

# Dyson–Schwinger equations in the theory of computation

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ABSTRACT. Following Manin’s approach to renormalization in the theory of computation, we investigate Dyson–Schwinger equations on Hopf algebras, operads and properads of flow charts, as a way of encoding self-similarity structures in the theory of algorithms computing primitive and partial recursive functions and in the halting problem.

## 1. Introduction

There are many deep connections between theoretical physics and information theory, and in particular the theory of computation, see for instance the account given in [1]. In the recent papers [24] and [25], Manin developed a new approach to the theory of computation and the halting problem, based on importing ideas and techniques from the Hopf-algebraic formulation of renormalization in perturbative quantum field theory, [3], [9], [17]. The purpose of this paper is to show that the Hopf algebra of flow charts introduced by Manin in [24] exhibits self-similarity structures given by solutions of combinatorial Dyson–Schwinger equations, defined as in perturbative renormalization, [2], [10], [11], [12], [21], [34]. These can be thought of as a notion of “equations of motion” in the theory of computation.

Dyson–Schwinger equations [7], [29] are a formulation of equations of motion in perturbative quantum field theory, expressed in the form of relations between Green functions. In recent years, an algebraic formulation of the combinatorial structure of perturbative renormalization for scalar field theories was developed, starting with the work of Kreimer on the Hopf-algebraic structure of renormalization [17], followed by the Connes–Kreimer formulation [3] of the BPHZ renormalization procedure, and the formulation in terms of a Riemann–Hilbert correspondence for categories of differential systems in Connes–Marcolli [5], [6]. The Connes–Kreimer formulation of renormalization was given a very general algebraic formulation in terms of Rota–Baxter algebras in the work of Ebrahimi-Fard, Guo and Kreimer [9]. Correspondingly, the Dyson–Schwinger equations were also given a combinatorial form, reflecting the Hopf-algebraic structure of renormalization, in the work of Bergbauer, Kreimer, and Yeats, [2], [21], [34] and more recently with an extensive study by Foissy, [10], [11], [12].

In this paper we investigate the formulation of Dyson–Schwinger equations in the context of Manin’s approach to renormalization and computation. Our main

application will be an alternative formulation of a possible “Galois theory of algorithms”, similar in spirit to the approach proposed by Yanofsky in [33], but more directly related to the Hopf algebra structure of algorithms.

In §2 we review some variants of the construction of a Hopf algebra of *flow charts*. These are diagrams, consisting of decorated planar rooted trees, that compute primitive recursive functions. In §3, we discuss how different version of the Hopf algebra, of flow charts or of “templates”, give rise to slightly different notions of *algorithms*, meant in the sense of [32] as intermediate levels between programs and functions. In §4 we discuss grafting operators and Dyson–Schwinger equations for flow charts, as a way of identifying certain self-similarity structures in the computation of primitive recursive functions. We review some known results about existence and uniqueness of solutions and on conditions under which the coordinates of the solution generate a Hopf subalgebra. We show that one can also consider ideals generated by the coordinates of solutions. Under the condition that the grafting operator is a cocycle, these can give rise to Hopf ideals. The quotient Hopf algebra determines a subgroup of the affine group scheme dual to the (commutative) Hopf algebra of flow charts. We propose this subgroup as a suitable notion of Galois group, following a point of view closely related to Yanofsky’s Galois theory of algorithm [33], although the groups we consider are of a different nature than those considered in [33]. The group we consider consists of those symmetries that are compatible with the self-similarity structure imposed by the Dyson–Schwinger equation. We then show an explicit example of a Dyson–Schwinger equation based on computation by binary trees that does not satisfy the Hopf algebra condition. In order to accommodate this type of example, that is relevant to the setting of computation, we show that it is convenient to take an operadic viewpoint on Dyson–Schwinger equations, which was already suggested in the work of Bergbauer–Kreimer [2]. In §5 we show that Dyson–Schwinger equations can be formulated in the operad setting, using a family of grafting operations, which is not required to satisfy a cocycle condition. We consider Dyson–Schwinger equations on an *operad of flow charts*. We prove the existence and uniqueness of solutions for these operadic Dyson–Schwinger equations and also for a *properad* version. The extension from operads to properads is motivated by the more recent formulation, given by Manin in [27], of a *properad of flow charts* based on directed (acyclic) graphs instead of rooted trees. Finally, §6 is more speculative in nature and it focuses on an algebraic Feynman rule associated to the halting problem proposed in [25]. We also formulate some general questions about the meaning of the BPHZ renormalization procedure and the role of Dyson–Schwinger equations in the halting problem.

## 2. Primitive recursive functions and the Hopf algebra of flow charts

We recall here the setup of Manin [24] on the Hopf algebra of flow charts. First recall from [23] the following facts about primitive recursive functions.

**2.1. Primitive recursive functions.** As in §V.2 of [23], we consider the class of *primitive recursive functions* as generated by the *basic functions*

- Successor  $s : \mathbb{N} \rightarrow \mathbb{N}$ ,  $s(x) = x + 1$ ;
- Constant  $c^n : \mathbb{N}^n \rightarrow \mathbb{N}$ ,  $c^n(x) = 1$  (for  $n \geq 0$ );
- Projection  $\pi_i^n : \mathbb{N}^n \rightarrow \mathbb{N}$ ,  $\pi_i^n(x) = x_i$  (for  $n \geq 1$ );

With the *elementary operations*

- Composition (substitution)  $\mathfrak{c}_{(m,n,p)}$ : for  $f : \mathbb{N}^m \rightarrow \mathbb{N}^n$ ,  $g : \mathbb{N}^n \rightarrow \mathbb{N}^p$ ,

$$g \circ f : \mathbb{N}^m \rightarrow \mathbb{N}^p, \quad \mathcal{D}(g \circ f) = f^{-1}(\mathcal{D}(g));$$

- Bracketing (juxtaposition)  $\mathfrak{b}_{(k,m,n_i)}$ : for  $f_i : \mathbb{N}^m \rightarrow \mathbb{N}^{n_i}$ ,  $i = 1, \dots, k$ ,

$$f = (f_1, \dots, f_k) : \mathbb{N}^m \rightarrow \mathbb{N}^{n_1 + \dots + n_k}, \quad \mathcal{D}(f) = \mathcal{D}(f_1) \cap \dots \cap \mathcal{D}(f_k);$$

- Recursion  $\mathfrak{r}_n$ : for  $f : \mathbb{N}^n \rightarrow \mathbb{N}$  and  $g : \mathbb{N}^{n+2} \rightarrow \mathbb{N}$ ,

$$h(x_1, \dots, x_n, 1) := f(x_1, \dots, x_n),$$

$$h(x_1, \dots, x_n, k+1) := g(x_1, \dots, x_n, k, h(x_1, \dots, x_n, k)), \quad k \geq 1,$$

where recursively  $(x_1, \dots, x_n, 1) \in \mathcal{D}(h)$  iff  $(x_1, \dots, x_n) \in \mathcal{D}(f)$   
and  $(x_1, \dots, x_n, k+1) \in \mathcal{D}(h)$  iff  $(x_1, \dots, x_n, k, h(x_1, \dots, x_n, k)) \in \mathcal{D}(g)$ .

**2.2. Hopf algebras of decorated rooted trees.** We recall here the construction of Hopf algebras of decorated rooted trees and of planar decorated rooted trees and their main properties, see [11] and references therein. Throughout this paper vector spaces and algebras are over a field  $\mathbb{K}$  of characteristic zero.

2.2.1. *Rooted trees.* A rooted tree is a finite graph  $\tau$ , whose geometric realization  $|\tau|$  is simply connected, defined by combinatorial data  $(F_\tau, V_\tau, v_\tau, \delta_\tau, j_\tau)$ , with  $F_\tau$  a set of half-edges (also called flags or tails),  $V_\tau$  the set of vertices with a distinguished element  $v_\tau \in V_\tau$  (the root), a boundary map  $\partial_\tau : F_\tau \rightarrow V_\tau$  that associates to half-edge its boundary vertex and an involution  $j_\tau : F_\tau \rightarrow F_\tau$ ,  $j_\tau^2 = 1$  that performs the matching of half-edges that determines the edges of  $\tau$ . The resulting graph has edges (or internal edges) given by pairs of half-edges matched by the involution, and tails (or external edges) given by half-edges that are fixed by the involution. We denote by  $E_\tau$  the set of internal edges and by  $E_\tau^{ext}$  the set of external edges.

2.2.2. *Orientations.* We consider rooted trees as endowed with the orientation that the root vertex as the *output*, namely where all edges are oriented in the direction of the unique path to the root.

2.2.3. *External edges.* We also assume that the root vertex has an outgoing half-edge and a number of incoming edges, while all other vertices have one outgoing edge and a number of incoming edges and possibly a number of incoming tails. Each leaf of the tree has an incoming tail.

2.2.4. *Planarity.* A *planar* rooted tree is a rooted tree  $\tau$  together with a fixed embedding  $\iota : |\tau| \hookrightarrow \mathbb{R}^2$  of its geometric realization in the plane.

2.2.5. *Decorations.* A (planar) decorated rooted tree is rooted tree and with a map  $\phi_V : V_\tau \rightarrow \mathcal{D}_V$  of the set of vertices to a set  $\mathcal{D}_V$  of vertex-labels and a map  $\phi_F : F_\tau \rightarrow \mathcal{D}_F$  to a set  $\mathcal{D}_F$  of labels of flags. The assignment of a decoration  $(\phi_V, \phi_F)$  to a rooted tree  $\tau$  is in general subject to constraints. For example, if  $f$  and  $f'$  in  $F_\tau$  are matched by the involution  $j_\tau$  to form an edge  $e \in E_\tau$ , then they should carry the same decoration  $\phi_F(f) = \phi_F(f')$ . This defines a labeling of edges. We will see other constraints below, in the specific case of the Hopf algebra of flow charts. We also consider the forgetful map that forgets the flag labels  $\mathcal{D}_F$  and retains the vertex labels  $\mathcal{D}_V$  and a further forgetful map to unlabeled trees.

2.2.6. *Noncommutative Hopf algebra of planar rooted trees.* Given a (planar) rooted tree  $\tau$ , an *admissible cut*  $C$  of  $\tau$  is a modification of the involution  $j_\tau$  that *cuts* a subset of internal edges  $e_i \in E_\tau$  into two flags  $f_i, f'_i$ . Namely, instead of having  $j_\tau(f_i) = f'_i$ , one modifies the involution so that the new  $\hat{j}(f_i) = f_i$  and  $\hat{j}(f'_i) = f'_i$ , where the set of edges  $e_i$  is chosen in such a way that every oriented path in  $\tau$  from a leaf vertex to the root contains at most one  $e_i$ . The new graph obtained using the involution  $\hat{j}$  is a forest

$$C(\tau) = \rho_C(\tau) \amalg \pi_C(\tau)$$

consisting of one (planar) rooted tree  $\rho_C(\tau)$  containing the root vertex of  $\tau$  and a finite disjoint union of other oriented planar trees  $\pi_C(\tau) = \amalg_i \pi_{C,i}(\tau)$ , where each tree  $\pi_{C,i}(\tau)$  has a single output external edge, to which we assign as root the boundary vertex of this output.

The *noncommutative* Hopf algebra of planar rooted trees  $\mathcal{H}^{nc}$  is defined, as an algebra, as the free algebra generated by the planar rooted trees, with the coalgebra structure given by the *admissible-cuts* coproduct

$$(2.1) \quad \Delta(\tau) = \tau \otimes 1 + 1 \otimes \tau + \sum_C \pi_C(\tau) \otimes \rho_C(\tau).$$

The coproduct is coassociative and the Hopf algebra is graded by the number of vertices of rooted trees, hence the antipode is defined inductively by the formula

$$S(x) = -x - \sum S(x')x'', \quad \text{for } \Delta(x) = x \otimes 1 + 1 \otimes x + \sum x' \otimes x''$$

with  $x', x''$  terms of lower degree. See [13], [15] for more details on this Hopf algebra. An element  $x$  in a Hopf algebra is primitive if  $\Delta(x) = 1 \otimes x + x \otimes 1$ .

2.2.7. *Commutative Hopf algebra of rooted trees.* In the Hopf algebraic approach to perturbative renormalization, as formulated by Kreimer [17] and Connes–Kreimer [3], one considers a *commutative* Hopf algebra of rooted trees decorated by Feynman graphs, with the admissible-cuts coproduct (2.1), or else a commutative Hopf algebra of Feynman graphs, with a coproduct that corresponds to inclusions of subgraphs (subdivergences) and quotient graphs.

In the case of planar rooted trees, one can also consider a commutative Hopf algebra  $\mathcal{H}$ , which is a quotient of  $\mathcal{H}^{nc}$  by imposing the commutativity of multiplication. The generators of  $\mathcal{H}$  are still planar rooted trees, but monomials in these generators no longer identify with *embedded* forests. The grading and antipode are as above, induced by those of  $\mathcal{H}^{nc}$ .

**2.3. Manin’s Hopf algebra of flow charts.** We consider, as in [24], the set of planar labelled rooted trees, where the label set  $\mathcal{D}_V$  of vertices is given by the set of elementary operations

$$\{\mathbf{c}_{(m,n,p)}, \mathbf{b}_{(k,m,n_i)}, \mathbf{r}_n\}$$

(composition, bracketing, recursion) and the label set  $\mathcal{D}_F$  of the flags is the set of primitive recursive functions. Notice that, because the inputs of the composition, bracketing and recursion operations are ordered, we need to work with *planar* rooted trees.

The data  $(\phi_V(v), \phi_F(f_i))$  of a labelings  $\phi_V(v) \in \mathcal{D}_V$  and  $\phi_F(f_i)$  of a vertex  $v \in V_\tau$  and of all the half-edges  $f_i \in F_\tau$  with  $\partial(f_i) = v$  are *admissible* if they satisfy the following conditions:

- If  $\phi_V(v) = \mathbf{c}_{(m,n,p)}$ , then  $v$  must have valence three; the labels  $h_1 = \phi_F(f_1)$  and  $h_2 = \phi_F(f_2)$  of the two incoming flags must have domains and ranges  $h_1 : \mathbb{N}^m \rightarrow \mathbb{N}^n$  and  $h_2 : \mathbb{N}^n \rightarrow \mathbb{N}^p$  and the outgoing flag must be labeled by the composition  $h_2 \circ h_1 = \mathbf{c}_{(m,n,p)}(h_1, h_2)$ .
- If  $\phi_V(v) = \mathbf{r}_n$ , then  $v$  must have valence three; the labels  $h_1 = \phi_F(f_1)$  and  $h_2 = \phi_F(f_2)$  of the two incoming flags must have domains and ranges  $h_1 : \mathbb{N}^n \rightarrow \mathbb{N}$  and  $h_2 : \mathbb{N}^{n+2} \rightarrow \mathbb{N}$  and the outgoing flag must be labeled by the recursion  $h = \mathbf{r}_n(h_1, h_2)$ .
- If  $\phi_V(v) = \mathbf{b}_{(k,m,n_i)}$  then  $v$  must have valence  $k + 1$  and all the functions  $h_i = \phi_F(f_i)$  associated to the incoming flags must have domain in the same  $\mathbb{N}^m$ . The outgoing flag must be labeled by the bracketing  $f = (f_1, \dots, f_k) = \mathbf{b}_{(k,m,n_i)}(f_1, \dots, f_k)$ .

DEFINITION 2.1. *An admissible labeling of a planar rooted tree  $\tau$  by primitive recursive functions is a collection of labelings*

$$\{(\phi_V(v), \phi_F(f_i)) \mid v \in V_\tau \text{ and } f_i \in F_\tau \text{ with } \partial(f_i) = v\}$$

*such that the admissibility condition above is satisfied at all vertices and such that  $\phi_F(f) = \phi_F(f')$  whenever  $j_\tau(f) = f'$ .*

*A vertex labeling of a planar rooted tree  $\tau$  by elementary operations is a collection of labeling  $\{\phi_V(v) \in \{\mathbf{c}, \mathbf{r}, \mathbf{b}_k\} \mid v \in V_\tau\}$  with the only constraint that labels of type  $\mathbf{c}$  and  $\mathbf{r}$  can only be assigned to vertices with two incoming flags and one outgoing flag and labels of type  $\mathbf{b}_k$  can be assigned to vertices with  $k$  incoming flags and one outgoing flag.*

Via the forgetful map that keeps the vertex labels and forgets the flag labels, a planar rooted tree with an admissible labeling by primitive recursive functions determines one with vertex labeling by elementary operations, while not all vertex labeling by elementary operations on a given planar rooted tree will admit a compatible flag labeling. Notice that, in the case of vertex labelings, we no longer distinguish between labels  $\mathbf{c}_{(m,n,p)}$ , with different choices of  $m$ ,  $n$  and  $p$ , since we do not fix the flag labels: all these vertex labels in  $\mathcal{D}_V$  correspond to the same label  $\mathbf{c}$ . Similarly, all the labels  $\mathbf{r}_n$  of admissible labeling correspond to the same label  $\mathbf{r}$  of vertex labeling, and all the  $\mathbf{b}_{(k,m,n_i)}$  correspond to the same  $\mathbf{b}_k$ . We retain the different  $\mathbf{b}_k$  as vertex labelings as those correspond to different valences of the vertex.

Thus, a planar rooted tree with an admissible labeling by primitive recursive functions can be interpreted as a *flow chart* that computes the output function starting from the functions associated to the incoming external edges according to the operations performed at the vertices.

A planar rooted tree  $\tau$  with a vertex labeling by elementary operations, on the other hand, should be thought of as a template of a possible computational scheme: when applied to a tuple of primitive recursive functions  $(h_1, \dots, h_k)$ , with  $k$  the number of incoming flags of  $\tau$ , it either computes an output functions  $h$ , if the inputs  $(h_1, \dots, h_k)$  together with the assigned vertex labels determine an admissible labeling of  $\tau$  by partial recursive function, or else we set its output to be the empty function.

We then define two slightly different versions of the Hopf algebra of flow charts, depending on the type of labeling that we want to use on trees.

DEFINITION 2.2. *The noncommutative Hopf algebra  $\mathcal{H}_{\text{flow},\mathcal{P}}^{\text{nc}}$  of flow charts is the free algebra generated by the planar rooted trees with admissible labelings by primitive recursive functions, with the coproduct (2.1). We also denote by  $\mathcal{H}_{\text{flow},\mathcal{P}}$  the commutative quotient of  $\mathcal{H}_{\text{flow},\mathcal{P}}^{\text{nc}}$ . Similarly, the noncommutative Hopf algebra  $\mathcal{H}_{\text{flow},\mathcal{V}}^{\text{nc}}$  is the free algebra generated by planar rooted trees with vertex labeling by elementary operations, again with the admissible-cuts coproduct. We denote by  $\mathcal{H}_{\text{flow},\mathcal{V}}$  its commutative quotient.*

Notice that if  $\tau$  has an admissible labeling, then the rooted trees  $\rho_C(\tau)$  and  $\pi_{C,i}(\tau)$  also inherit induced admissible labelings, so the Hopf algebras are well defined. There is a Hopf algebra homomorphism  $\mathcal{H}_{\text{flow},\mathcal{P}}^{\text{nc}} \rightarrow \mathcal{H}_{\text{flow},\mathcal{V}}$ , determined by forgetting the flag labels, but as observed above not all generators of  $\mathcal{H}_{\text{flow},\mathcal{V}}$  are in the image of this map. Since the rooted trees in  $\mathcal{H}_{\text{flow},\mathcal{V}}$  have no labels attached to flags, we can represent them as rooted trees without external edges. This will be implicitly done in the description of the grafting operators in Lemma 4.1 below.

In the following, whenever a statement applies to both types of tree decorations on the Hopf algebras of flow charts, we will drop the  $\mathcal{P}$  and  $\mathcal{V}$  subscripts.

2.3.1. *Binary versus  $k$ -ary operations.* It is possible to reformulate the Hopf algebra of flow charts by using only *binary* trees. In fact, it is shown in §2 of [32] that the description of primitive recursive functions in terms of basic functions (successor, constants and projections) and elementary operations (composition, bracketing and recursion), as recalled above, can be reformulated as repeated applications of *binary* operations of these same forms. In fact, the composition and recursion operations already are binary (they label vertices with two input flags and one output) so one only needs to verify that the bracketing operation  $\mathfrak{b}_{(k,m,n_i)}$  assigning to  $k$  functions  $f_i : \mathbb{N}^m \rightarrow \mathbb{N}^{n_i}$ ,  $i = 1, \dots, k$  their juxtaposition  $f = (f_1, \dots, f_k) : \mathbb{N}^m \rightarrow \mathbb{N}^{n_1 + \dots + n_k}$  is a composition

$$\mathfrak{b}_{(k,m,n_i)} = \mathfrak{b}_{(2,m,n_1,n_2+\dots+n_k)} \circ \dots \circ \mathfrak{b}_{(2,m,n_{k-1},n_k)}.$$

This replaces a single vertex with  $k$  inputs and one output with  $k$  vertices with two inputs and one output. Thus, without loss of generality in the algorithmic representation of primitive recursive functions, one can consider a sub-Hopf algebra  $\mathcal{H}_{\text{flow}}^{\text{nc},0} \subset \mathcal{H}_{\text{flow}}^{\text{nc}}$  generated by *binary* planar rooted tree with an admissible labeling by primitive recursive functions, and the corresponding commutative quotient  $\mathcal{H}_{\text{flow}}^0 \subset \mathcal{H}_{\text{flow}}$ .

On the other hand, one can also extend the two binary operations of composition and recursion to  $k$ -ary operations for arbitrary  $k \geq 2$ . In the case of composition, we can introduce a new set of vertex labels  $\mathfrak{c}_{(k,n_i)}$  that correspond to the  $k$ -ary compositions  $\mathfrak{c}_{(k,m,n_i)}(h_i) = h_k \circ \dots \circ h_1$  of functions  $h_i : \mathbb{N}^{n_{i-1}} \rightarrow \mathbb{N}^{n_i}$ , for  $i = 1, \dots, k$ , with  $n_0 = m$ . In the case of recursion, for  $k \geq 1$ , one can consider  $(k+1)$ -ary recursions where the recursion depends on  $k$  initial conditions. Namely, we define  $\mathfrak{r}_{k+1,n}$  as the operation that computes the function  $h = \mathfrak{r}_{k,n}(h_1, \dots, h_{k+1}) : \mathbb{N}^{n+1} \rightarrow \mathbb{N}$ , from the input  $h_i : \mathbb{N}^n \rightarrow \mathbb{N}$  for  $i = 1, \dots, k$

and  $h_{k+1} : \mathbb{N}^{2n+1} \rightarrow \mathbb{N}$ , as

$$\begin{aligned}
 (2.2) \quad & h(x_1, \dots, x_n, 1) = h_1(x_1, \dots, x_n), \\
 & \vdots \\
 & h(x_1, \dots, x_n, k) = h_k(x_1, \dots, x_n), \\
 & h(x_1, \dots, x_n, k + \ell) = h_{k+1}(x_1, \dots, x_n, h_1(x_1, \dots, x_n), \dots, h_k(x_1, \dots, x_n), k + \ell - 1), \\
 & \text{for } \ell \geq 1
 \end{aligned}$$

These  $k$ -ary operations are clearly reducible to compositions of binary operations, hence any primitive recursive function that can be computed using these operations can also be computed by the presentation in terms of binary compositions and recursions. However, as we shall discuss more in detail below, it is convenient to include these explicitly in the structure of the Hopf algebra of flow charts in order to have well defined grafting operators that give rise to Hochschild cocycles and to Dyson–Schwinger equations.

Thus, we propose another minor modification of Manin’s Hopf algebra of flow charts, as follows.

**DEFINITION 2.3.** *We denote by  $\mathcal{H}_{\text{flow}, \mathcal{P}'}^{nc}$  and  $\mathcal{H}_{\text{flow}, \mathcal{Y}'}^{nc}$  the noncommutative Hopf algebras of flow charts where the vertex label sets is enlarged to contain all the  $k$ -ary compositions and recursions  $\mathfrak{c}_{(k, n_i)}$  and  $\mathfrak{r}_{k, n}$ , and the admissibility condition is restated accordingly. We denote by  $\mathcal{H}_{\text{flow}, \mathcal{P}'}$  and  $\mathcal{H}_{\text{flow}, \mathcal{Y}'}$  their commutative quotients.*

### 3. Flow charts, templates, and algorithms

In the above, we have introduced different versions of rooted trees with vertex and edge decorations. We discuss here briefly the meaning of the objects considered above from the point of view of computation.

**3.1. Flow charts versus templates.** A rooted tree with vertices decorated by elementary operations and edges compatibly decorated by recursive functions is a *flow chart*. It describes a concrete realization of a primitive recursive function (the output) in terms of a computation that starts with other primitive recursive functions, possibly basic (the inputs of the tree) with the operations performed at the vertices.

By extension, we also refer to arbitrary elements of the Hopf algebra  $\mathcal{H}_{\text{flow}, \mathcal{P}'}^{nc}$  as “flow charts”. These are formal sums of forests, with a planar embedding, decorated as above by primitive recursive functions at edges and elementary operations at vertices. In this extended sense a flow charts computes a formal sum of ordered sequences of primitive recursive functions, given by the collection of the outputs of each of the trees.

We also considered a different type of objects, given by rooted trees where only the vertices are decorated by elementary operations, while the edges are undecorated. These are computational architectures where the inputs are open: once a set of primitive recursive functions is assigned as input of such a tree, one obtains a primitive recursive function as output, computed according to the operations specified at the vertices. We refer to these trees with vertex decorations as *templates*.

By extension we also refer to an arbitrary element of the Hopf algebra  $\mathcal{H}_{\text{flow}, \nu}^{nc}$  as a *template*. These are formal sums of forests with vertex decorations by elementary operations.

The difference between working with flow charts (vertex and edge decorations) or with what we call here templates (vertex decorations only) will become more transparent when one considers implementing certain “reasonable equivalence relations”, which describe, in the general philosophy of Yanofsky [32], an intermediate level of *algorithms*, in between programs (flow charts) and primitive recursive functions.

**3.2. The notion of algorithm.** Yanofsky’s work [32] provides an in depth discussion of the notion of *algorithm*. To summarize quickly his general approach, one considers equivalence classes consisting of different programs that implement “the same algorithm”. In this view, algorithms describe any intermediate level between the programs (the flow charts in our language above) and the primitive recursive functions they compute. Rather than fixing precisely what the intermediate level should be, Yanofsky allows for some freedom, by considering all the possible quotients of the set of programs by “reasonable equivalence relations” and organizing all these intermediate levels into a hierarchy governed by a suitable kind of Galois correspondence, [32], [33].

In Section 3 of [32], Yanofsky gives a list of examples of “reasonable equivalence relations” one can impose on programs. These include: associativity of composition, projections acting as identity, distributivity of composition over bracketing, associativity of bracketing, almost-commutativity of bracketing, functoriality of bracketing, idempotency of twists, Reidemeister moves, relation between recursion and bracketing, relation between recursion and composition, relation between recursion and the successor function.

Our viewpoint here follows closely this idea of Yanofsky, but with one important difference. We consider the existence of a Hopf algebra structure on flow charts as an essential part of the data and we therefore require that a notion of “reasonable equivalence relation” should include the compatibility with the Hopf algebra structure. This means requiring that equivalence relations among flow charts should define Hopf ideals.

Notice how some of the relations listed above from [32] reflect properties (associativity, distributivity, almost-commutativity) of the elementary operations, which are independent of the input functions, while others involve relations between the elementary operations and certain specific input functions (projections, successor). The different nature of these two types of relations will be reflected, in our viewpoint, in the fact that the first type of relations can be implemented both at the level of the Hopf algebra of flow charts  $\mathcal{H}_{\text{flow}, \mathcal{P}}^{nc}$  and at the level of the Hopf algebra of templates  $\mathcal{H}_{\text{flow}, \nu}^{nc}$ , while the second type of relation can only be implemented in  $\mathcal{H}_{\text{flow}, \mathcal{P}}^{nc}$ .

To see a concrete example of both cases, let us first consider the relation that corresponds to associativity of composition, as stated in §3.1.1 of [32]. This implies

identifying trees of the form

$$(3.1) \quad \tau_1 = \begin{array}{c} \diagup \quad \diagdown \\ \bullet \mathbf{c} \\ \text{\scriptsize } f \quad \text{\scriptsize } g \\ | \\ \text{\scriptsize } f \circ g \\ | \\ \diagup \quad \diagdown \\ \bullet \mathbf{c} \\ \text{\scriptsize } f \quad \text{\scriptsize } h \\ | \\ \text{\scriptsize } f \circ g \circ h \end{array} \cong \begin{array}{c} \diagup \quad \diagdown \\ \bullet \mathbf{c} \\ \text{\scriptsize } g \quad \text{\scriptsize } h \\ | \\ \text{\scriptsize } g \circ h \\ | \\ \diagup \quad \diagdown \\ \bullet \mathbf{c} \\ \text{\scriptsize } f \quad \text{\scriptsize } h \\ | \\ \text{\scriptsize } f \circ g \circ h \end{array} = \tau_2$$

In order to implement this relation at the level of the Hopf algebra, we need to construct a Hopf ideal. The ideal  $\mathcal{I} = (\tau_1 - \tau_2)$  in the underlying algebra would not, by itself, be a Hopf ideal, because the non-primitive part of the coproduct  $\Delta(\tau_1 - \tau_2)$  would not be in  $\mathcal{H} \otimes \mathcal{I} \oplus \mathcal{I} \otimes \mathcal{H}$ , but if we also consider the relation identifying the trees

$$(3.2) \quad \tau'_1 = \begin{array}{c} | \\ \bullet \mathbf{c} \\ \text{\scriptsize } f \\ \diagdown \quad \diagup \\ \bullet \mathbf{c} \\ \text{\scriptsize } f \quad \text{\scriptsize } g \\ | \\ \text{\scriptsize } f \circ g \end{array} \cong \begin{array}{c} | \\ \bullet \mathbf{c} \\ \text{\scriptsize } g \\ \diagdown \quad \diagup \\ \bullet \mathbf{c} \\ \text{\scriptsize } f \quad \text{\scriptsize } g \\ | \\ \text{\scriptsize } f \circ g \end{array} = \tau'_2$$

and the ideal  $\mathcal{I} = (\tau_1 - \tau_2, \tau'_1 - \tau'_2)$ , then the coproduct satisfies  $\Delta(\mathcal{I}) \subset \mathcal{H} \otimes \mathcal{I} \oplus \mathcal{I} \otimes \mathcal{H}$ . The resulting Hopf ideal is the appropriate way, in our setting, to implement the “associativity of composition” relation.

Notice that the relations defining  $\mathcal{I}$  do not depend on particular input functions being assigned to the incoming edges, hence they not only define a Hopf ideal in the Hopf algebra  $\mathcal{H}_{\text{flow}, \mathcal{P}}^{nc}$  of flow charts, but they also define a Hopf ideal in the Hopf algebra  $\mathcal{H}_{\text{flow}, \mathcal{V}}^{nc}$  of templates.

We then consider, as a second example, the relation defined in §3.1.2 of [32], which represents the fact that projection  $\pi^n = (\pi_i^n) : \mathbb{N}^n \rightarrow \mathbb{N}^n$  is the identity. This relation is implemented by identifying, for any primitive recursive function  $f : \mathbb{N}^n \rightarrow \mathbb{N}^m$ , the decorated rooted trees

$$(3.3) \quad \begin{array}{c} \diagup \quad \diagdown \\ \bullet \mathbf{c} \\ \text{\scriptsize } f \\ | \\ \text{\scriptsize } f \end{array} \cong \begin{array}{c} \diagup \quad \diagdown \\ \bullet \mathbf{c} \\ \text{\scriptsize } f \quad \text{\scriptsize } f \\ | \\ \text{\scriptsize } f \end{array} \cong \begin{array}{c} | \\ \bullet \mathbf{c} \\ \text{\scriptsize } f \end{array}$$

As in the previous case, one can associate to this relation a Hopf ideal in  $\mathcal{H}_{\text{flow}, \mathcal{P}}^{nc}$ . However, in this case, the relation involves not only a property of the elementary operations that label the vertices but also of a specific choice of input functions, hence it does not define a Hopf ideal in  $\mathcal{H}_{\text{flow}, \mathcal{V}}^{nc}$ .

Thus, from the point of view of computation and algorithms, the difference between working with flow charts or with templates corresponds to two different notions of algorithms. In the first case algorithms are equivalence classes of flow charts under relations that can involve both properties of the elementary operations that hold for arbitrary input functions and properties that hold for specific input functions, while in the case of templates, one only allows the first type of relations in the definition of algorithms.

**3.3. Symmetries: a physics motivated perspective.** We will be discussing, especially in relation to Dyson–Schwinger equations, notions of *symmetry* in the context of Hopf algebras. The main motivating example for the notion of symmetry we adopt is again coming from the setting of quantum field theory, where gauge symmetries are implemented at the quantum level in the form of Ward (or Slavonov–Taylor) identities. It was shown by van Suijlekom in [30] that the Slavonov–Taylor identities define a Hopf ideal in the Connes–Kreimer Hopf algebra of Feynman graphs. It is therefore a natural approach from the physics perspective to require that a good notion of symmetry in a setting where the main objects of interest (Feynman graphs, flow charts, etc.) are described in terms of a Hopf algebra should be expressed in terms of Hopf ideals. We adopt this viewpoint in the following, in §4.5, where we propose a different approach to Yanofsky’s Galois theory of algorithms, using Hopf ideals in the Hopf algebra of flow charts as a good notion of equivalence relations defining intermediate levels (algorithms) between flow charts and primitive recursive functions, and in §4.4.1, where we consider Hopf ideals associated to solutions of Dyson–Schwinger equations.

#### 4. Dyson–Schwinger equations in the Hopf algebra of flow charts

Recall here the formulation of combinatorial Dyson–Schwinger equations in Hopf algebras of decorated rooted trees, following [2], [34] and especially [10], [11], [12]. We then focus on the specific case of the Hopf algebras of flow charts described above.

**4.1. Insertion operators for decorated planar rooted trees.** We work here with the Hopf algebra of flow charts  $\mathcal{H}_{\text{flow}, \mathcal{V}'}^{nc}$ . By analogy with the setting of perturbative renormalization [2], [34], we define, for each type of elementary operation  $\delta \in \{\mathfrak{b}, \mathfrak{c}, \mathfrak{r}\}$ , a grafting operator  $B_\delta^+$  defined in the following way. Given a monomial  $T$  in  $\mathcal{H}_{\text{flow}, \mathcal{V}'}^{nc}$ , which consists of a forest with vertex labeling by elementary operations, we define  $B_\delta^+(T)$  as the sum of planar graphs obtained by adding a new root vertex with a number of incoming flags equal to the number of trees in  $T$  and a single output flag, with the vertex decorated by  $\delta$ .

4.1.1. *Hochschild 1-cocycles.* In the case of perturbative renormalization, it is well known [2] that the analogous grafting operators  $B_\delta^+$ , with vertex decorations  $\delta$ , define Hochschild 1-cocycles.

The 1-cocycle condition for the operator  $B_\delta^+$  then consists of the property that

$$(4.1) \quad \Delta B_\delta^+ = (id \otimes B_\delta^+) \Delta + B_\delta^+ \otimes 1.$$

LEMMA 4.1. *The operators  $B_\delta^+$  on  $\mathcal{H}_{\text{flow}, \mathcal{V}'}^{nc}$ , with  $\delta \in \{\mathfrak{b}, \mathfrak{c}, \mathfrak{r}\}$ , satisfy the 1-cocycle condition (4.1).*

PROOF. This case works exactly as in the renormalization setting, see [2]: we reproduce the argument here for the reader’s convenience. For an element  $x$  in a Hopf algebra  $\mathcal{H}$  with coproduct

$$\Delta(x) = 1 \otimes x + x \otimes 1 + \sum x' \otimes x'',$$

where  $x'$  and  $x''$  denote lower degree terms, we use the notation

$$\tilde{\Delta}(x) := \sum x' \otimes x''.$$

One then sees easily that the cocycle condition (4.1) is equivalent to

$$(4.2) \quad \tilde{\Delta} B_\delta^+ = (id \otimes B_\delta^+) \tilde{\Delta} + id \otimes B_\delta^+(1),$$

where, as in the renormalization case, with  $B_\delta^+(1) = v_\delta$  is a single vertex with assigned label  $\delta$ . The argument is then as in the original case, with the first term in the right hand side of (4.2) accounting for the admissible cuts where the root vertex remains attached to the  $\rho_C(T)$  part of an admissible cut of  $T$ , and the second term counts the case where the admissible cut separates the root vertex completely.  $\square$

Notice that, for the cocycle condition to hold, we need to be able to assign the same type of label (**b**, **c**, or **r**) to vertices of arbitrary valence, hence the reason for including in the set of vertex decorations all the  $k$ -ary versions of composition and recursion and working with the version  $\mathcal{H}_{\text{flow}, \mathcal{V}'}^{nc}$  of the Hopf algebra of flow charts.

It is possible to define other interesting grafting operators in the Hopf algebras of flow charts, which, however, do not satisfy the cocycle condition. These can still be used to construct Dyson–Schwinger equations, but the set of solutions does not define a Hopf subalgebra. This will be discussed in an example in the following subsection.

**4.2. Systems of combinatorial Dyson–Schwinger equations.** Given a graded Hopf algebra  $\mathcal{H}$ , we consider the direct product  $\overline{\mathcal{H}} = \prod_{n=0}^\infty \mathcal{H}_n$ , whose elements we write as infinite sums  $x = \sum_{n=0}^\infty x_n$ , with  $x_n \in \mathcal{H}_n$ , endowed with the associative product and coproduct induced by those of  $\mathcal{H}$ . Following [10], (see also [2], [34]) given a formal power series

$$P(t) = \sum_{k=0}^\infty a_k t^k, \quad \text{with } a_0 = 1,$$

and a Hochschild 1-cocycle  $B^+$  on the Hopf algebra  $\mathcal{H}$ , the associated Dyson–Schwinger equation is given by

$$(4.3) \quad X = B^+(P(X)).$$

One considers this as an equation in  $\overline{\mathcal{H}}$  and interprets the infinite sum  $P(X) = \sum_k a_k X^k$  as an element in  $\overline{\mathcal{H}}$ .

One should read (4.3) as a fixed point equation for the nonlinear transformation

$$X \mapsto B^+(P(X))$$

of  $\overline{\mathcal{H}}$ . Solutions to the fixed point equation are elements of  $\overline{\mathcal{H}}$  that exhibit a self-similarity property with respect to the operation of first applying  $X \mapsto P(X)$  (in the case of a polynomial, this operation heuristically replaces  $X$  by a sum of multiple copies of itself weighted by the coefficients of the polynomial) and then grafting them together according to the cocycle  $B^+$ .

In fact, the equation (4.3) has a unique solution given by an element  $X = \sum_k x_k$  with  $x_k \in \mathcal{H}_k$  whose homogeneous components are determined inductively by the procedure

$$(4.4) \quad x_{n+1} = \sum_{k=1}^n \sum_{j_1+\dots+j_k=n} a_k B^+(x_{j_1} \cdots x_{j_k}),$$

starting with  $x_1$  given by  $B^+(1)$ , see Lemma 2 of [2] and Proposition 2 of [10].

Notice that we can still define the Dyson–Schwinger equation (4.3) for other types of grafting operators that do not necessarily satisfy the 1-cocycle condition (4.1), such as the examples discussed in the previous subsection. The existence and uniqueness of solutions, of the form (4.4) is also still valid. Where the 1-cocycle condition is crucially used is in showing that the associative subalgebra of  $\mathcal{H}$  generated by the components  $x_k$  of the solution is also a Hopf subalgebra, see Theorem 3 of [2].

4.2.1. *Systems of Dyson–Schwinger equations.* As discussed in the work of Foissy [11], in cases like our  $\mathcal{H}_{\text{flow}, \nu'}^{nc}$ , where one has different vertex labels  $\delta \in \{\mathfrak{b}, \mathfrak{c}, \mathfrak{r}\}$ , one can consider more complicated *systems* of Dyson–Schwinger equations, involving the grafting operators  $B_\delta^+$  on  $\mathcal{H}_{\text{flow}, \nu'}^{nc}$ .

For  $\delta \in \{\mathfrak{b}, \mathfrak{c}, \mathfrak{r}\}$ , we consider the grafting cocycles  $B_\delta^+$  on the Hopf algebra  $\mathcal{H}_{\text{flow}, \nu'}^{nc}$ , as above. We also consider data of three non-constant formal power series in three variables  $X = (X_\delta)$

$$F_\delta(X) = \sum_{k_1, k_2, k_3} a_{k_1, k_2, k_3}^{(\delta)} X_{\mathfrak{b}}^{k_1} X_{\mathfrak{c}}^{k_2} X_{\mathfrak{r}}^{k_3}.$$

To these data we associate the system of Dyson–Schwinger equations as in [11]

$$(4.5) \quad X_\delta = B_\delta^+(F_\delta(X)).$$

As shown in Proposition 5 of [11], these systems of equations also have a unique solution of the form  $X_\delta = \sum_\tau x_\tau \tau$ , with the sum over all planar rooted trees with root decorated by  $\delta$ , with coefficients

$$(4.6) \quad x_\tau = \left( \prod_{k=1}^3 \frac{(\sum_{l=1}^{m_k} p_{\delta, l})!}{\prod_{l=1}^{m_k} p_{\delta, l}!} \right) a_{\sum_{k=1}^3 p_{1, k}, \sum_{k=1}^3 p_{2, k}, \sum_{k=1}^3 p_{3, k}}^{(\delta)} x_{\tau_{1,1}}^{p_{1,1}} \cdots x_{\tau_{3, m_3}}^{p_{3, m_3}},$$

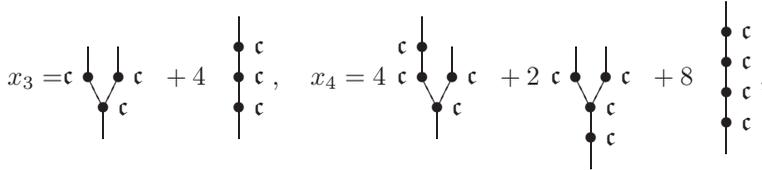
when

$$\tau = B^+(\tau_{1,1}^{p_{1,1}} \cdots \tau_{1, m_1}^{p_{1, m_1}} \cdots \tau_{3,1}^{p_{3,1}} \cdots \tau_{3, m_3}^{p_{3, m_3}}).$$

**4.3. Dyson–Schwinger solutions: a computational perspective.** The solution  $X_\delta = \sum_\tau x_\tau \tau$  of a Dyson–Schwinger equation  $X_\delta = B_\delta^+(F_\delta(X))$  as above, is a formal infinite linear combination of templates  $\tau$ , with coefficients  $x_\tau$  given by recursively defined computable functions of the coefficients  $a_{k_1, k_2, k_3}^{(\delta)}$  of  $F_\delta$ . It takes an input  $f^{\text{input}} = (f_\tau^{\text{input}})_\tau$  consisting of a sequence of primitive recursive functions, where  $f_\tau^{\text{input}} = (f_{\tau,1}^{\text{input}}, \dots, f_{\tau, n_\tau}^{\text{input}})$  are the  $n_\tau$  inputs of a template  $\tau$  given by a tree with  $n_\tau$  leaves. It produces an output given by a formal linear combination  $\sum_\tau x_\tau f_\tau^{\text{output}}$ , where  $f_\tau^{\text{output}}$  is the output of the template  $\tau$  for the given input  $f_\tau^{\text{input}}$ . More precisely, to avoid dealing with infinite linear combinations, one can consider the individual components of the solution obtained by recursively solving the Dyson–Schwinger equation.

For simplicity, we look at how this works in a toy-model example of Dyson–Schwinger equation. Let us restrict to only one possible type of vertex labeling by composition  $\mathfrak{c}$ , and consider an equation of the form  $X = 1 + B_{\mathfrak{c}}^+(X^2)$ . The components  $x_n$  of the solution are then obtained recursively as

$$x_0 = 1, \quad x_1 = \bullet \mathfrak{c}, \quad x_2 = \begin{array}{c} \bullet \mathfrak{c} \\ | \\ \bullet \mathfrak{c} \end{array},$$



and so on. More realistically, one would have to count also templates with other possible labelings of vertices besides  $\mathfrak{c}$ , but for simplicity let us focus on this case. Then, for example, we interpret  $x_3$  as a template that, for a given input  $(f_1, f_2)$ , consisting of a pair of primitive recursive functions computes the output function  $6f_1 \circ f_2 + 8f_1$ . Notice that the self-similar structure of the Dyson–Schwinger equation is not apparent when one looks at one individual component of the solution but only in the overall recursive relation  $x_{n+1} = \sum_{k=0}^n B_{\mathfrak{c}}^+(x_k x_{n-k})$ .

**4.4. Hopf subalgebras and Hopf ideals.** Bergbauer and Kreimer showed that, under the assumption that  $B_{\delta}^+$  is a Hochschild cocycle, solutions of Dyson–Schwinger equations determine a sub-Hopf algebra. Namely, they considered Dyson–Schwinger equation in a Hopf algebra  $\mathcal{H}$  of the form

$$(4.7) \quad X = 1 + \sum_{n=1}^{\infty} c_n B_{\delta}^+(X^{n+1}).$$

One then considers the associative algebra  $\mathcal{A}$  defined as the *subalgebra* of  $\mathcal{H}$  generated by the components  $x_n$ , with  $n \geq 0$  of the unique solution of this Dyson–Schwinger equations. In Theorem 3 of [2], they showed that  $\mathcal{A}$  is in fact a *sub-Hopf* algebra, by inductively using the cocycle condition (4.1) to see that

$$(4.8) \quad \Delta(x_n) = \sum_{k=0}^n \Pi_k^n \otimes x_k, \quad \text{where} \quad \Pi_k^n = \sum_{j_1 + \dots + j_{k+1} = n-k} x_{j_1} \cdots x_{j_{k+1}}.$$

This result was extended by Foissy to the more general form (4.3) of Dyson–Schwinger equations in [10] and to systems of Dyson–Schwinger equations in [11]. In Theorem 4 of [10] it is shown that, for an equation of the form (4.3), the subalgebra  $\mathcal{A}$  spanned by the solutions is a Hopf subalgebra if and only if the formal series  $P(t)$  satisfies the differential equation

$$(4.9) \quad (1 - \alpha\beta t)P'(t) = \alpha P(t), \quad \text{with} \quad P(0) = 1,$$

for some  $\alpha, \beta \in \mathbb{K}$ , again under the assumption that  $B^+$  satisfies the cocycle condition (4.1). Similarly, in [11] combinatorial conditions and conditions on the multivariable series  $F_i$  are identified that completely characterize when the solutions of that systems of Dyson–Schwinger equations generate a Hopf algebra, under the assumption that the  $B^+$  satisfies the cocycle condition (4.1).

**4.4.1. Ideals of Dyson–Schwinger solutions.** It is natural to consider not only the associative *subalgebra*  $\mathcal{A}$  spanned by the components of the solutions of Dyson–Schwinger equations, but also the *ideal*  $\mathcal{I}$  spanned by the  $(x_n)_{n \geq 1}$ . As in the case of the subalgebra  $\mathcal{A}$ , it is also natural to ask for conditions that will ensure that it is a Hopf subalgebra, so with the ideal  $\mathcal{I}$  of solutions it is natural to ask for conditions ensuring that it is a Hopf ideal, so that it makes sense to define a quotient Hopf algebra, playing the role of the ring of functions on the variety of solutions.

As an example, we consider here the Dyson–Schwinger equation (4.3), with  $B^+ = B_{\flat}^+$  the grafting operator to a vertex labelled with the bracketing operation.

Assuming that parallel computations that do not feed into each other can be inputted in arbitrary order at the vertex, we can regard this grafting operator as descending to the commutative quotient  $\mathcal{H}_{\text{flow}, \mathcal{V}'}$  of the noncommutative Hopf algebra  $\mathcal{H}_{\text{flow}, \mathcal{V}'}^{nc}$ . This heuristic rationale for passing to the commutative quotient does not take into account a proper definition of a notion of parallel processes, which would require considering the notion of leveled trees and, in the binary case, Hopf algebra structures related to permutohedra instead of associahedra. We refer the reader to [14], [16], [22].

Let  $X = \sum_{n \geq 1} x_n$  be the unique solution of (4.3), obtained as in (4.4), with  $x_1 = B_{\mathfrak{b}}^+(1) = v_{\mathfrak{b}}$  a single vertex labeled with the  $\mathfrak{b}$  operation. We define  $\mathcal{I}$  to be the ideal in the commutative algebra  $\mathcal{H}_{\text{flow}, \mathcal{V}'}$  generated by the  $x_n$  with  $n \geq 1$ .

The same conditions that ensures that  $\mathcal{A}$  is a Hopf subalgebra also ensure that  $\mathcal{I}$  is a Hopf ideal. We verify it in the case of (4.7). The more general case of [10] is similar.

LEMMA 4.2. *The ideal  $\mathcal{I}$  generated by the components  $x_n$  with  $n \geq 1$  of the solution of (4.7) is a Hopf ideal.*

PROOF. Elements of the ideal  $\mathcal{I}$  in  $\mathcal{H} = \mathcal{H}_{\text{flow}, \mathcal{V}'}$  are finite sums  $\sum_{m=1}^M h_m x_{k_m}$ , with  $h_m \in \mathcal{H}$  and  $x_k$  the coordinates of the unique solution of (4.7). The condition that  $\mathcal{I}$  is a Hopf ideal is that  $\Delta(\mathcal{I}) \subset \mathcal{I} \otimes \mathcal{H} \oplus \mathcal{H} \otimes \mathcal{I}$ . Using the formula (4.8) for the coproduct of the elements  $x_k$ , we see that we obtain a sum of terms of which the primitive part  $1 \otimes x_k + x_k \otimes 1$  is in  $\mathcal{H} \otimes \mathcal{I} \oplus \mathcal{I} \otimes \mathcal{H}$  and all the other terms  $\hat{\Delta}(x_k)$  are in  $\mathcal{I} \otimes \mathcal{I}$ . Then the coproducts  $\Delta(h_m x_{k_m})$  will also be in  $\mathcal{H} \otimes \mathcal{I} \oplus \mathcal{I} \otimes \mathcal{H}$ .  $\square$

Notice that we can also work with the noncommutative Hopf algebra  $\mathcal{H}_{\text{flow}, \mathcal{V}'}^{nc}$  and the two sided ideal with elements of the form  $\sum_m h_m x_{k_m} \ell_m$ , with  $h_m, \ell_m \in \mathcal{H}_{\text{flow}, \mathcal{V}'}^{nc}$ , but for the considerations that follow (see §4.5 below) it is more natural to work with the commutative quotient.

**4.5. Yanofsky's Galois theory of algorithms.** As we have seen above, one can pass to commutative quotients  $\mathcal{H}_{\text{flow}}$  of the noncommutative Hopf algebras of flow charts. The meaning of passing to the commutative quotient  $\mathcal{H}_{\text{flow}}$  is best expressed as follows: one can think of monomials in  $\mathcal{H}_{\text{flow}}^{nc}$  as diagrams for parallel computations of a certain number of outputs (one for each planar tree in the forest). Imposing commutativity then corresponds to the reasonable expectation that if one has two parallel computations that do not feed one into the other, then computing them in either order will not affect the result. However, when one introduces operations that graft the output of one tree to the input of another, with  $\mathfrak{c}$  and  $\mathfrak{r}$  labels at the vertices, imposing the commutativity relation is no longer appropriate, as the result of grafting would not be well defined as a *planar* tree. In the case of grafting with  $\mathfrak{c}$  label, we can still pass to the commutative quotient, as argued above.

4.5.1. *Commutative Hopf algebras and affine group schemes.* An advantage of working with commutative Hopf algebras is that they are dual to affine group schemes. Thus, the Hopf algebra  $\mathcal{H}_{\text{flow}}$  determines an affine group scheme  $G_{\text{flow}}$  such that, for any commutative  $\mathbb{K}$ -algebra  $A$ , one obtains a group  $G_{\text{flow}}(A) = \text{Hom}(\mathcal{H}_{\text{flow}}, A)$ , the set of homomorphisms of  $\mathbb{K}$ -algebras, with the product  $(\phi_1 \star$

$\phi_2)(x) = \langle \phi_1 \otimes \phi_2, \Delta(x) \rangle$  dual to the coproduct and the inverse determined by the antipode.

4.5.2. *Galois groups of algorithms.* The idea elaborated by Yanofsky in the recent work [33], of a Galois theory of algorithms, aims at considering all possible “reasonable sets of relations” on programs that would correspond to “implementing the same algorithm”, viewing “algorithms” as an intermediate level between the labelled planar rooted trees and the primitive recursive functions they compute. Our point of view here is similar and directly inspired by the approach proposed by Yanofsky, with the main difference that we seek to implement relations at the level of the Hopf algebra  $\mathcal{H}_{\text{flow}}^{nc}$  of flow charts.

The automorphism groups that arise in Yanofsky’s setting are typically products of symmetric groups and the notion of equivalence relations considered does not take into account any additional structure on flow charts, such as Manin’s Hopf algebra point of view. Moreover, a possible drawback of the current formulation of the approach of [33] lies in the fact that different symmetry groups can have the same orbit space and give rise to the same equivalence on algorithms. We propose instead to consider a different approach where equivalence relations on algorithms are imposed at the level of the Hopf algebra of flow charts.

Our proposal for an alternative approach to a Galois theory of algorithms is to take into account the Hopf algebra structure as part of the construction. By this we mean defining as “reasonable relations” between the planar rooted trees of  $\mathcal{H}_{\text{flow}}^{nc}$  those that correspond to passing to quotient Hopf algebras. Each such quotient is determined by a Hopf ideal  $\mathcal{I}$  in  $\mathcal{H}_{\text{flow}}^{nc}$ , the kernel of the map to the quotient Hopf algebra  $\mathcal{H}_{\mathcal{I}}$ .

In the commutative setting  $\mathcal{H}_{\text{flow}}$  one can view, dually, the quotient Hopf algebras  $\mathcal{H}_{\mathcal{I}}$  as sub-group schemes  $G_{\mathcal{I}}$  of the affine group scheme of characters  $G_{\text{flow}}$ . The group schemes  $G_{\mathcal{I}}$  obtained in this way replace, in our approach, the basic objects of interest in Yanofsky’s Galois theory of algorithms.

4.5.3. *Hopf ideals and Galois groups.* In light of Yanofsky’s point of view, it is natural to consider ideals in  $\mathcal{H}_{\text{flow}}$  generated by solutions of Dyson–Schwinger equations, and the resulting quotient Hopf algebras, rather than sub-Hopf algebras. The Hopf ideals and quotient Hopf algebras define, dually, the affine group schemes that provide our proposed notion of Galois groups of algorithms.

As a heuristic geometric interpretation, in the Hopf algebraic setting, we can think of the quotient  $\mathcal{H}/\mathcal{I}$  of a Hopf algebra by a Hopf ideal as an analog of the algebro-geometric quotient  $R/I$  of a polynomial ring  $R$  by a prime ideal  $I$ , which provides the coordinate ring of the affine variety  $V$  defined by  $I$ . In the case of the Hopf ideal  $\mathcal{I}$  determined by the components of the solution of the Dyson–Schwinger equation, we can therefore think of  $\mathcal{H}/\mathcal{I}$  as the “ring of function” of the variety of solutions. The associated group scheme  $G_{\mathcal{I}}$  consists of those symmetries in  $G_{\text{flow}}$  that are compatible with the self-similarity structure described by the Dyson–Schwinger equation.

PROPOSITION 4.3. *A Dyson–Schwinger equation (4.7) in the commutative Hopf algebra  $\mathcal{H}_{\text{flow}}$  determines a Galois group  $G_{\mathcal{P}} \subset G_{\text{flow}}$ .*

This is simply a reformulation of Lemma 4.2, with our notion of what we consider to be an appropriate notion of Galois group. A similar result holds for

Dyson–Schwinger equations (4.3) with (4.9). While, on the one hand, our point of view takes into account more information, coming from the Hopf algebra structure, at another level it is less structured than Yanofsky’s setting: in fact, we work here with the Hopf algebra where flow charts only have vertex decorations, so that the input functions (and the corresponding edge labelings through the graph) are not specified, while in Yanofsky’s setting the underlying functions are assigned and preserved by the symmetries.

**4.6. A non-Hopf example with binary trees.** In the setting of Dyson–Schwinger equations described above, it is crucial that we consider the planar rooted trees in  $\mathcal{H}_{\text{flow}, \mathcal{V}}^{nc}$  without external edges. In fact, this allows us to define the grafting operators  $B_\delta^+$  as grafting any number of trees to the *same* root consisting of a single vertex decorated by  $\delta$ . If we reintroduce external edges, then the graphs consisting of a single vertex but with different numbers of incoming edges are viewed as different generators in the Hopf algebra and they no longer define a cocycle  $B_\delta^+$ .

To see this more explicitly, we focus on the case of binary trees, keeping track of external edges. Consider the following primitive elements in the Hopf algebra  $\mathcal{H}_{\text{flow}, \mathcal{V}}^{nc, 0}$ :

$$(4.10) \quad \begin{array}{c} \diagup \\ \bullet \\ \diagdown \\ | \\ \tau \end{array}, \quad \begin{array}{c} \diagup \\ \bullet \\ \diagdown \\ | \\ \mathfrak{b} \end{array}, \quad \begin{array}{c} \diagup \\ \bullet \\ \diagdown \\ | \\ \mathfrak{c} \end{array}$$

For each of these graphs  $\tau$  one can define a grafting operation  $B_\tau^+$  by setting the result equal to zero on monomials  $T$  that cannot be grafted to  $\tau$ . This  $B_\tau^+$  clearly does *not* satisfy the cocycle condition because one would have to have  $B_\tau^+(1) = 0$ , but then the expressions  $\tilde{\Delta} B_\tau^+$  and  $(id \otimes B_\tau^+) \tilde{\Delta}$ , when applied to elements  $T = \tau_1 \amalg \tau_2$  that are not in the kernel, would differ by a term corresponding to the cut that separates  $\tau$  from  $T$ . On the other hand, if  $B_\tau^+(1) = \tau$  as this case would require, then the cocycle condition fails on elements in the kernel of  $B_\tau^+$ .

Another way of defining a grafting operator  $B_\tau^+$  with  $\tau$  one of the trees (4.10) is to define, for a monomial  $T = \tau_1 \amalg \cdots \amalg \tau_n$  the grafting  $B_\tau^+(T)$  to be obtained by just gluing one of the outputs of  $T$  (say, the first one) to one of the two inputs of  $\tau$  (say, the first one), with the result of grafting taken to be  $\tau$  itself if  $T = \emptyset$ . We use this choice in the explicit example below. Again, however, one sees that the grafting defined in this way does *not* satisfy the cocycle condition. In fact, consider the case of a monomial  $T$  consisting of a forest with more than one component. Then only one of the components of  $T$  is glued to  $\tau$  by  $B_\tau^+$ . This means that, in the computation of the term  $\tilde{\Delta} B_\tau^+$  there is not just one but several admissible cuts that separate  $\tau$  from  $T$ , namely all the admissible cuts that include a cut of the edge used for grafting together with an arbitrary cut of the remaining components not connected to  $\tau$ . Thus, the difference between  $\tilde{\Delta} B_\tau^+$  and  $(id \otimes B_\tau^+) \tilde{\Delta}$  does not consist only of the term  $id \otimes B_\delta^+(1) = id \otimes \tau$ .

However, one can still consider a Dyson–Schwinger equation of the form

$$X = 1 + \sum_{\delta} B_{\tau_\delta}^+(X^2),$$

where the sum is on the three labels  $\delta \in \{\mathbf{b}, \mathbf{c}, \mathbf{r}\}$  of the vertex of the tree  $\tau_\delta$  as in (4.10). This admits an explicit solution, as in the general case above, which are of the form

$$x_{n+1} = \sum_{k=0}^n \sum_{\delta} B_{\tau_\delta}^+(x_k x_{n-k}).$$

In  $\overline{\mathcal{H}}_{\text{flow}, \nu}^0$ , one can see that the first few terms are of the form  $x_0 = I$  and

$$\begin{aligned}
 x_1 &= \begin{array}{c} \diagup \\ \bullet \mathbf{r} \\ \diagdown \\ | \end{array} + \begin{array}{c} \diagup \\ \bullet \mathbf{b} \\ \diagdown \\ | \end{array} + \begin{array}{c} \diagup \\ \bullet \mathbf{c} \\ \diagdown \\ | \end{array} \\
 x_2 &= \sum_a B_+^a(2x_1 x_0) \\
 &= 2 \begin{array}{c} \diagup \\ \bullet \mathbf{r} \\ | \\ \bullet \mathbf{r} \\ \diagdown \\ | \end{array} + 2 \begin{array}{c} \diagup \\ \bullet \mathbf{b} \\ | \\ \bullet \mathbf{r} \\ \diagdown \\ | \end{array} + 2 \begin{array}{c} \diagup \\ \bullet \mathbf{c} \\ | \\ \bullet \mathbf{r} \\ \diagdown \\ | \end{array} \\
 &+ \begin{array}{c} \diagup \\ \bullet \mathbf{r} \\ | \\ \bullet \mathbf{b} \\ \diagdown \\ | \end{array} + 2 \begin{array}{c} \diagup \\ \bullet \mathbf{b} \\ | \\ \bullet \mathbf{b} \\ \diagdown \\ | \end{array} + 2 \begin{array}{c} \diagup \\ \bullet \mathbf{c} \\ | \\ \bullet \mathbf{b} \\ \diagdown \\ | \end{array} \\
 &+ \begin{array}{c} \diagup \\ \bullet \mathbf{r} \\ | \\ \bullet \mathbf{c} \\ \diagdown \\ | \end{array} + 2 \begin{array}{c} \diagup \\ \bullet \mathbf{b} \\ | \\ \bullet \mathbf{c} \\ \diagdown \\ | \end{array} + 2 \begin{array}{c} \diagup \\ \bullet \mathbf{c} \\ | \\ \bullet \mathbf{c} \\ \diagdown \\ | \end{array}
 \end{aligned}$$

One can continue with the successive coefficients, such as  $x_3 = \sum_a B_+^a(x_1^2 + 2x_2x_0)$ , using the facts that for a vertex labeled by recursion, the order of inputs always matters, as well as for composition, while we can assume that order of inputs does not matter for bracketing as a consequence of the nature of parallel computation.

This example shows that, if we want to work with binary trees with external edges, which is a very natural choice from the point of view of computation theory, and we want to obtain a grafting operation that is non-trivial on monomials of arbitrary degree, we would have to extend the grafting to include, besides the

primitive two-valent trees (4.10), also a choice of binary trees with larger numbers of input flags. This cannot be accommodated in the usual setting of combinatorial Dyson–Schwinger equations based on Hochschild cocycles in Hopf algebras. However, we shall see below that such generalizations exist naturally when one reformulates Dyson–Schwinger equations in the operadic setting.

## 5. Operadic viewpoint

Instead of using a Hopf algebraic setup to describe flow charts and primitive recursive functions, one can also adopt an operadic viewpoint, as suggested by Manin in [27].

Let us consider, as above, planar rooted trees  $\tau$  with vertex labeling by the elementary operations of type  $\{\mathbf{b}, \mathbf{c}, \mathbf{r}\}$ , oriented so that each tree has a single outgoing flag attached to the root vertex and a certain number of incoming flags.

**DEFINITION 5.1.** *We define the operad of flow charts  $\mathcal{O}_{\text{flow}, \mathcal{V}}$  by setting  $\mathcal{O}(n)$  to the  $\mathbb{K}$ -vector space spanned by labelled planar rooted trees with  $n$  incoming flags and the operad composition operations*

$$\circ_{\mathcal{O}} : \mathcal{O}(n) \otimes \mathcal{O}(m_1) \otimes \cdots \otimes \mathcal{O}(m_n) \rightarrow \mathcal{O}(m_1 + \cdots + m_n)$$

are given on generators  $\tau \otimes \tau_1 \otimes \cdots \otimes \tau_n$  by grafting the output flag of the tree  $\tau_i$  to the  $i$ -th input flag of the tree  $\tau$ .

As pointed out in [2], Dyson–Schwinger equations admit a natural operadic interpretation. Namely, given a formal series  $P(t) = 1 + \sum_{k=1}^{\infty} a_k t^k$  and a collection  $\beta = (\beta_n)$  with  $\beta_n \in \mathcal{O}(n)$ , we consider the equation

$$(5.1) \quad X = \beta(P(X)),$$

where  $X = \sum_k x_k$  is a formal sum with  $x_k \in \mathcal{O}(k)$ . In the right hand side we have  $\beta(P(X))_1 = 1 + \beta_1 \circ x_1$ , where 1 is the identity in  $\mathcal{O}(1)$ , and for  $n \geq 2$

$$(5.2) \quad \beta(P(X))_n = \sum_{k=1}^n \sum_{j_1 + \cdots + j_k = n} a_k \beta_k \circ (x_{j_1} \otimes \cdots \otimes x_{j_k}),$$

with  $x_{j_1} \otimes \cdots \otimes x_{j_k} \in \mathcal{O}(j_1) \otimes \cdots \otimes \mathcal{O}(j_k)$ , so that the composition  $\beta_k \circ_{\mathcal{O}} (x_{j_1} \otimes \cdots \otimes x_{j_k}) \in \mathcal{O}(n)$ , since  $j_1 + \cdots + j_k = n$ .

**PROPOSITION 5.2.** *For  $\mathcal{O} = \mathcal{O}_{\text{flow}, \mathcal{V}}$  the operad of flow charts, if  $a_1 \beta_1 \neq 1 \in \mathcal{O}(1)$ , the operadic Dyson–Schwinger equation (5.1) has a unique solution  $X \in \prod_{n \geq 1} \mathcal{O}(n)$  given inductively by*

$$(5.3) \quad (1 - a_1 \beta_1) \circ x_{n+1} = \sum_{k=2}^{n+1} \sum_{j_1 + \cdots + j_k = n+1} a_k \beta_k \circ (x_{j_1} \otimes \cdots \otimes x_{j_k}).$$

**PROOF.** First observe that in the operad of flow charts  $\mathcal{O}(1)$  is one-dimensional, since it is spanned by the trees with a single vertex, one incoming and one outgoing flag, with the operations  $\mathbf{b}$ ,  $\mathbf{c}$  or  $\mathbf{r}$  at the vertex, but in the case of a single input, all of these operations are the identity, so that elements of  $\mathcal{O}(1)$  are scalar multiples of the identity map, viewed as operations on primitive recursive functions. Thus, the element  $\beta_1 \in \mathcal{O}(1)$  is a scalar multiple of the identity and, if we assume that  $a_1 \beta_1 \neq 1$  in  $\mathcal{O}(1)$ , then  $1 - a_1 \beta_1$  is invertible. Since we have  $\beta(P(X))_1 = 1 + \beta_1 \circ x_1$

as the term in  $\mathcal{O}(1)$  in the right hand side of (5.1), we obtain  $(1 - a_1\beta_1)x_1 = 1$ , which fixes  $x_1 = (1 - a_1\beta_1)^{-1} \in \mathcal{O}(1)$ . At the next step we have, from (5.2),  $x_2 = a_1\beta_1 \circ x_2 + a_2\beta_2 \circ (x_1 \otimes x_1)$ , which gives  $x_2 = (1 - a_1\beta_1)^{-1}(a_2\beta_2 \circ (x_1 \otimes x_1))$ , with  $x_1$  as above. At the  $(n + 1)$ -st step,  $x_{n+1}$  is then determined uniquely by (5.3) in terms of the coefficients  $a_k$  and the elements  $\beta_k$  for  $k = 1, \dots, n + 1$ .  $\square$

Let then  $\mathcal{O}_{\beta,P}(n)$  denote the  $\mathbb{K}$ -linear span of all the compositions  $x_k \circ (x_{j_1} \otimes \dots \otimes x_{j_k})$ , for  $k = 1, \dots, n$  and  $j_1 + \dots + j_k = n$ , with  $x_k$  the coordinates of the solution  $X = \sum_k x_k$  of the Dyson–Schwinger equation determined by  $\beta = (\beta_n)$  and  $P(t)$ . The  $\mathcal{O}_{\beta,P}(n)$  form a sub-operad under the composition maps induced from  $\mathcal{O}(n)$ .

For instance, for a choice of one of the operations  $\delta \in \{\mathbf{b}, \mathbf{c}, \mathbf{r}\}$ , we can take  $a_1 \neq 1$  and the element  $\beta_k$  given by the tree with a single vertex marked by  $\delta$ , with  $k$  incoming flags and one outgoing flag. This choice gives an operadic reformulation of the Hopf-theoretic Dyson–Schwinger equations with the cocycles  $B_\delta^+$ . However, we can now consider also more general operadic Dyson–Schwinger equations for different choices of  $\beta = (\beta_n)$  that do not correspond to Hochschild cocycles in the Hopf algebra setting.

Restricting to the sub-operad  $\mathcal{O}_{\beta,P}(n)$  generated by solutions of the Dyson–Schwinger equation should be regarded as considering those operations (formal combinations of flow charts) that satisfy a self-similarity property with respect to the transformation  $X \mapsto \beta(P(X))$ .

5.0.1. *The case of binary trees.* In particular, we can now revisit the example of the binary trees in this operadic setting. Let  $\mathcal{O}_{\text{flow},\mathcal{V}}^0(n)$  be the  $\mathbb{K}$ -linear space spanned by the planar *binary* trees with vertex labeling by the binary operations  $\mathbf{b} = \mathbf{b}_2, \mathbf{r}, \mathbf{c}$ . These form a sub-operad of  $\mathcal{O}_{\text{flow},\mathcal{V}}$  with the induced composition maps.

For a fixed  $\delta \in \{\mathbf{b}, \mathbf{r}, \mathbf{c}\}$  let  $\beta_{1,\delta}$  be the binary tree with a single vertex, two incoming and one outgoing flag, as in (4.10). We then consider a choice  $\beta_n$  of a binary tree with  $n$  input flags, for each  $n \geq 1$ , with a fixed choice of vertex labelings in  $\{\mathbf{b}, \mathbf{r}, \mathbf{c}\}$ , not necessarily all of the same type. Each such choice of  $\beta = (\beta_n)$ , together with a choice of  $P(t)$ , now determines an operadic Dyson–Schwinger equation of the form (5.1), in the binary setting, with a corresponding solution as in Proposition 5.2.

5.0.2. *Systems of Dyson–Schwinger equations in operads.* By analogy to what happens in Hopf algebras, [11], one can pass from the case of a single Dyson–Schwinger equation to systems of Dyson–Schwinger equations. In the operadic setting a system of Dyson–Schwinger equations would be determined by the data of formal series  $F_\delta(t_1, t_2, t_3) = \sum_{k_1, k_2, k_3} a_{k_1, k_2, k_3}^{(\delta)} t_1^{k_1} t_2^{k_2} t_3^{k_3}$  and of elements  $\beta^{(\delta)} = (\beta_n^{(\delta)})$ , where for  $\delta \in \{\mathbf{b}, \mathbf{r}, \mathbf{c}\}$ , the elements  $\beta_n^{(\delta)} \in \mathcal{O}(n)$  are realized by rooted trees with root vertex marked by  $\delta$ . The operadic systems of Dyson–Schwinger equations are then of the form

$$X_\delta = \beta^{(\delta)}(F_\delta(X)).$$

The explicit form of the solutions is more cumbersome than in the case of a single equation in Proposition 5.2 above, and it follows a pattern similar to the derivation of the explicit solutions of Hopf theoretic systems of Dyson–Schwinger equations (4.5) as in [11]. These systems detect more elaborate forms of self-similarity in the

operad of flow charts, involving different families of grafting operations simultaneously.

5.0.3. *Operads and Properads.* As mentioned in [27], one can extend the Hopf algebra of flow charts by considering, instead of labelled embedded rooted trees, more general labelled embedded acyclic graphs, namely graphs endowed with an acyclic orientation. Correspondingly, the operad of flow charts would be replaced by a properad, where the compositions extend from grafting output and input flags of trees to grafting outputs and inputs of acyclic graphs.

The notion of *properad* was introduced in [31] as an intermediate notion between operad and prop. Namely, a properad parameterizes operations with varying numbers of inputs and outputs that are labelled by connected acyclic graphs, while the operad case uses trees (varying number of inputs and a single output) and props allow for disconnected graphs.

More precisely, we consider the properad  $\mathcal{P}_{\text{flow}, \mathcal{V}}$  where  $\mathcal{P}(m, n)$  is the  $\mathbb{K}$ -vector space spanned by planar connected directed (acyclic) graphs with  $m$  incoming flags and  $n$  outgoing flags, with vertices decorated by operation that include the elementary  $\mathfrak{b}$ ,  $\mathfrak{c}$ ,  $\mathfrak{r}$ , with  $m$  inputs and one output, but also include now a chosen set of additional operations with  $m$  inputs and  $n$  outputs. As we have seen in the case of extending the operations  $\mathfrak{c}$  and  $\mathfrak{r}$  from binary to  $m$ -ary inputs, in this case also such additional operations with  $m$  inputs and  $n$  outputs acting on primitive recursive functions can be decomposed in terms of the elementary operations: the resulting operations associated to vertices with  $m$  incoming and  $n$  outgoing flags can be regarded as “macros”, in the sense discussed in §2 of [33]. The properad has composition operations

$$\mathcal{P}(m, n) \otimes \mathcal{P}(j_1, k_1) \otimes \cdots \otimes \mathcal{P}(j_\ell, k_\ell) \rightarrow \mathcal{P}(j_1 + \cdots + j_\ell, n), \quad \text{for } k_1 + \cdots + k_\ell = m.$$

5.0.4. *Dyson–Schwinger equations in properads.* We can define Dyson–Schwinger equations for properads, extending the operadic case considered above. Namely we consider again a choice of a formal power series  $P(t) = 1 + \sum_k a_k t^k$  and a collection  $\beta = (\beta_{m,n})$  of elements with  $\beta_{m,n} \in \mathcal{P}(m, n)$ . We then consider the equation

$$(5.4) \quad X = \beta(P(X))$$

as in (5.1), where now the right hand side has  $(m, n)$ -component in  $\mathcal{P}(m, n)$  given by

$$(5.5) \quad \beta(P(X))_{m,n} = \sum_{k=1}^m a_k \sum_{\substack{j_1 + \cdots + j_k = m \\ i_1 + \cdots + i_k = \ell}} \beta_{\ell,n} \circ (x_{j_1, i_1} \otimes \cdots \otimes x_{j_k, i_k}).$$

In order to construct solutions of the properad Dyson–Schwinger equations, we introduce transformations  $\Lambda_n = \Lambda_n(a, \beta)$  with

$$\Lambda_n(a, \beta) : \oplus_{k=1}^n \mathcal{P}(n, k) \rightarrow \oplus_{k=1}^n \mathcal{P}(n, k), \quad \text{with } \Lambda_n(a, \beta)_{ij} = a_j \beta_{j,i}.$$

**THEOREM 5.3.** *If for all  $n \geq 1$  the transformation  $I - \Lambda_n(a, \beta)$  is invertible, with  $I$  the identity on  $\oplus_{k=1}^n \mathcal{P}(n, k)$ , then the properad Dyson–Schwinger equation*

(5.4) admits a unique solution, given by  $x_{1,1} = \Lambda_1^{-1}$ , and for  $m < n$  by

$$(5.6) \quad x_{m,n} = \sum_{k=1}^m a_k \beta_{k,n} \circ \left( \sum_{\ell=1}^k \sum_{\substack{j_1+\dots+j_\ell=m \\ i_1+\dots+i_\ell=k}} x_{j_1,i_1} \otimes \dots \otimes x_{j_\ell,i_\ell} \right),$$

while the remaining components  $x_{m,n}$  with  $m \geq n$  are determined by

$$(5.7) \quad Y_n(x) = (I - \Lambda_n)^{-1} \Lambda_n V^{(n)}(x),$$

where  $Y_n(x)^t = (x_{n,1}, \dots, x_{n,n})$  and  $V^{(n)}(x)^t = (V^{(n)}(x)_j)_{j=1,\dots,n}$

$$V^{(n)}(x)_j = \sum_{k=2}^n \sum_{\substack{r_1+\dots+r_k=n \\ s_1+\dots+s_k=j}} x_{r_1,s_1} \otimes \dots \otimes x_{r_k,s_k}.$$

PROOF. By (5.5) we see that we have  $x_{1,1} = 1 + a_1 \beta_{1,1} \circ x_{1,1}$ , which can be solved for  $x_{1,1}$  provided the invertibility of  $1 - \Lambda_1 = 1 - a_1 \beta_{1,1}$  holds. The components  $x_{m,n}$  with  $m < n$  are determined by (5.6), where the right hand side only involves components  $x_{r,s}$  with  $r < m$  and  $s < n$ . The equations for the remaining entries  $x_{m,n}$  with  $n \leq m$  break into systems of the form

$$(I - \Lambda_m) \begin{pmatrix} x_{m,1} \\ \vdots \\ x_{m,m} \end{pmatrix} = \sum_j (\Lambda_m)_{ij} \begin{pmatrix} \sum_{k=2}^m \sum_{\substack{r_1+\dots+r_k=m \\ s_1+\dots+s_k=j}} x_{r_1,s_1} \otimes \dots \otimes x_{r_k,s_k} \\ \vdots \end{pmatrix}_j.$$

Provided that the invertibility condition for the transformation  $I - \Lambda_m$  holds, these systems provide unique values for all the components  $x_{m,n}$  with  $n \leq m$ , as a function of the components  $x_{r,s}$  with  $r < m$ , that we have already been determined.  $\square$

Notice that one does not expect the invertibility of  $I - \Lambda_n(a, \beta)$  to hold in general, as  $\Lambda_n(a, \beta)$  involves compositions with the elements  $\beta_{j,i}$ . The simplest example in which one has invertibility can be obtained by taking  $\beta_{m,n} = 0$  for  $m \neq n$  and  $I - \beta_{n,n}$  an invertible transformation in  $\mathcal{P}(n, n)$ .

As in the operad case, we can define  $\mathcal{P}_{\beta,P}(m, n)$  to be the  $\mathbb{K}$ -vector spaced spanned by all the properad compositions

$$x_{k,n} \circ (x_{j_1,i_1} \otimes \dots \otimes x_{j_\ell,i_\ell})$$

with  $j_1 + \dots + j_\ell = m$  and  $i_1 + \dots + i_\ell = k$  and where all the  $x_{m,n}$  are components of the solution to the properad Dyson–Schwinger equation determined by the choice of  $P(t)$  and of  $\beta = (\beta_{m,n})$ , with the invertibility condition of Theorem 5.3. The  $\mathcal{P}_{\beta,P}(m, n)$  form a sub-properad of  $\mathcal{P}_{\text{flow}, \mathcal{V}}$  with the induced composition.

## 6. Renormalization of the halting problem

In [25], Manin adapted the formalism of perturbative renormalization to the halting problem in the theory of computation. To summarize Manin’s approach, the basic idea is that one considers non-computable functions as an analog of divergent Feynman integrals. Renormalization consists of a procedure of extraction of finite values from divergences and in this setting Manin proposes that a suitable form of renormalization should lead to a procedure of extraction of a “computable part” from a non-computable function.

While this whole program is, at present, still in its early stages and some of the ideas involved are still speculative, we attempt here to give a sense of the general picture and formulate some of the main questions, as we see them. We will also discuss how Dyson–Schwinger equations may be accommodated in this general circle of ideas. While most of what is discussed in this section is more speculative than the previous parts of the paper, we regard this viewpoint as an important motivation for further developments.

**6.1. The halting problem: regularization and renormalization.** We describe here a setting for renormalization of the halting problem which is a minor elaboration on the procedure described in §3 of [25]. We do not add any significant new development beyond what is outlined as a possible program in Manin’s work [25], but we will try to identify some specific questions that we hope to return to in future work.

6.1.1. *Partial recursive functions and Hopf algebras.* We first pass from the class of primitive recursive functions considered above to the larger class of partial recursive functions. These have a similar presentation as the one recalled in §2 above, with the same basic functions (successor, constants, and projections) and, in addition to the three elementary operations  $\mathfrak{c}$ ,  $\mathfrak{b}$ ,  $\mathfrak{r}$  of composition, bracketing and recursion, an additional  $\mu$  operation that assigns to an input function  $f : \mathbb{N}^{n+1} \rightarrow \mathbb{N}$  assigns an output

$$h : \mathbb{N}^n \rightarrow \mathbb{N}, \quad h(x_1, \dots, x_n) = \min\{x_{n+1} \mid f(x_1, \dots, x_{n+1}) = 1\},$$

with domain

$$\mathcal{D}(h) = \{(x_1, \dots, x_n) \mid \exists x_{n+1} \geq 1 : f(x_1, \dots, x_{n+1}) = 1\},$$

with  $(x_1, \dots, x_n, k) \in \mathcal{D}(f), \forall k \leq x_{n+1}$ . According to Church’s thesis the partial recursive functions obtained in this way are exactly the semi-computable functions, namely those for which there exists a program that, for all  $x \in \mathcal{D}(f)$ , computes  $f(x)$  and which either computes zero or runs for an infinite time when  $x \notin \mathcal{D}(f)$ , see §V of [23]. Notice that it is the presence of the additional operation  $\mu$  that allows for the construction, starting from the basic functions, of functions that are only partially defined and semi-computable.

Correspondingly, we enrich the Hopf algebras of flow charts by additional vertex decorations by the  $\mu$  operations. As in the case of the other operations, this requires defining a version of the  $\mu$  operation for vertices of valence greater than two. For example, a possible choice may be to use the higher-arity bracketing operation and an application of  $\mu$  coordinatewise. We shall still use the notations  $\mathcal{H}_{\text{flow}, \mathcal{R}'}^{nc}$  and  $\mathcal{H}_{\text{flow}, \mathcal{V}'}^{nc}$ , for the resulting Hopf algebras with admissible decorations by partial recursive functions or with vertex decorations, respectively, and  $\mathcal{H}_{\text{flow}, \mathcal{R}'}$  and  $\mathcal{H}_{\text{flow}, \mathcal{V}'}$  for their commutative quotients.

6.1.2. *Partial recursive functions and the halting problem.* As in [25], we regard the halting problem as a particular instance of the general problem of recognizing whether a given  $k \in \mathbb{N}$  belongs to the domain of definition  $\mathcal{D}(f)$  of a partial recursive function  $f$ . The halting problem corresponds to the case where, given a partial recursive function  $f$  and a program  $P$  computing it, one considers the partial recursive function  $t(P)$ , the computation time, with domain  $\mathcal{D}(t(P)) = \mathcal{D}(f)$ , see [25], §3.1.(a) and §2.1, and also §IX of [23].

6.1.3. *Algebraic renormalization.* The basic algebraic formalism for perturbative renormalization, in its formulation as given in [9] (see also §5 of [28]), consists of an *algebraic Feynman rule*

$$(6.1) \quad \phi : \mathcal{H} \rightarrow \mathcal{B}$$

which is a *morphism of commutative algebras* from a *commutative Hopf algebra*  $\mathcal{H}$  to a *Rota–Baxter algebra*  $\mathcal{B}$  of weight  $-1$ .

Recall that a commutative Rota–Baxter algebra of weight  $\lambda$  is a commutative, associative algebra  $\mathcal{B}$  endowed with a linear operator  $T$  on it that satisfies the identity

$$(6.2) \quad T(x)T(y) = T(xT(y)) + T(T(x)y) + \lambda T(xy).$$

The field of convergent Laurent series in one variable with the projection  $T$  onto the polar part is an example of Rota–Baxter algebra of weight  $\lambda = -1$ , which is widely used in renormalization theory, [3], [5]. The Rota–Baxter operator  $T$  of weight  $-1$  determines a splitting of  $\mathcal{B}$  into two commutative unital subalgebras,  $\mathcal{B}_+ = (1 - T)\mathcal{B}$  and  $\mathcal{B}_-$  given by  $T\mathcal{B}$  with a unit adjoined. The fact that  $\mathcal{B}_\pm$  are indeed algebras and not just vector spaces follows from the Rota–Baxter identity (6.2).

The BPHZ renormalization of the algebraic Feynman rule (6.1) is then obtained by showing the existence of a multiplicative factorization

$$\phi = (\phi_- \circ S) \star \phi_+, \quad \text{with} \quad \phi_\pm : \mathcal{H} \rightarrow \mathcal{B}_\pm,$$

obtained inductively via the BPHZ preparation formula (see [3], [9])

$$(6.3) \quad \phi_-(x) = -T(\phi(x) + \sum \phi_-(x')\phi(x'')), \text{ and } \phi_+(x) = (1 - T)(\phi(x) + \sum \phi_-(x')\phi(x'')),$$

where  $x \in \mathcal{H}$ , with  $\Delta(x) = 1 \otimes x + x \otimes 1 + \sum x' \otimes x''$ . The factorization is unique if normalized by  $\epsilon_- \circ \phi_- = \epsilon$ , where  $\epsilon_- : \mathcal{B}_- \rightarrow \mathbb{C}$  is the augmentation map and  $\epsilon$  is the counit of  $\mathcal{H}$ . Again, it is the Rota–Baxter property of  $(\mathcal{B}, T)$  that ensures that  $\phi_\pm$  obtained in this way are algebra homomorphisms. The homomorphism  $\phi_+ : \mathcal{H} \rightarrow \mathcal{B}_+$  is the *renormalized* algebraic Feynman rule, while  $\phi_- : \mathcal{H} \rightarrow \mathcal{B}_-$  is the *counterterm*, the divergence.

6.1.4. *An algebraic Feynman rule for the halting problem.* We recall here a proposal made by Manin in [25] of a possible construction of an algebraic Feynman rule in the setting of the halting problem.

We let  $\mathcal{B}$  be the algebra of functions  $\Phi : \mathbb{N}^m \rightarrow \mathcal{M}(D)$  from  $\mathbb{N}^m$ , for some  $k$ , to the algebra  $\mathcal{M}(D)$  of analytic functions in the unit disk  $D = \{z \in \mathbb{C} : |z| < 1\}$ . Since the functions  $\Phi$  can have different domains  $\mathbb{N}^k$ , the sums and products are defined by extending them to functions  $\Phi(k_1, k_2, \dots, k_n, \dots)$  from a common domain  $\mathbb{N}^\infty$ , which depend on only finitely many variables,  $\Phi = \Phi \circ \pi_m^\infty$ , with  $\pi_m^\infty : \mathbb{N}^\infty \rightarrow \mathbb{N}^m$ . The Rota–Baxter operator  $T$  on  $\mathcal{B}$  is given by componentwise (for each  $\underline{k}$ ) projection onto the polar part at  $z = 1$ . Thus,  $\mathcal{B}_+$  consists of functions from  $\mathbb{N}^\infty$  (depending only on finitely many  $k_i$ ) to  $\mathcal{M}(D)_+$ , the algebra of meromorphic functions in  $D$  that extend continuously at  $z = 1$ , while  $\mathcal{B}_-$  consists of functions that are either constant or have a pole at  $z = 1$ .

We define algebraic Feynman rules for the Hopf algebras  $\mathcal{H}_{\text{flow}, \mathcal{R}'}$  and  $\mathcal{H}_{\text{flow}, \mathcal{V}'}$ . In the case of  $\mathcal{H}_{\text{flow}, \mathcal{R}'}$ , where both vertices and flags in the trees are labeled by

admissible labelings, there is a well defined partial recursive function  $f$  associated to each labelled rooted tree  $\tau$  in  $\mathcal{H}_{\text{flow}, \mathcal{R}'}$ , namely the function  $f$  labeling the unique outgoing flag of  $\tau$ . Let us first assume that the partial recursive function  $f$  has range in  $\mathbb{N}$ . We extend this  $f : \mathbb{N}^m \rightarrow \mathbb{N}$  to a function  $\bar{f} : \mathbb{N}^m \rightarrow \mathbb{Z}_{\geq 0}$  that computes  $f(x)$  at  $x \in \mathcal{D}(f)$  and takes value 0 at  $x \notin \mathcal{D}(f)$ .

We assign to a tree  $\tau$  that computes  $f$  an element in  $\Phi_\tau(\underline{k}, z)$  in  $\mathcal{B}$ , defined (as in §3 of [25]) by

$$(6.4) \quad \Phi_\tau(\underline{k}, z) = \Phi(\underline{k}, f, z) := \sum_{n \geq 0} \frac{z^n}{(1 + n\bar{f}(\underline{k}))^2}.$$

The resulting function  $\Phi_\tau(\underline{k}, \cdot) \in \mathcal{M}(D)$  has a pole at  $z = 1$  iff  $\underline{k} \notin \mathcal{D}(f)$ . We extend this definition multiplicatively to cases where the output of the flow chart  $\tau$  is a partial recursive function  $f : \mathbb{N}^k \rightarrow \mathbb{N}^\ell$ , by setting  $\Phi_\tau(\underline{k}, z) = \prod_{j=1}^\ell \Phi(\underline{k}, f_j, z)$ . In this way, the resulting function has a pole at  $z = 1$  if at least one of the  $f_j$  does, that is, iff  $\underline{k} \notin \mathcal{D}(f_1) \cap \dots \cap \mathcal{D}(f_\ell) = \mathcal{D}(f)$ , consistently with the above. We then extend the definition multiplicatively to the case of a monomial  $\tau_1 \cdots \tau_n$  in  $\mathcal{H}_{\text{flow}, \mathcal{R}'}$  and additively to linear combinations of monomials. Thus, we obtain an algebra homomorphism  $\Phi : \mathcal{H}_{\text{flow}, \mathcal{R}'} \rightarrow \mathcal{B}$ .

We write  $\Phi(\underline{k}, f, z) := \Phi_f(\underline{k})(z)$  for the function  $\Phi_f(\underline{k}) \in \mathcal{M}(D)$  of  $z \in D$ , associated to a given  $\underline{k} \in \mathbb{N}^\infty$  and depending on the output  $f$  of a given flow chart, with  $f$  a partial recursive functions. If we denote by  $\tau$  the decorated planar rooted tree describing the flowchart, we equivalently write  $\Phi(\underline{k}, \tau, z)$  or  $\Phi_{\underline{k}, \tau}(z)$ .

The specific form (6.4) of the proposed Feynman rule does not appear to be so important, as much as the two properties:

- The function  $\Phi(\underline{k}, f, \cdot) \in \mathcal{M}(D)$  has a pole at  $z = 1$  iff  $\underline{k} \notin \mathcal{D}(f)$
- The function  $\Phi(\underline{k}, f, \cdot)$  depends on  $\underline{k}$  through the values  $\bar{f}(\underline{k})$  and these values can be reconstructed from the function  $\Phi(\underline{k}, f, z)$ .

Any other choice of a function with these properties would be a suitable Feynman rule according to [25]. Thus, it would seem perhaps more convenient to consider the set of all such functions, possibly up to a suitable equivalence relation. We will discuss more below why it may be necessary to introduce an equivalence relation on Feynman rules. This has an analog in the quantum field theory setting, where the classification of divergences in terms of a Galois theory based on the Tannakian formalism was obtained in [6] precisely by introducing a suitable equivalence relation on algebraic Feynman rules. We will not develop this approach further in the present paper, but we hope to return to it elsewhere.

In Manin's idea a possible renormalization of the halting problem should provide a way of extracting a computable part from the halting problem for partial recursive functions that accounts for "subdivergences" created by subroutines within the flowcharts. Though, at present, how exactly this should be achieved is not yet fully understood, in the procedure described in [25] one would proceed by applying the BPHZ factorization (6.3) to the chosen algebraic Feynman rule. The result of factorization consists of two new algebra homomorphisms from the same Hopf algebra to the two parts of the target Rota–Baxter algebra. The main question is then the interpretation of both pieces in terms of computability and partial recursive functions.

When we apply the BPHZ procedure to  $\Phi$  defined as in (6.4), we obtain

$$(6.5) \quad \Phi_-(\underline{k}, f_\tau, z) = -T(\Phi(\underline{k}, f_\tau, z) + \sum_C \Phi_-(\underline{k}, f_{\pi_C(\tau)}, z)\Phi(\underline{k}, f_{\rho_C(\tau)}, z)).$$

Because  $f = f_\tau$  is the output function computed by the flow chart, that is, the label of the outgoing flag of  $\tau$ , we have  $f_{\rho_C(\tau)} = f_\tau$ , since these trees have the same root and outgoing flag with the same label. Thus, the expression above is simply of the form

$$(6.6) \quad \Phi_-(\underline{k}, f_\tau, z) = -T \left( \Phi(\underline{k}, f_\tau, z) \left( 1 + \Phi_-(\underline{k}, \sum_C f_{\pi_C(\tau)}, z) \right) \right).$$

This has the effect of considering not only the pole at  $z = 1$  that is caused by  $f_\tau$  itself, but also those that come from subdivergences created by all the partial recursive functions  $f_{\pi_C(\tau)}$  computed in the intermediate steps of the computation performed by the flow chart  $\tau$ , which are the outputs of the “pruned parts”  $\pi_C(\tau)$  of the admissible cuts on  $\tau$ .

This suggests that the main role of the negative piece of the Birkhoff factorization should be understood as a way of detecting the presence of a pole not only in the function  $\Phi(\underline{k}, f, z)$  itself, but in the analogous functions associated to the results  $f_{\pi_C(\tau)}$  and  $f_{\rho_C(\tau)}$  computed by subroutines of the program  $\tau$  that computes  $f$ .

The role of “subdivergences” in programs computing partial recursive functions is interesting, for example because one can have programs that compute total functions as output, but involve partial functions, so that the graph itself may be divergence-free and still contain subdivergences, as the next example (suggested to us by Joachim Kock) shows.

EXAMPLE 6.1. The identity function can be computed by a tree with one  $\mathfrak{c}$  node and one  $\mu$  node, as a composition of a successor function followed by a partial predecessor function,  $\mu(|y + 1 - x|)$ , undefined at 0 and  $x - 1$  for  $x > 0$ .

Thus, replacing the polar part of  $\Phi(\underline{k}, f, z)$  with the polar part  $\Phi_-(\underline{k}, f_\tau, z)$  can be understood as a way to take into account all possible (and possibly hidden) subdivergences.

It is more difficult to obtain a good interpretation of the positive piece of the Birkhoff factorization. The analogy with quantum field theory suggests that it is not the value itself of  $\Phi_+(\underline{k}, f_\tau, z)$  that matters. After all, even in the original physical setting, there is an intrinsic ambiguity in the renormalized values, coming from the action of the renormalization group. The main question of Manin’s approach here appears to be the following (loosely formulated).

QUESTION 6.2. Is there a new function  $f_{\text{ren}}$ , which is primitive recursive, such that

$$(6.7) \quad \Phi_+(\underline{k}, f_\tau, z) = \Phi(\underline{k}, f_{\text{ren}}, z),$$

namely, is  $\Phi_+$  equal to the original  $\Phi$  of another function, this time primitive instead of partial recursive?

If one could answer this question positively, then one could think of the new function  $f_{\text{ren}}$  as an “extraction of a computable piece” from the partial recursive function  $f$  (and the program  $\tau$  that computes it). Upon close inspection, it is easy

to see that, as naively formulated in Question 6.2, it may not be possible in general to obtain such an  $f_{\text{ren}}$ . However, one can perhaps formulate a better question of whether an identification (6.7) holds after considering the algebraic Feynman rules up to an equivalence relation, as with the gauge transformations on germs of meromorphic functions considered in the quantum field theory setting in [6]. We will not develop this issue further in the present paper, but we leave it as an open question.

6.1.5. *The case of  $\mathcal{H}_{\text{flow}, \mathcal{V}'}$ .* We make here a sketch of a possible suggestion for a variation on the construction above, adapted to define an algebraic Feynman rule on  $\mathcal{H}_{\text{flow}, \mathcal{V}'}$ , where the flags are not labelled and only vertices are. In this case, in order to define something like (6.4), we need a way of assigning inputs to the flow charts. A reasonable choice is to use the basic functions as input. We then define the algebraic Feynman rule as

$$(6.8) \quad \Phi(\underline{k}, \tau, z) := \prod_{\sigma} \Phi(\underline{k}, f_{\tau, \sigma}, z),$$

where the product is over  $\sigma$  ranging over the set of all possible functions  $\sigma : E_{\text{ext}}^{\text{in}}(\tau) \rightarrow \{s, c, \pi\}$  that label the incoming external edges of  $\tau$  by basic functions (successor, constant, or projection). The target of  $\sigma$  contains all the possible basic functions (of type  $s$ ,  $c$ , or  $\pi$ ). Since this is an infinite set (there are infinitely many possible choices of projection functions  $\pi_i^n$ ), if one wants to have a finite product in (6.8) one may have to restrict  $\sigma$  by making a priori choices of smaller sets of basic functions as target, or else introduce further conditions on the class of functions  $\Phi$  that would ensure convergence. The partial recursive function  $f_{\tau, \sigma}$  is the output of the flow chart given by the tree  $\tau$  with inputs assigned by  $\sigma$ . Each  $\Phi(\underline{k}, f_{\tau, \sigma}, z)$  is computed as in (6.4), and (6.8) is extended to arbitrary elements of the Hopf algebra as before. In this setting, the function  $\Phi_{\tau}(\underline{k}, z)$  has a pole at  $z = 1$  iff there is at least a choice of a basic input  $\sigma$  for which  $\underline{k} \notin \mathcal{D}(f_{\tau, \sigma})$ .

With this setting, the BPHZ formula for the algebraic Feynman rule (6.8) becomes more interesting than in the case of (6.6). We have again (6.5), in the form

$$(6.9) \quad \Phi_{-}(\underline{k}, \tau, z) = -T(\Phi(\underline{k}, \tau, z) + \sum_C \Phi_{-}(\underline{k}, \pi_C(\tau), z) \Phi(\underline{k}, \rho_C(\tau), z)),$$

with  $\Phi(\underline{k}, \rho_C(\tau), z)$  and  $\Phi(\underline{k}, \pi_C(\tau), z)$  again computed as in (6.8), but this time it is no longer true that  $f_{\tau} = f_{\rho_C(\tau)}$  because on the tree  $\rho_C(\tau)$  we are using the new inputs given by basic functions and not the input coming from the output of  $\pi_C(\tau)$ . The divergence (6.9) here combines the divergences of the  $f_{\tau, \sigma}$  with divergences coming from partial recursive functions  $f_{\pi_C(\tau), \sigma}$  and  $f_{\rho_C(\tau), \sigma}$  that flowcharts  $\pi_C(\tau)$  and  $\rho_C(\tau)$  compute starting from inputs of basic functions.

**6.2. Dyson–Schwinger equations for the halting problem.** Finally, we tentatively propose a possible role for Dyson–Schwinger equations in the context of Manin’s approach to the halting problem. These remarks are also, at this stage, very speculative, as they would depend for a more precise interpretation upon a better understanding of the main questions about the Birkhoff factorization discussed earlier in this section.

If we work with the Feynman rule described above on the Hopf algebra  $\mathcal{H}_{\text{flow}, \mathcal{V}'}$  with only vertex decorations, then we can readily consider Dyson–Schwinger equations as discussed in §4 above. However, if we work with the definition of the algebraic Feynman rule  $\Phi$  described in (6.4) on the Hopf algebra  $\mathcal{H}_{\text{flow}, \mathcal{R}'}$  of flow charts with flag decorations, we need to extend the grafting operators appropriately that are needed to define Dyson–Schwinger equations from  $\mathcal{H}_{\text{flow}, \mathcal{V}'}$  to  $\mathcal{H}_{\text{flow}, \mathcal{V}'}$ . This can be done by assigning as output the empty function  $f = \emptyset$ , mapped to the constant function  $\Phi(\underline{k}, \emptyset, z) \equiv 1$ , as the output of any tree  $\tau$  that is obtained through a grafting where the flag labels don’t match.

We can then consider a Dyson–Schwinger equation of the form (4.3), with  $B^+ = \sum_{\delta} B_{\delta}^+$ , with  $\delta \in \{\mathfrak{b}, \mathfrak{c}, \mathfrak{r}, \mu\}$  or a more general system of Dyson–Schwinger equations (4.5), with the condition that the components  $x_m$  of the unique solution span a Hopf subalgebra of  $\mathcal{H}_{\text{flow}, \mathcal{R}'}$ . In such case, we can restrict the algebraic Feynman rule  $\Phi$  constructed above to this Hopf subalgebra and still perform the BPHZ renormalization.

One can then reformulate the same questions that we have briefly discussed regarding Manin’s approach to the halting problem, of how to interpret the resulting negative and positive parts of the Birkhoff factorization. In this case, the fact that one is restricting attention to the Hopf subalgebra generated by the components of the solution to the Dyson–Schwinger equation means that the negative part  $\Phi_{-}(\underline{k}, f_{x_m}, z)$  (or  $\Phi_{-}(\underline{k}, f_{x_m, \sigma}, z)$  in the case of  $\mathcal{H}_{\text{flow}, \mathcal{V}'}$ ) will now account for subdivergences that belong to the same Hopf subalgebra determined by the Dyson–Schwinger equation. This can be thought of, heuristically, as measuring the “amount of non-computability” that can be produced by sub-flow-charts that are characterized by a certain self-similarity property (defined by the Dyson–Schwinger equation). An analog in this setting of Question 6.2 on the interpretation of the positive part of the Birkhoff factorization can also be formulated in the Hopf subalgebra determined by the Dyson–Schwinger equation.

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