- abelian variety  $J_a^{D-2}(Q_{1,\epsilon} \cap Q_{2,\epsilon} \cap Q_{3,\epsilon})$  isomorphic to Prym variety  $\text{Prym}(\tilde{C}/C)$
- Prym variety  $\text{Prym}(\tilde{C}/C)$  with  $C$  discriminant divisor of quadric fibration associated to the triple intersection and  $\tilde{C}$  étale double cover of curve  $C$



- if motive  $\mathfrak{h}(Q_{1,\epsilon} \cap Q_{2,\epsilon} \cap Q_{3,\epsilon})_{\mathbb{Q}}$  were mixed-Tate it would be sum of powers of Lefschetz motive  $\mathbb{L}$ , hence only even dimensional cohomology
- but first cohomology

$$\begin{aligned} H^1(\mathfrak{h}(Q_{1,\epsilon} \cap Q_{2,\epsilon} \cap Q_{3,\epsilon})_{\mathbb{Q}}) &= H^1(\mathfrak{h}^1(\mathrm{Prym}(\tilde{C}/C))_{\mathbb{Q}}) \\ &= H^1(\mathrm{Prym}(\tilde{C}/C)) \neq 0 \end{aligned}$$

- so non-mixed-Tate because of term

$$\mathfrak{h}^1(J_a^{D-2}(Q_{1,\epsilon} \cap Q_{2,\epsilon} \cap Q_{3,\epsilon}))_{\mathbb{Q}} \otimes \mathbb{L}^{\otimes(D-1)}$$

**Conclusion:** the Feynman motive  $M^c(\mathbb{P}^{2D} \setminus Q_{(\Gamma,m)})_{\mathbb{Q}}$  of the Sunset Graph is **non-mixed-Tate** (for generic non-zero mass parameters)

## Prym varieties and intermediate Jacobians

- A. Beauville, *Variétés de Prym et Jacobiennes intermédiaires*. Ann. Sci. de l'ENS 10 (1977) 309–391.

## Prym varieties

- $\pi : \tilde{C} \rightarrow C$  étale double covering of curves
- pullback maps on the Jacobians  $\pi^* : J \rightarrow \tilde{J}$
- norm map  $N : \tilde{J} \rightarrow J$  (project points of divisor)
- Prym variety is the kernel of the norm map (largest abelian subvariety of  $\tilde{J}$  on which norm map is trivial)

## Other periods

- same differential form  $\eta_{\alpha,\epsilon}$
- taking the derivative with respect to the mass parameter
- this raises powers of the edge propagators

$$\eta_{\alpha_1, \dots, \alpha_n} = \frac{\omega}{\prod_{i=1}^n q_i^{\alpha_i}}$$

- these are solutions of differential system satisfied by the Feynman integral

**Question:** provide an upper bound estimative for the dimension of the space of these Feynman integrals, through a bound on the dimension of the space of periods on the Feynman motive

**Dimension bound:** upper bound given by  $7 + \dim H^1(\text{Prym}(\tilde{C}/C))$

- estimate dimension of space of periods in middle cohomology
- periods are pairing between de Rham and Betti cohomology:

$$H_{dR}^{2D}(\mathbb{P}^{2D} \setminus Q_{(\Gamma, m)}) \times H_B^{2D}(\mathbb{P}^{2D} \setminus Q_{(\Gamma, m)}) \longrightarrow \mathbb{C}$$

- estimate dimension of middle cohomology  $H^{2D}(\mathbb{P}^{2D} \setminus Q_{(\Gamma, m)})$
- take  $\mathbb{P} = \mathbb{P}^{2D}$ ,  $U = \mathbb{P}^{2D} \setminus (Q_{1,\epsilon} \cup Q_{2,\epsilon})$ ,  $V = \mathbb{P}^{2D} \setminus Q_{3,\epsilon}$ ,  
 $U_{13} = \mathbb{P}^{2D} \setminus (Q_{1,\epsilon} \cap Q_{3,\epsilon})$ ,  $U_{23} = \mathbb{P}^{2D} \setminus (Q_{2,\epsilon} \cap Q_{3,\epsilon})$ , and  
 $Q_{123} = Q_{1,\epsilon} \cap Q_{2,\epsilon} \cap Q_{3,\epsilon}$
- Mayer-Vietoris triangle with  $U \cap V = \mathbb{P} \setminus Q_{(\Gamma, m)}$  gives long exact sequence in cohomology

- in a long exact sequence  $\dots \rightarrow V^{r-1} \rightarrow V^r \rightarrow V^{r+1} \rightarrow \dots$   
dimensions  $\dim(V^r) \leq \dim(V^{r-1}) + \dim(V^{r+1})$

- long exact sequence in cohomology

$$\dots \rightarrow H^r(U \cup V) \rightarrow H^r(U) \oplus H^r(V) \rightarrow H^r(U \cap V) \rightarrow H^{r+1}(U \cup V) \rightarrow \dots$$

- this gives estimate on dimensions

$$\dim H^{2D}(\mathbb{P}^{2D} \setminus Q_{r,m}) \leq \dim H^{2D}(U) + \dim H^{2D}(V) + \dim H^{2D+1}(U \cup V)$$

$$\dim H^{2D+1}(U \cup V) \leq \dim H^{2D+1}(U_{13}) + \dim H^{2D+1}(U_{23}) + \dim H^{2D+2}(\mathbb{P} \setminus Q_{12})$$

- given complete intersection  $Z_1 \cap \dots \cap Z_c =: Z \subset \mathbb{P}$   
codimension  $c$ , Gysin long exact sequence in cohomology:

$$\dots \rightarrow H^{r-2c}(Z) \rightarrow H^r(\mathbb{P}) \rightarrow H^r(\mathbb{P} \setminus Z) \rightarrow H^{r-2c+1}(Z) \rightarrow H^{r+1}(\mathbb{P}) \rightarrow \dots$$

- for  $U_{ij}$  either  $U_{13}$  or  $U_{23}$  with  $Q_{i,\epsilon} \cap Q_{j,\epsilon}$  codimension 2

$$H^{r-4}(Q_{i,\epsilon} \cap Q_{j,\epsilon}) \rightarrow H^r(\mathbb{P}) \rightarrow H^r(U_{ij}) \rightarrow H^{r-3}(Q_{i,\epsilon} \cap Q_{j,\epsilon}) \rightarrow H^{r+1}(\mathbb{P})$$

- for  $r = 2D + 1$  estimate:

$$\dim H^{2D+1}(U_{ij}) \leq \dim H^{2D+1}(\mathbb{P}) + \dim H^{2D-2}(Q_{i,\epsilon} \cap Q_{j,\epsilon})$$

$$\leq \dim H^{2D-2}(Q_{i,\epsilon} \cap Q_{j,\epsilon})$$

- by motivic decomposition have  $\mathfrak{h}^{2D-2}(Q_{i,\epsilon} \cap Q_{j,\epsilon})_{\mathbb{Q}} \simeq \mathbb{L}^{D-1}$
- so get  $\dim H^{2D-2}(Q_{i,\epsilon} \cap Q_{j,\epsilon}) = 1$
- previous estimates give

$$\dim H^{2D+1}(U \cup V) \leq 2 + \dim H^{2D+2}(\mathbb{P} \setminus Q_{123})$$

$$\dim H^{2D}(\mathbb{P}^{2D} \setminus Q_{(r,m)}) \leq 2 + \dim H^{2D}(U) + \dim H^{2D}(V) + \dim H^{2D+2}(\mathbb{P} \setminus Q_{123})$$

- estimate  $\dim H^{2D+2}(\mathbb{P} \setminus Q_{123})$  via Gysin exact sequence

$$H^{r-6}(Q_{123}) \rightarrow H^r(\mathbb{P}) \rightarrow H^r(\mathbb{P} \setminus Q_{123}) \rightarrow H^{r-6+1}(Q_{123})$$

- taking  $r = 2D + 2$  and using  $\dim H^{2D-3}(Q_{123}) = \dim H^1(\text{Prym}(\tilde{C}/C))$

$$\dim H^{2D+2}(\mathbb{P} \setminus Q_{123}) \leq \dim H^1(\text{Prym}(\tilde{C}/C)) + 1$$

- estimate  $\dim H^{2D}(V)$  using Gysin exact sequence for  $Q_{3,\epsilon}$  codimension 1

$$\cdots \rightarrow H^{i-2}(Q_{3,\epsilon}) \rightarrow H^i(\mathbb{P}) \rightarrow H^i(V) \rightarrow H^{i-2+1}(Q_{3,\epsilon}) \rightarrow \cdots$$

- get estimate

$$\dim H^{2D}(V) \leq \dim H^{2D-1}(Q_{3,\epsilon}) + \dim H^{2D}(\mathbb{P})$$

- by motivic decomposition  $Q_{3,\epsilon}$  has no odd cohomology so get  $\dim H^{2D}(V) \leq 1$
- estimate  $\dim H^{2D}(U)$  with Mayer-Vietoris long exact sequence for open cover  $\{\mathbb{P} \setminus Q_{1,\epsilon}, \mathbb{P} \setminus Q_{2,\epsilon}\}$  of  $\mathbb{P} \setminus (Q_{1,\epsilon} \cap Q_{2,\epsilon})$

$$\dim H^{2D}(U) \leq \dim H^{2D+1}(\mathbb{P} \setminus (Q_{1,\epsilon} \cap Q_{2,\epsilon})) + \dim H^{2D}(\mathbb{P} \setminus Q_{1,\epsilon}) + \dim H^{2D}(\mathbb{P} \setminus Q_{2,\epsilon}).$$

- as for  $U_{13}$  and  $U_{23}$  have  $\dim H^{2D+1}(U_{12}) \leq 1$
- similarly to  $V$  have  $\dim H^{2D}(\mathbb{P} \setminus Q_{i,\epsilon}) \leq 1$
- so get  $\dim H^{2D}(U) \leq 3$
- combining all these estimates get middle cohomology bounded by  $7 + \dim H^1(\text{Prym}(\tilde{C}/C))$



**Massless case:** motive of the sunset graph known to be mixed-Tate

- P. Aluffi, M. Marcolli, *Feynman motives of banana graphs*, Commun. Number Theory Phys. 3 (2009), no. 1, 1–57.

**Why previous argument does not apply?**

- quadrics  $Q_i$  in  $\mathbb{P}^{LD-1}$  (no  $x$  coordinate)
- still deform so quadrics smooth and transverse
- even dimensional quadrics  $Q_{i,\epsilon}$  odd dimensional  $\mathbb{P}^{2D-1}$
- motivic decomposition (Bernardara-Tabuada)

$$h(Q_{i,\epsilon} \cap Q_{j,\epsilon})_{\mathbb{Q}} = \begin{cases} \mathbb{L}^{\otimes i/2} & 0 \leq i \leq 2d, i \text{ even} \\ h^1(J_a^{D-2}(Q_{i,\epsilon} \cap Q_{j,\epsilon}))_{\mathbb{Q}} & i = d \\ 0 & \text{otherwise} \end{cases}$$

- now find distinguished triangles with two terms  
non-mixed-Tate: cannot say anything about third

**One-loop triangle graph:** massive case known to be mixed-Tate

- S. Bloch, D. Kreimer, *Mixed Hodge structures and renormalization in physics*. Comm. Number Theory Phys. 2 (2008), no. 4, 637–718.

**Why previous argument does not apply?**

- if zero external momenta, momentum conservation at vertices: same momentum entering through an edge exists through the other
- three Feynman quadrics in same momentum variable  $k \in \mathbb{A}^D$
- momentum conservation condition for the deformations forces identification  $\bar{T}_{i,\epsilon} = \bar{T}_{j,\epsilon}$  on consecutive edges
- no deformation satisfying momentum conservation can achieve transversality