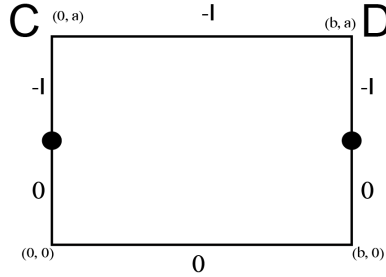


Physics 106c – Problem Set 1 — Solutions

April 26, 2012

Problem 1



With charge conservation and Maxwell equation, we can write down $\nabla \times \vec{K} = \frac{\partial K_x}{\partial x} + \frac{\partial K_y}{\partial y} = 0$ and $\nabla \vec{E} = 0$ and introduce ϕ here.

$$K_x = -\frac{\partial \phi}{\partial y}, \quad K_y = \frac{\partial \phi}{\partial x} \quad (1)$$

Plug K_x and K_y into the $\nabla \vec{E} = \nabla(\rho \vec{K}) = 0$ and one will get $\nabla \phi^2 = 0$. The boundary condition (as in the figure) can be obtained by $I = \int_0^a K_x dy$. The problem is solved in last term HW2. After some algebra, the ϕ can be solved as the following. For $\frac{a}{2} < y < a$

$$\phi(x, y) = -I + \frac{2I}{\pi} \sum_{n \text{ odd}} \frac{\sin(\frac{n\pi x}{b}) \sinh(\frac{n\pi(a-y)}{b})}{n \sinh(\frac{n\pi a}{2b})} \quad (2)$$

for $0 < y < a/2$,

$$\phi(x, y) = -\frac{2I}{\pi} \sum_{n \text{ odd}} \frac{\sin(\frac{n\pi x}{b}) \sinh(\frac{n\pi y}{b})}{n \sinh(\frac{n\pi a}{2b})} \quad (3)$$

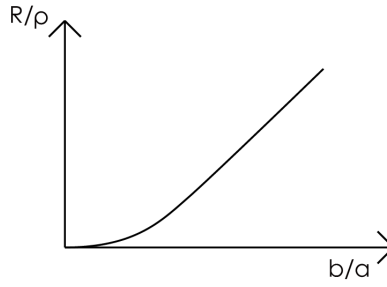
So the voltage $V_{cd} = \int_0^b E_x dx$ is

$$V_{cd} = \rho \frac{4I}{\pi} \sum_{n \text{ odd}} \frac{1}{n \sinh(\frac{n\pi a}{2b})} \quad (4)$$

and the resistance R equals

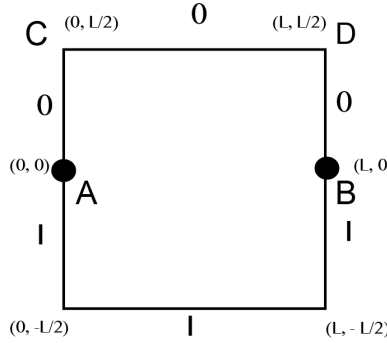
$$R_{cd} = \frac{4\rho}{\pi} \sum_{n \text{ odd}} \frac{1}{n \sinh(\frac{n\pi a}{2b})} \quad (5)$$

The figure below is the plot of R_{cd} vs b/a . In the limit $b/a \rightarrow \infty$, $R_{cd} \rightarrow \rho \frac{4}{\pi} \frac{2b}{\pi a} \sum_{n \text{ odd}} \frac{1}{n^2} = \rho \frac{8b}{\pi^2 a} \frac{\pi^2}{8} = \rho b/a$.



Problem 2

To determine the boundary condition for ϕ , one can start from the boundary condition of K . (i) $K_x(0, y) = K_x(L, y) = 0$, for $y \neq 0$, $y \in [L/2, -L/2]$. (ii) $K_y(x, L/2) = K_y(x, -L/2) = 0$, for $x \in [0, L]$. (iii) $I = \int_{-L/2}^{L/2} K_x dy$. With $K_x = -\frac{\partial \phi}{\partial y}$ and $K_y = \frac{\partial \phi}{\partial x}$, one can set $\phi = 0$ in upper plane and $\phi = I$ in lower plane.



Since there's no time-varying B field, and $E_x = \rho_{xx}K_x$, $E_y = \rho_{yy}K_y$, one will get

$$\frac{\partial K_y}{\partial x} - \frac{\rho_{xx}}{\rho_{yy}} \frac{\partial K_x}{\partial y} = 0 \quad (6)$$

$$\frac{\partial^2 \phi}{\partial x^2} - \frac{\rho_{xx}}{\rho_{yy}} \frac{\partial^2 \phi}{\partial y^2} = 0 \quad (7)$$

Here we can rescale the $y' = \sqrt{\frac{\rho_{yy}}{\rho_{xx}}} y = \eta y$. From Problem 1, we can immediate have the ϕ which is.

$$\phi(x, y) = -\frac{I}{\pi} \sum_{n \text{ odd}} \frac{\sin(\frac{n\pi x}{L}) \sinh(\frac{n\pi \eta (L/2 - y)}{L})}{n \sinh(\frac{n\pi \eta}{2})} \quad (8)$$

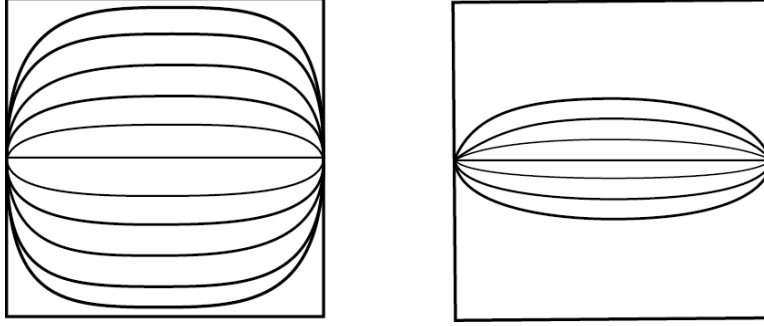
The potential differences V_{CD} and R are

$$V_{CD} = \sqrt{\rho_{xx}\rho_{yy}} \frac{4I}{\pi} \sum_{n \text{ odd}} \frac{1}{n \sinh(\frac{n\pi \eta}{2})} \quad (9)$$

$$R_{CD} = \sqrt{\rho_{xx}\rho_{yy}} \sum_{n \text{ odd}} \frac{4}{n\pi \sinh(\frac{n\pi \eta}{2})} \quad (10)$$

When $\rho_{xx} = 1$, $\rho_{yy} = 0.2$, $R_1 = 0.802658$, In another case, $\rho_{xx} = 0.2$, $\rho_{yy} = 1$, $R_1 = 0.0340057$. So there is a difference between R_1/R_2 , ρ_{xx}^1/ρ_{xx}^2 and ρ_{yy}^1/ρ_{yy}^2 . For the first case (the left figure), since ρ_{xx} is larger than ρ_{yy} current wants to flow to y direction more than in the second case(the

right figure). So the current which entered in plate goes longer path, and as a result effectively it feels more resistance. Therefore, considering such phenomena, R_1/R_2 is bigger than ρ_{xx}^1/ρ_{xx}^2 which we just consider x -direction.



Problem 3

$$\vec{J} = \sigma(\vec{E} + \vec{v} \times \vec{B}), \quad \vec{B} = (0, 0, B)$$

$$j_i = \sigma(E_i + \epsilon_{ijk} v_j B_k), \quad \vec{j} = ne\vec{v} \Rightarrow v_i = \frac{j_i}{ne}$$

$$\begin{cases} j_1 = \sigma E_1 + \epsilon_{123} v_2 B \\ j_2 = \sigma(E_2 + \epsilon_{213} v_1 B) \end{cases}$$

↓

$$\begin{cases} j_1 = \sigma E_1 + \frac{\sigma B}{ne} j_2 \\ j_2 = \sigma E_2 - \frac{\sigma B}{ne} j_1 \end{cases}$$

↓

$$j_1 = \sigma E_1 + \frac{\sigma B}{ne} (\sigma E_2 - \frac{\sigma B}{ne} j_1)$$

$$j_1 \left(1 + \left(\frac{\sigma B}{ne}\right)^2\right) = \sigma \left(E_1 + \frac{\sigma B}{ne} E_2\right)$$

$$j_1 = \frac{1 \cdot \sigma}{1 + \left(\frac{\sigma B}{ne}\right)^2} \left(E_1 + \frac{\sigma B}{ne} E_2\right)$$

$$\begin{aligned} j_2 &= \sigma E_2 - \frac{\sigma B}{ne} j_1 = \sigma E_2 - \frac{\sigma}{1 + \left(\frac{\sigma B}{ne}\right)^2} \left(\frac{\sigma B}{ne} E_1 + \frac{\sigma B}{ne} E_2\right) = \\ &= \frac{\sigma}{1 + \left(\frac{\sigma B}{ne}\right)^2} \left(E_2 - \frac{\sigma B}{ne} E_1\right) \end{aligned}$$

$$\Rightarrow \begin{pmatrix} j_1 \\ j_2 \end{pmatrix} = \frac{\sigma}{1 + \left(\frac{\sigma B}{ne}\right)^2} \begin{pmatrix} 1 & \frac{\sigma B}{ne} \\ -\frac{\sigma B}{ne} & 1 \end{pmatrix} \begin{pmatrix} E_1 \\ E_2 \end{pmatrix}$$

$$\Rightarrow \begin{pmatrix} \epsilon_{xx} & \epsilon_{xy} \\ \epsilon_{yx} & \epsilon_{yy} \end{pmatrix} = \frac{\sigma}{1 + \left(\frac{\sigma B}{ne}\right)^2} \begin{pmatrix} 1 & \frac{\sigma B}{ne} \\ -\frac{\sigma B}{ne} & 1 \end{pmatrix}.$$

$$3. \quad \sigma_{xx} = \frac{\sigma}{1 + \left(\frac{\sigma B}{ne}\right)^2}$$

$$\begin{pmatrix} \rho_{xx} & \rho_{xy} \\ \rho_{yx} & \rho_{yy} \end{pmatrix} = \begin{pmatrix} \sigma_{xx} & \sigma_{xy} \\ \sigma_{yx} & \sigma_{yy} \end{pmatrix}^{-1} = \begin{pmatrix} \frac{1}{\sigma} & -\frac{B}{ne} \\ \frac{B}{ne} & \frac{1}{\sigma} \end{pmatrix}$$

$\rho_{xx} = \frac{1}{\sigma}$, If we fix B (finite), let $\sigma \rightarrow \infty$

ρ_{xx}, σ_{xx} both $\rightarrow 0$

$\rho_{xx} \rightarrow 0$ just because resistivity is inverse of conductivity

$\sigma_{xx} \rightarrow 0$ because when $\sigma \rightarrow 0$, from $\vec{J} = \sigma \left(\vec{E} + \frac{\vec{J}}{ne} \times \vec{B} \right)$

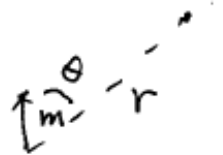
$\Rightarrow \vec{E} + \frac{\vec{J} \times \vec{B}}{ne} \rightarrow 0 \Rightarrow \vec{J} \perp \vec{E} \Rightarrow$ diagonal components of

$\vec{J}$ must vanish.

Problem 4

In free space, we have

$$\vec{B}_0 = \frac{\mu_0 \vec{m}}{4\pi r^3} (2\cos\theta \hat{r} + \sin\theta \hat{\theta})$$



For linear magnetic materials, H doesn't change

$$\begin{aligned}\vec{B} &= \frac{\vec{B}_0 H}{H_0} = \vec{B}_0 (H \chi_m) \\ &= \frac{\mu_0 (H \chi_m)}{4\pi r^3} (2\cos\theta \hat{r} + \sin\theta \hat{\theta})\end{aligned}$$